



# Ferroelectric Domain Walls Physics and Function

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# Selected References

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Disclaimer: This topic is a rich one! There is not time cover all of it in detail, some great work is thus not included here. No offense is meant, just a result of the time limits! For those new to the field, know there is more beyond these slides to explore...



# Question?

Imagine we are walking around the Corsican countryside and find a piece of magnetite ( $\text{Fe}_3\text{O}_4$ ) which is magnetic (specifically ferrimagnetic). Could we stick this piece of  $\text{Fe}_3\text{O}_4$  to the refrigerator?



Why not?

Domains – One of the most important considerations for property evolution in ferroics!





# Ferroic Materials

- **Ferroic crystals** possess two or more orientation states or **domains** and under the right driving force, the **domain walls** move, switching the crystal from one domain state to another
  - Force can be stress ( $X$ ), electric field ( $E$ ), magnetic field ( $H$ ), or a combination thereof
- **Ferroelectric**  $\rightarrow$  Spontaneous polarization  $P_s$
- **Ferroelastic**  $\rightarrow$  Spontaneous strain  $x_s$
- **Ferromagnetic**  $\rightarrow$  Spontaneous magnetization  $M_s$
- Not necessary that the orientation states differ in primary quantities ( $x$ ,  $P$ ,  $M$ ) for the appropriate field to develop a driving force for domain motion
  - Example  $\rightarrow$  Domains have different orientations of elastic compliance tensor, suitably chosen stress can produce different strains in 2 domains
  - Same stress may act upon the difference in induced strain to produce wall motion, domain reorientation
  - This is called *ferrobielastic* response as opposed to *ferroelasticity*



# What is a Ferroelectric Material?

## The Hierarchy of Materials: Symmetry

32 Crystal Classes → 11 possess center of symmetry → No polar properties



Of remaining 21 → 20 exhibit polarity when subjected to a stress (*Piezoelectric*)



Of these 20 → 10 show unique polar axis → possess spontaneous polarization (*Polar*)

- Temperature dependence → changes in polarization result in flow of charge to-and-from surface (*Pyroelectric*)
- When a *polar* material has two or more orientation states in the absence of an  $E$  and can be shifted from one to another of these states by an electric field → *Ferroelectric*

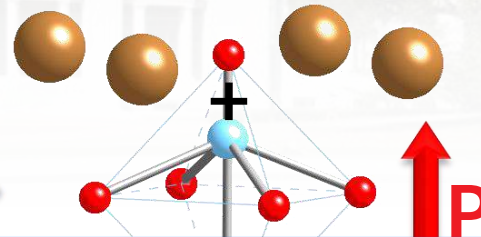
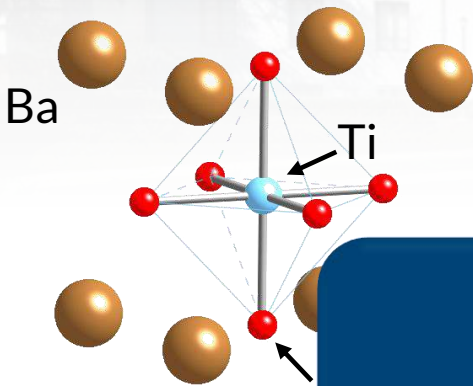


# What is a Ferroelectric Material?

Paraelectric

Ferroelectric

Switchable Polarization



Energy

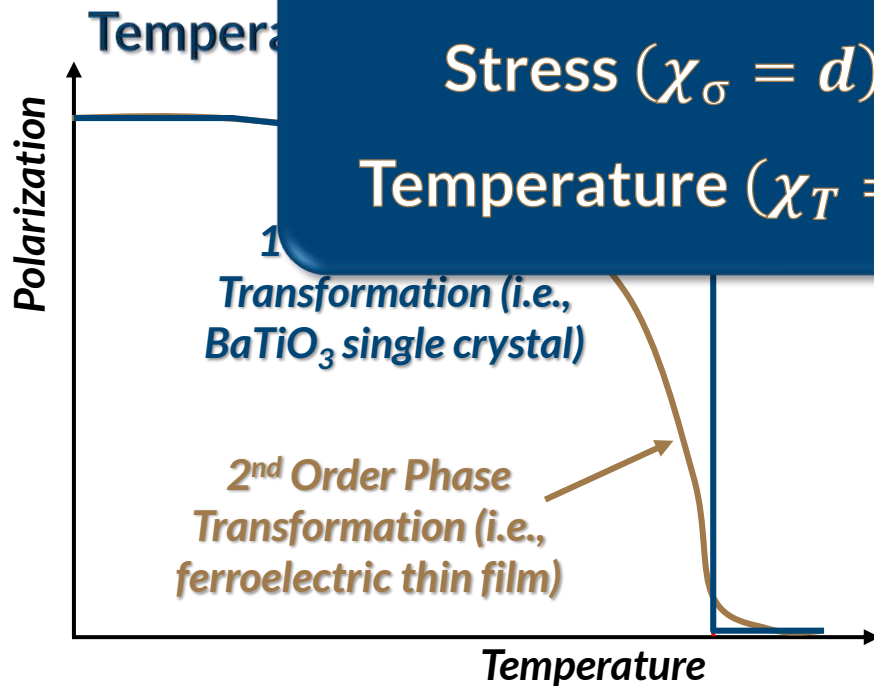
Displacement

$$\text{Susceptibility } \chi_{\alpha} = \frac{dP}{d\alpha}$$

Electric Field ( $\chi_E = \epsilon$ )  $\rightarrow$  Dielectric

Stress ( $\chi_{\sigma} = d$ )  $\rightarrow$  Piezoelectric

Temperature ( $\chi_T = \pi$ )  $\rightarrow$  Pyroelectric



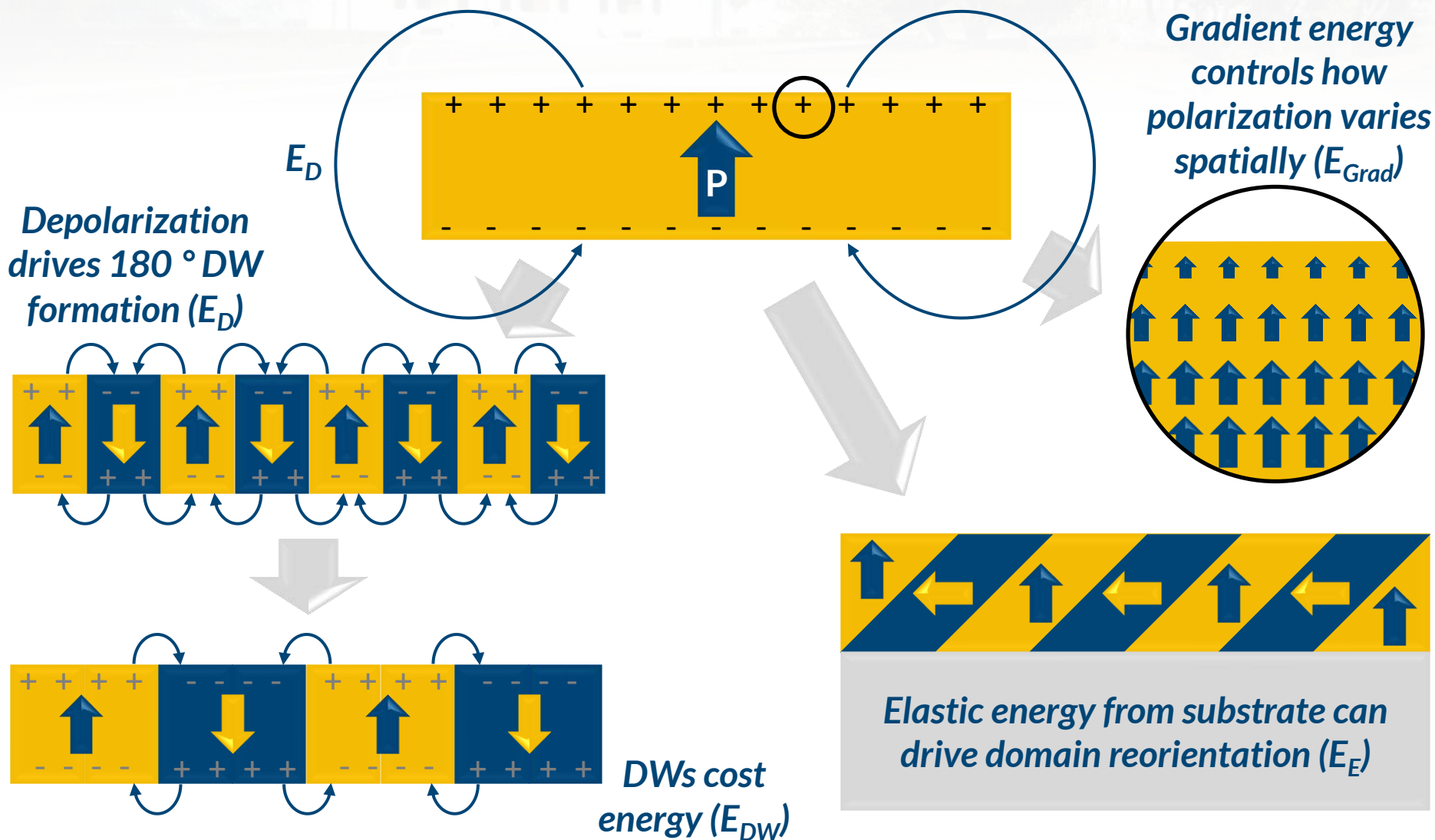
Electric-Field Dependence





# Energies in Ferroelectrics

$$E_{Total} = E_P + E_D + E_{DW} + E_{Grad} + E_E$$





# Ferroelectric Domains

- We are concerned with how polarization changes at inhomogeneities...
- Recall that  $\mathbf{D} = \epsilon_0 \mathbf{E} + \mathbf{P}$  and that  $\mathbf{P}$  arises from both the polarizability of the material in the presence of field  $\mathbf{P}_E = \chi \mathbf{E}$  and from the spontaneous alignment of dipoles ( $\mathbf{P}_S$ )
- If there is a spatial variation of the  $\mathbf{D}$  the free charge density  $\rho$  must satisfy Poisson's Equation  $\rightarrow \nabla \mathbf{D} = \rho$  so that  $\nabla \mathbf{E} = \frac{1}{\epsilon \epsilon_0} (\rho - \nabla \mathbf{P}_S)$
- In an ideal infinite ferroelectric  $\mathbf{P}_S$  is uniform so that  $\nabla \mathbf{E} = \frac{\rho}{\epsilon \epsilon_0}$  as in ordinary dielectrics
- At a surface where  $\mathbf{P}_S$  decreases to zero (or near defects where  $\mathbf{P}_S$  may differ from perfect xtal)  $\nabla \mathbf{P}_S$  acts as the source for a  $\mathbf{E}_D$
- This field can be compensated by the flow of free charge within the crystal  $\rightarrow \rho = \int_0^t \sigma \cdot \mathbf{E} dt$  (where  $\sigma$  is the electrical conductivity)





# Ferroelectric Domains

- Alternatively, free charge in the surrounding medium can compensate for the depolarization field at the crystal surface
  - In air this is usually slow compared to crystal conductivity
- Accumulation of surface charge by conduction satisfies the requirement that the  $\mathbf{E}$  vanish both outside the crystal and inside the bulk (but not necessarily just below the crystal surface)
- The energy associated with this depolarization field is  $W_E = \int_V \mathbf{D} \cdot \mathbf{E} dV$ 
  - Energy is zero for a totally compensated crystal in  $\infty$
  - In an insulating xtal (and environment),  $\infty$  is reached very slowly and at short times, the energy  $W_E$  can reach very high values
- When we undergo the FE transition, there are at least 2 equivalent directions along which the spontaneous polarization may occur
- In order to minimize  $W_E$  different regions of the xtal polarize in each of these directions



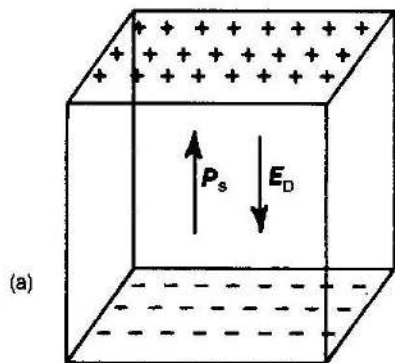
# Ferroelectric Domains

- Each volume of uniform polarization is a *domain*
- The depolarizing field ( $E_D$ ) that arises during cooling from PE phase is usually enough to *prevent any polarization in a virgin crystal*
- Boundaries separating domains are domain walls  $\rightarrow$  they have an energy associated with them ( $E_{DW}$ )
- Final domain configuration is determined by minimizing an appropriate free energy including both of these energy terms
- Final domain structures rely on things like...
  - crystal symmetry | electrical conductivity | defect structure
  - magnitudes of the polarization and elastic and dielectric compliances
  - history of the crystal preparation | sample geometry | etc.



# Ferroelectric Domains

- The onset of spontaneous polarization results in the appearance of surface charge density and an accompanying depolarization field ( $E_D$ )



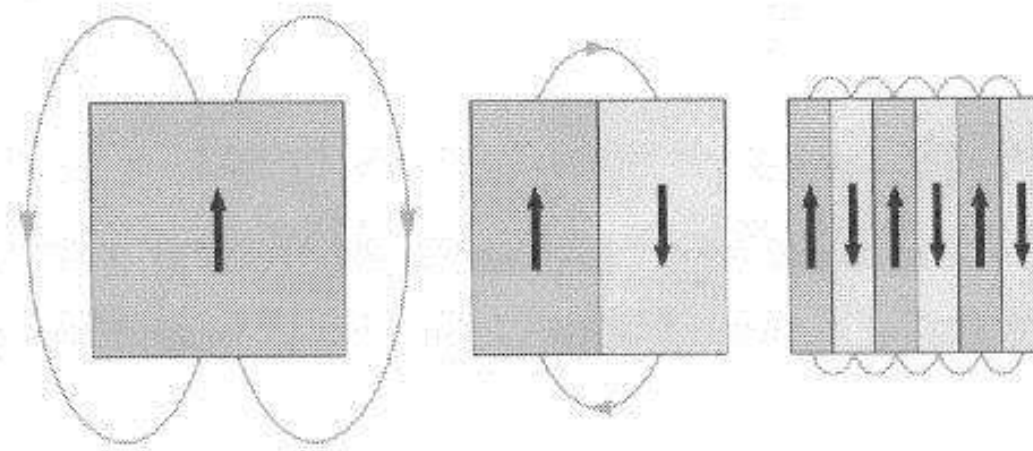
- Thus we can write  $E_{Tot} = E_{Ext} + E_D$  where  $E_D$  points opposite to  $E_{Ext}$
- The energy associated with the  $P$  in the  $E_D$  can be minimized by driving the crystal into regions of oppositely polarized regions  $\rightarrow$  Domains
- Domains result in a reduction in  $E_D$  and in electrostatic energy for the system, but cost the energy to create a domain wall (generally  $180^\circ$  domain walls)





# Ferroelectric Domains

- Configuration of domains follows a head-to-tail condition in order to avoid discontinuities in the polarization at the domain boundary  $\rightarrow \nabla \vec{P} = \sigma$



- Build-up of domain walls, elastic stress fields, and free charge carriers all counteract the process of domain formation (don't forget vacancies, defects, dopants, etc.)
- A single domain structure can be achieved by applying the right  $E \rightarrow$  domains aligned along the  $E$  direction grow at the expense of the others



# Ferroelectric Domains: Free Energy

- Static domain configuration of a FE is obtained by minimizing the total free energy of the crystal including the energy associated with the crystal surfaces and domain walls

- We can write:  $G_1 = G_1^0 + \int \left( \frac{\alpha}{2} D^2 + \frac{\gamma}{4} D^4 \right) dV + W_W + W_E$

- Where we neglect higher order powers

- **Depolarization Energy ( $W_E$ )**

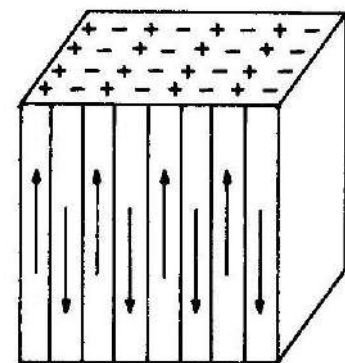
- To minimize the depolarization energy at the crystal surfaces before compensation by free charge takes place domains form throughout the volume of the crystal
- Without domain formation the energy associated with the crystal surfaces can be greater than the energy associated with the ferroelectric ordering → could turn FE order off!
- Magnitude of  $W_E$  depends on the crystal and domain geometry



# Ferroelectric Domains: Free Energy

## • Domain Wall Energy ( $W_W$ )

- If the energy per unit area of a domain wall is  $\sigma$  then for a domain geometry like this  $\rightarrow W_W = (\sigma/d)V$
- Minimizing the energy  $W_E + W_W$  yields the =m value of the domain width as  $d = \left(\frac{\sigma t}{\epsilon^* P_0^2}\right)^{1/2}$  where  $P_0$  must be obtained from self-consistent solution free energy
- Domains of  $d$  greater than this are suppressed by the  $W_E$ , domains smaller are suppressed by the  $W_W$
- $d$  varies quadratically with thickness  $\rightarrow$  thickness  $\downarrow$ ,  $d$  approaches domain thickness  $\rightarrow W_E$  can no longer be minimized by domain formation, FE may cease to exist
- Contributions to  $W_W \rightarrow$ 
  - $W_E$  due to the  $\nabla P$  at domain boundaries
  - Dipolar energy (misalignment of dipoles on either side of DW)
  - Elastic energy

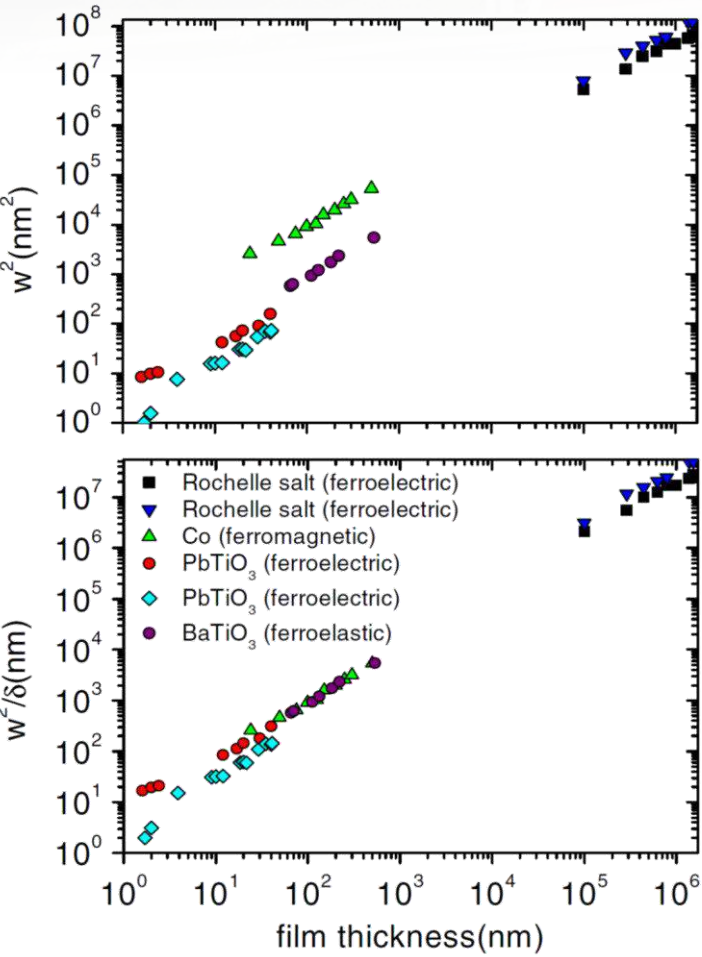






# Ferroelectric Domains: Kittel's Law

• Kittel's law  $\rightarrow$  Domain width  $w = \sqrt{\frac{\sigma}{U}} d$  where  $\sigma$  is the energy density per unit area of the wall,  $U$  is the volume energy density of the domain, and  $d$  is the thickness



- Comparisons between stripe domains of different ferroic materials shows...
  - All of them scale with the same *square root dependence* of domain width on film thickness;
  - Kittel's law holds true for ferroelectrics down to small thickness;
  - When the square of the domain size is normalized by the domain-wall thickness, the different ferroics fall on the same curve

**One can engineer domain structures (periodicities) with film thickness**

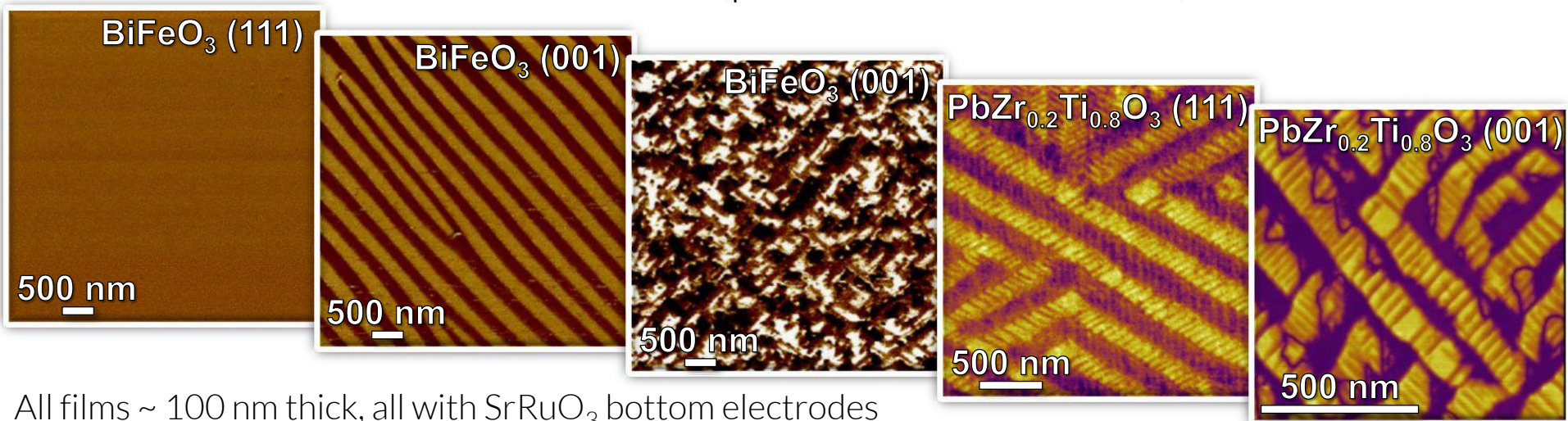
Catalan et al., J. Mater. Sci. 44, 5307 (2009)



# Ferroelectric Domains: Result

Combining all effects...

- Domain-wall thickness: 1-10 nm (2-20 unit cells)  $\rightarrow$  100s nm in FM
- Energy of domain wall:  $\sim 10$  erg/cm<sup>2</sup> ( $10^{-6}$  J/cm<sup>2</sup>)  $\rightarrow$  1-10 erg/cm<sup>2</sup> in FM
- FE domain walls are much less wide than magnetic domain walls  $\rightarrow$ 
  - Magnetic exchange energy is much larger than the elastic energy and slow rotation of the magnetization vector occurs over hundreds of unit cells
  - Magnetization magnitude remains constant, FE polarization reduces to zero in some domain walls (paraelectric character)



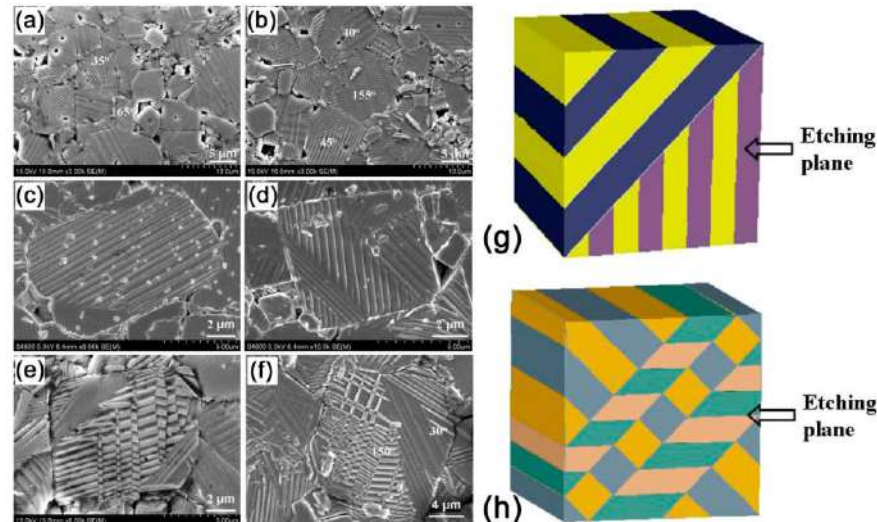
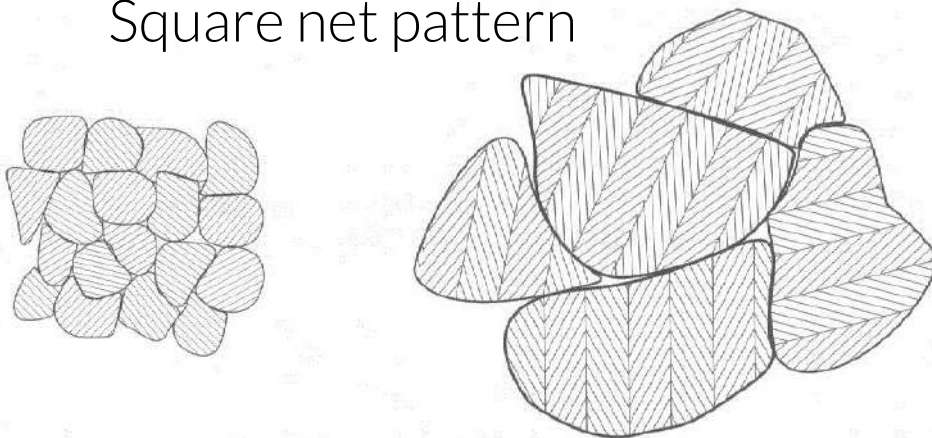
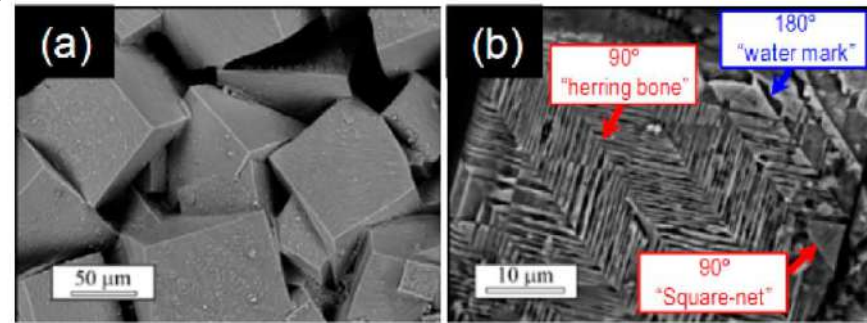




# Ferroelectric Domains

- Polycrystalline, bulk ceramics → pattern of domains is quite different because the domain structure in each grain is formed under elastic, clamped conditions by surrounding domains
- Domain Wall Types in Single Xtals
  - Tetragonal – 90° (ferroelastic), 180° (ferroelectric)
  - Rhombohedral – 71°, 109° (ferroelastic), 180° (ferroelectric)
- Domain Structures in Poly-Xtals
  - Herringbone (more common) and Square net pattern

Alikin et al., Materials 10, 47 (2017)

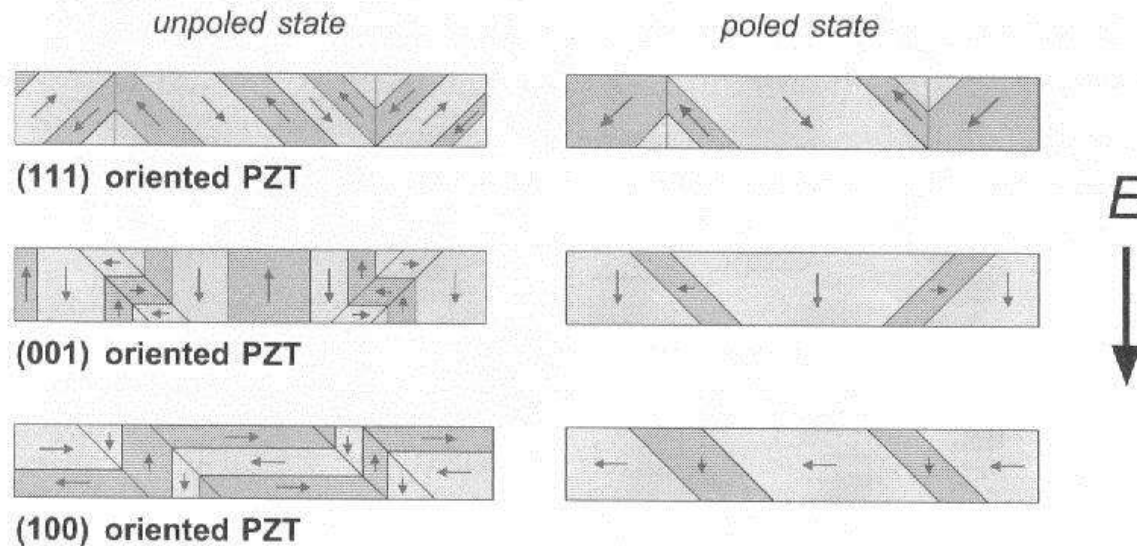




# Ferroelectric Domains

## Thin Film Domain Structures

- Ultra-thin films – Single-domain
- Films  $> \sim 100$  nm – Polydomain
  - Example tetragonal  $\text{PbZr}_{1-x}\text{Ti}_x\text{O}_3$

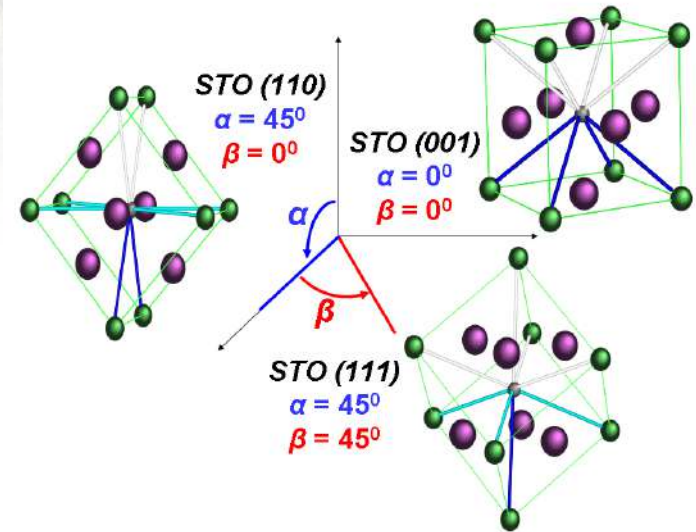


- Varies with film/substrate orientation, applied field, etc.
- Predominant structure consists of  $90^\circ$

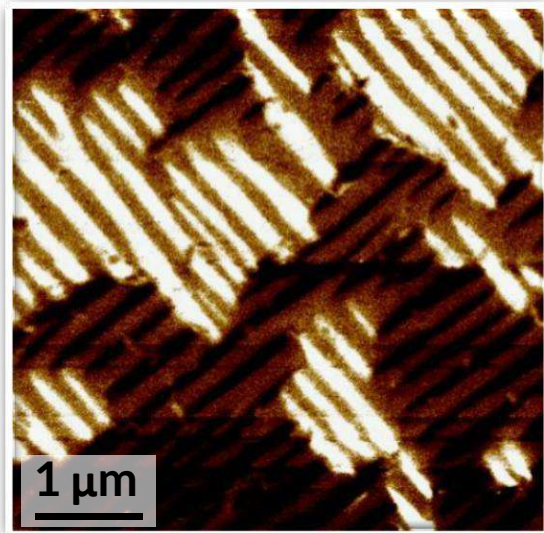
# Ferroelectric Domains

## Thin Film Domain Structures

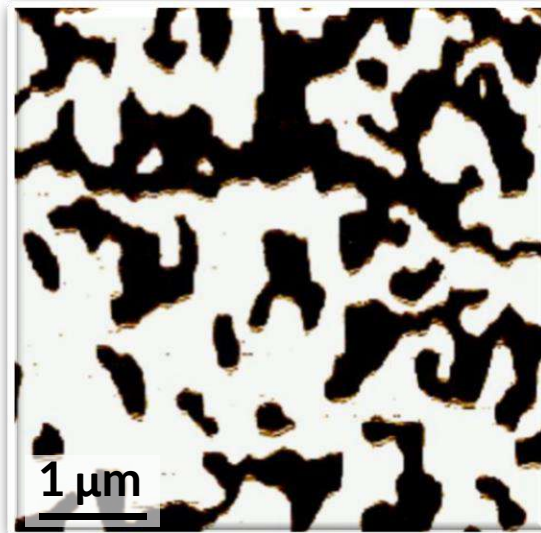
- Example: rhombohedral  $\text{BiFeO}_3$ 
  - (001) – Predominantly  $71^\circ$  domain walls (could have all types)
  - (110) – Predominantly  $71^\circ$  domain walls (could have all types)
  - (111) – Can only have  $180^\circ$  domain walls



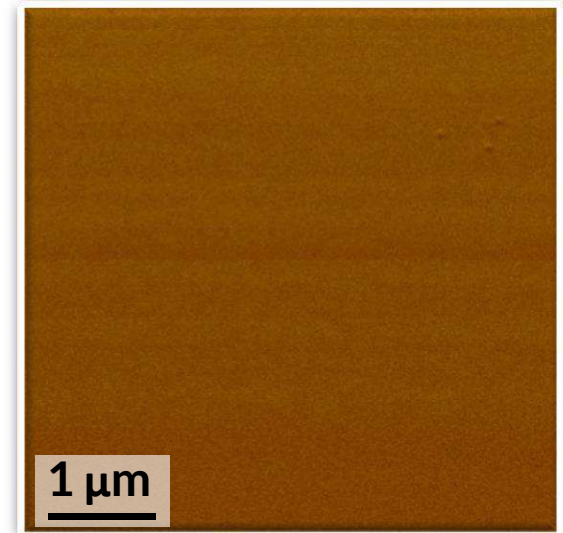
$\text{BiFeO}_3 / \text{SrTiO}_3$  (001)



$\text{BiFeO}_3 / \text{SrTiO}_3$  (110)



$\text{BiFeO}_3 / \text{SrTiO}_3$  (111)





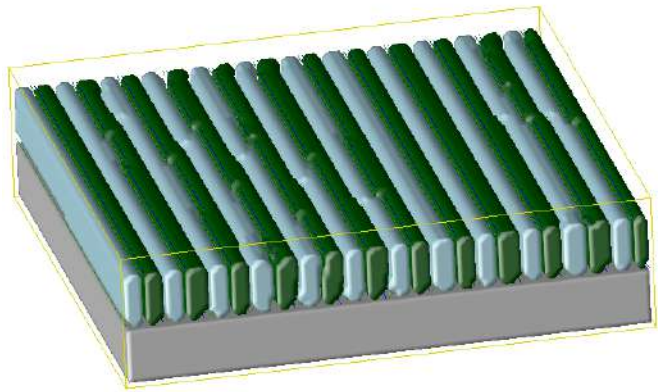


# Controlling Ferroelectric Domain

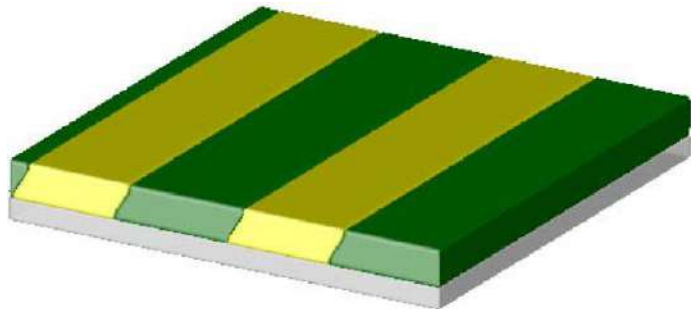
## Interplay between Strain and Electrostatics

S. K. Streiffer, et al. *J. Appl. Phys.* **83**, 2742 (1998)

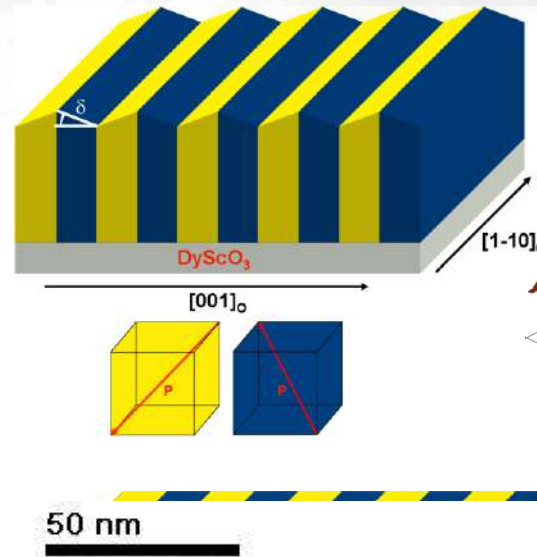
Phase Field Models  
(Yulan Li, Long-Qing Chen,  
Penn State)



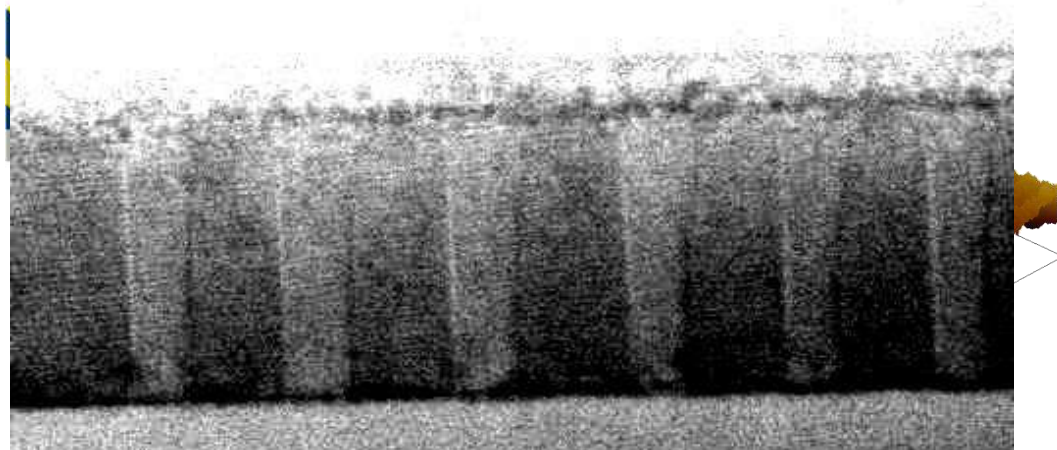
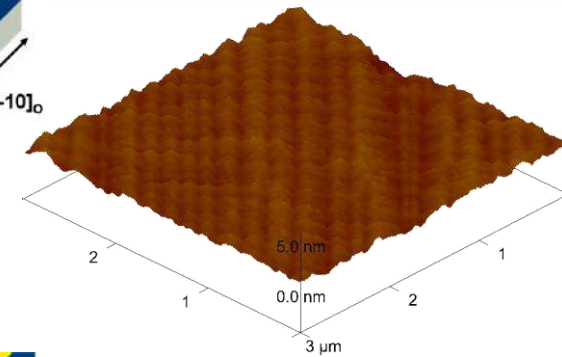
*Open circuit boundary condition*



*Short circuit boundary condition*



*109° Domain Wall  
Arrays*

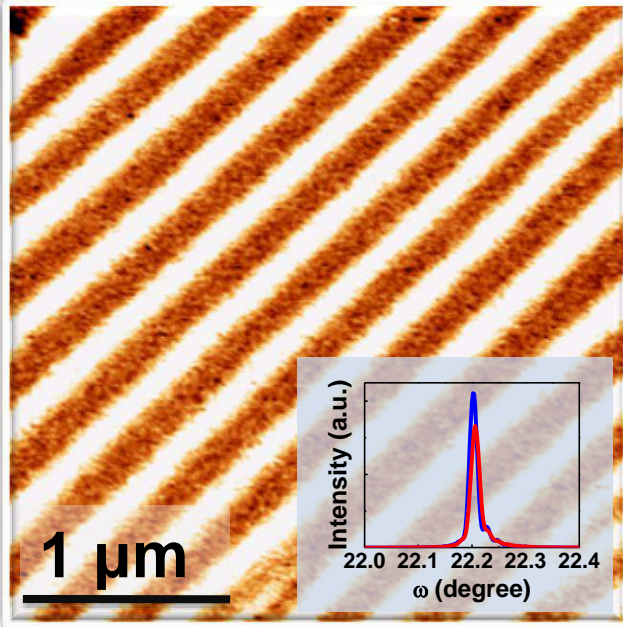


Chu et al., *Nano Lett.* **9**, 1726 (2009)



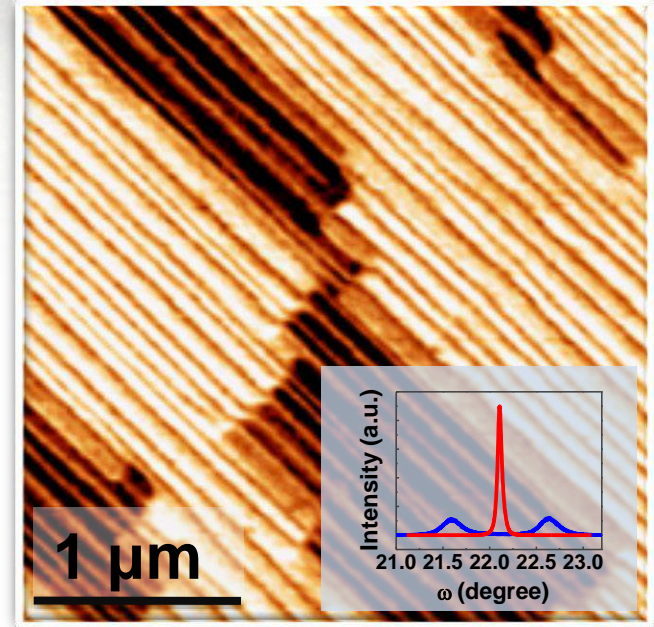
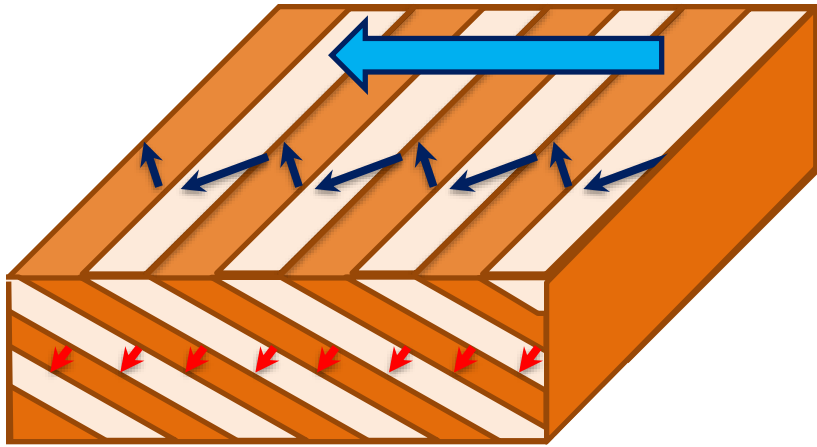


# Model Systems – Ordered DW Arrays



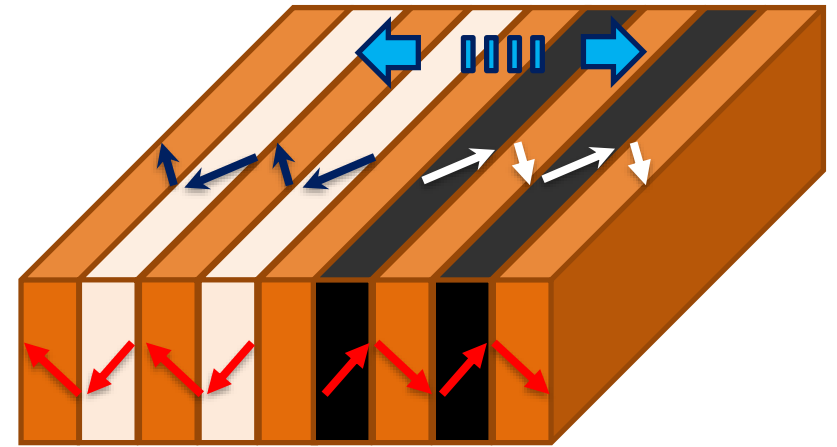
71° Domain Walls

$\vec{P}_{\text{net}}$



109° Domain Walls

$\vec{P}_{\text{net}}$







# Engineering Domain Structures: $\text{BiFeO}_3$

- $\text{BiFeO}_3$**
- Film thickness
  - Growth mechanism
  - Growth rate

$\text{BiFeO}_3$

- Bottom Electrode**
- Material
  - Film thickness
  - Growth mechanism

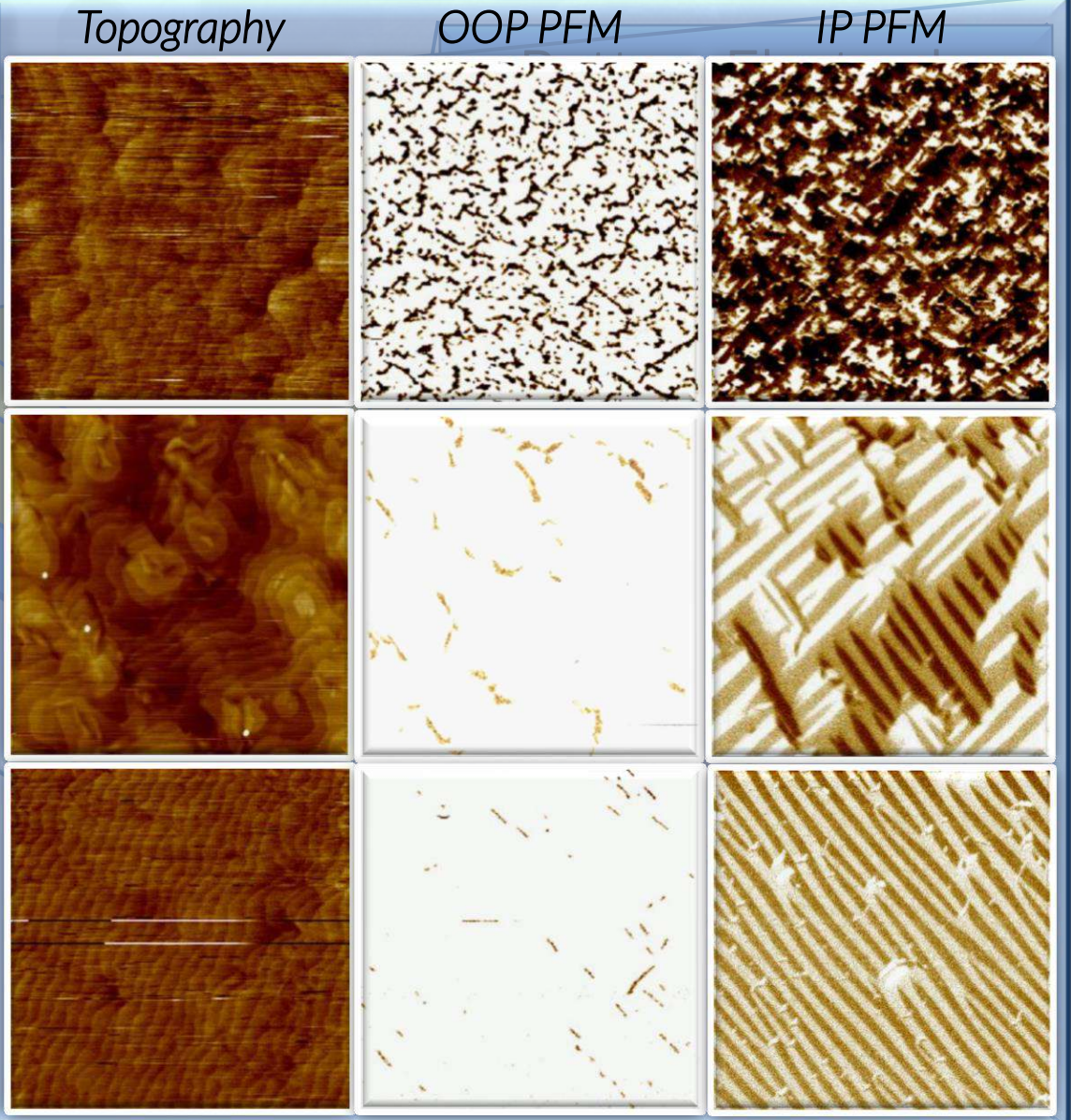
- Substrate**
- Material
  - Orientation
  - Miscut

- Controlling domain structures in BFO**
- $\text{BiFeO}_3$  – Growth rate, underlying bottom electrode structure
  - $\text{SrRuO}_3$  – Structural variants
  - **Substrates** - vicinality
    - Miscut

Island Growth

Step Bunching

Step Flow

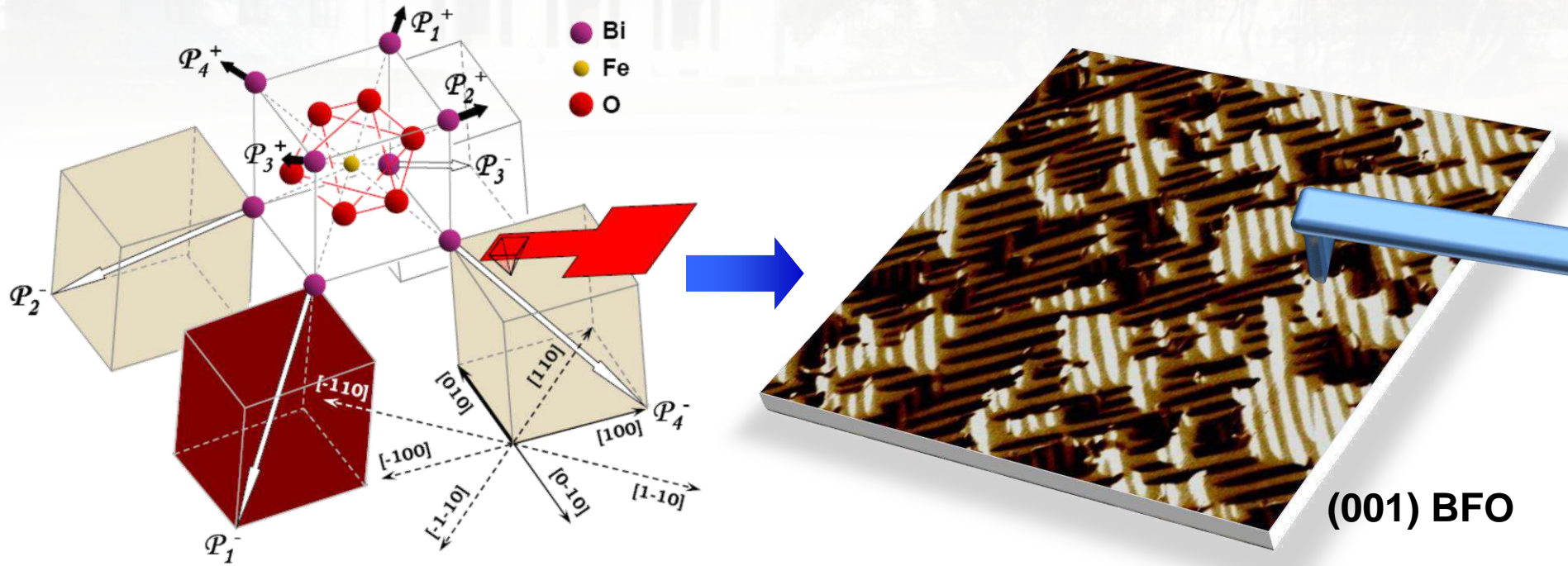






# Imaging Ferroelectric Domains

## Piezoresponse Force Microscopy (PFM)



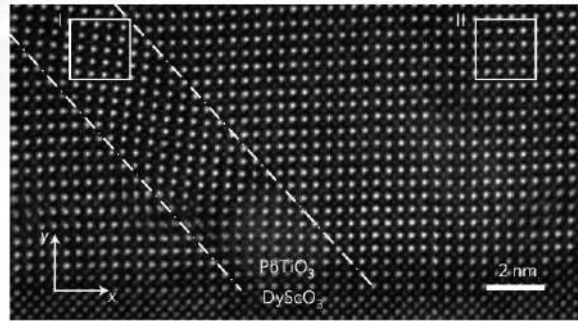
OOP Contrast	IP Contrast		Possible Orientations
	[010]	[0-10]	
Dark (up)	Dark (left)	Light (right)	$[-111]$ , $[-1-11]$
Light (down)	Light (right)	Dark (left)	$[1-1-1]$ , $[11-1]$

S. Kalinin, A. Gruverman, *Scanning Probe Microscopy* (Vol II), Springer: New York (2007); Kalinin et al., *IEEE Trans. Ultrason. Ferroelect. Freq. Cont.* **53**, 2226 (2006); Soergel *J. Phys. D: Appl. Phys.* **44**, 464003 (2011); Gruverman et al., *Nature Commun.* **10**, 1661 (2019)



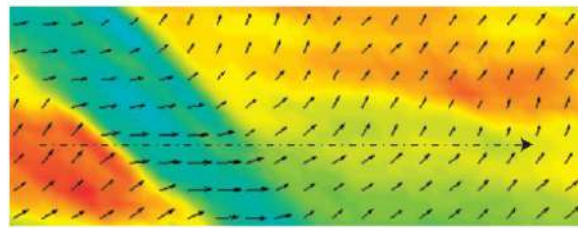
# Direct Imaging of Domain Walls

## (Scanning) Transmission Electron Microscopy (STEM)



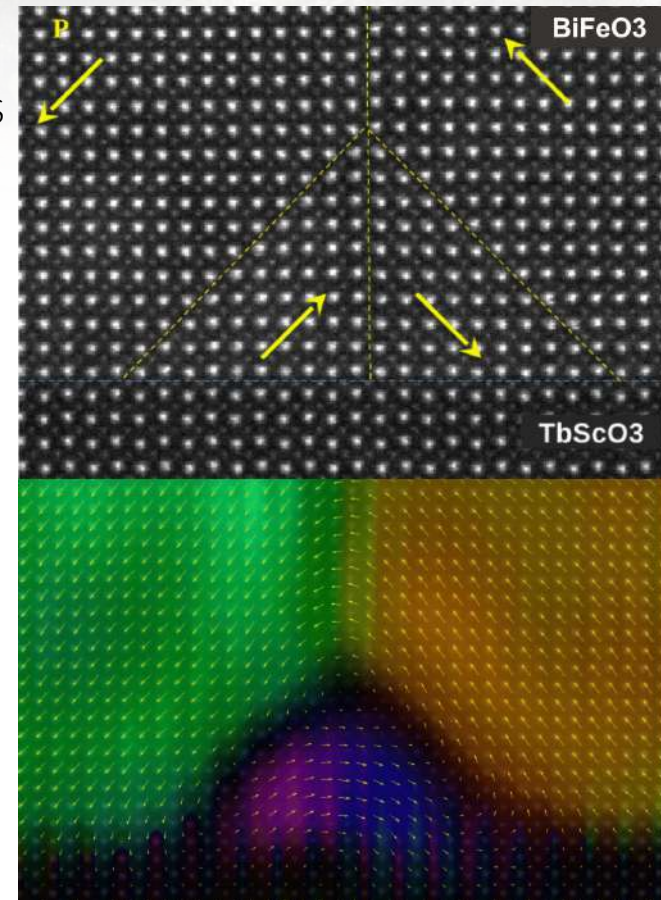
PbTiO<sub>3</sub> → Strain gradients (flexoelectricity) @ domain walls results in polarization rotation

Catalan *et al.*, *Nature Mater.* **10**, 963 (2011)



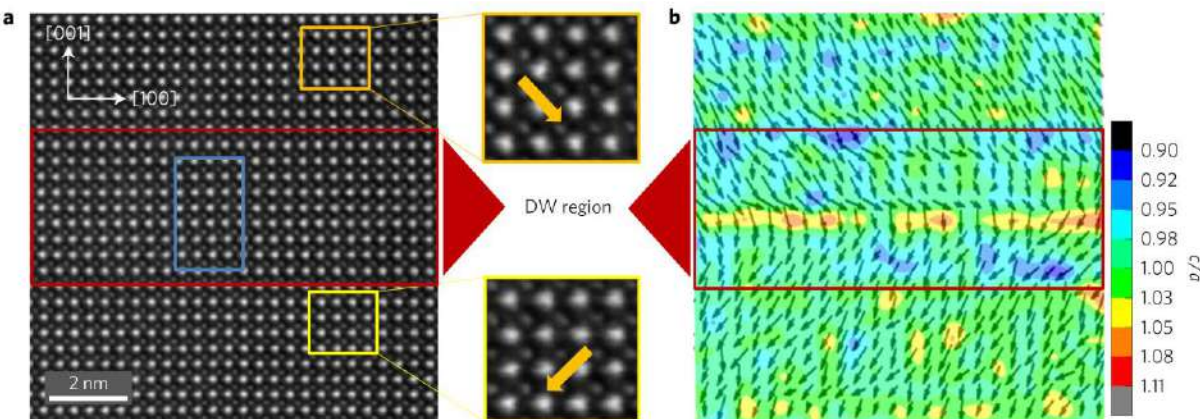
BiFeO<sub>3</sub> → Atomic resolution polarization mapping, unusual nanodomain structures

Nelson *et al.*, *Nano Lett.* **11**, 828 (2011)



BiFeO<sub>3</sub> → Direct mapping of defects/charge accumulation at domain walls that results in domain-wall conductivity

Rojac *et al.*, *Nature Mater.* **16**, 322 (2017)

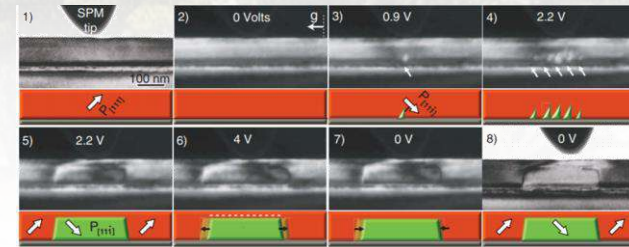




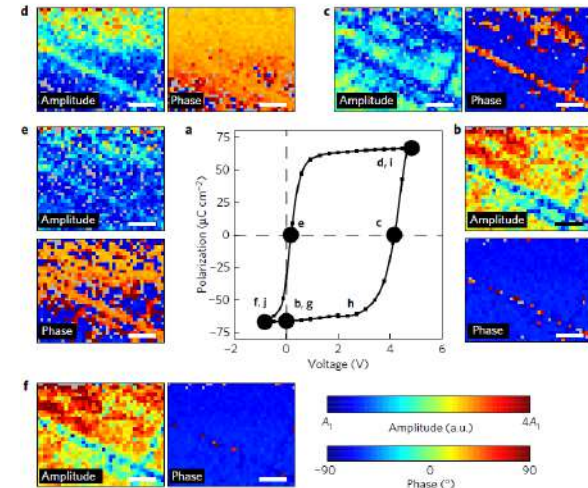


# In operando Studies of Domains

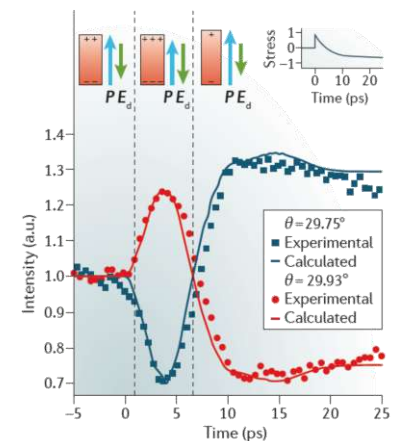
- **STEM** → Can probe ferroelectrics at the timescales of real processes important for device operation; increasing frame rates, data-acquisition rates during high-resolution imaging to probe ferroelectric switching  
 Nelson *et al.*, *Science* **334**, 968 (2011); Gao *et al.*, *Nature Commun.* **2**, 591 (2011); Winkler *et al.*, *Nano Lett.* **14**, 3617 (2014)



- **PFM** → Band excitation is changing the way FEs are studied under bias; real-time, *in operando* investigations of FE switching with nanoscale resolution and with the quantification and extraction of extensive data  
 Jesse, Kalinin *J. Phys. D* **44**, 464006 (2011); Agar *et al.*, *Nature Mater.* **15**, 549 (2016)



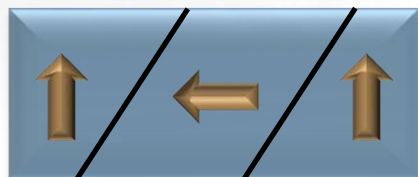
- **Synchrotron** → Undergoing rapid development; bright-light sources enable *in operando* studies of devices under applied fields, study of ultrathin films, the evolution of FEs under environmental conditions, and ultrafast probing of dynamic responses  
 Evans *et al.*, *Am. Ceram. Soc. Bull.* **92**, 18 (2013); Do *et al.*, *Nature Mater.* **3**, 365 (2004); Fong *et al.*, *Phys. Rev. Lett.* **96**, 127501 (2006); Grigoriev *et al.*, *Phys. Rev. Lett.* **96**, 187601 (2006); Highland *et al.*, *Phys. Rev. Lett.* **105**, 167601 (2010); Jo *et al.*, *Phys. Rev. Lett.* **107**, 055501 (2011); Highland *et al.*, *Phys. Rev. Lett.* **107**, 187602 (2011); Daranciang *et al.*, *Phys. Rev. Lett.* **108**, 087601 (2012); Wen *et al.*, *Phys. Rev. Lett.* **110**, 037601 (2013); Hruszkewycz *et al.*, *Phys. Rev. Lett.* **110**, 177601 (2013);



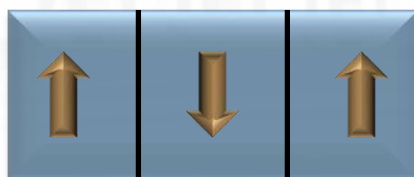


# Domain-Wall Contributions to Properties

## Domain Walls in Ferroelectrics



90° DW



180° DW

$$\chi_\alpha = \frac{d\langle P \rangle}{d\alpha} = \frac{d(\phi_c P_s)}{d\alpha}; \alpha = T, E, \sigma$$

$$= \phi_c \frac{dP_s}{d\alpha} + P_s \frac{d\phi_c}{d\alpha}$$

Intrinsic

Extrinsic

Domain wall displacement →  
Extrinsic contribution to  $\chi$

## Total Pyroelectric Coefficient ( $\pi$ )

$$\pi = \pi_i + \pi_e + \pi_s$$

Intrinsic  $\pi \rightarrow \Delta T$   
gives rise to  $\Delta P_s$   $\phi_c \left( \frac{dP_s}{dT} \right)$

Primary  
Pyroelectric  
Effect

Extrinsic  $\pi \rightarrow$  Temp. dependent  
movement of domain walls  $P_s \left( \frac{d\phi_c}{dT} \right)$

Secondary  $\pi \rightarrow$  Diff. in thermal  
expansion between film/sub.  
 $\propto d_{31} (\alpha_f - \alpha_s)$

Zook and Liu, J. Appl. Phys. 49, 4604 (1978)

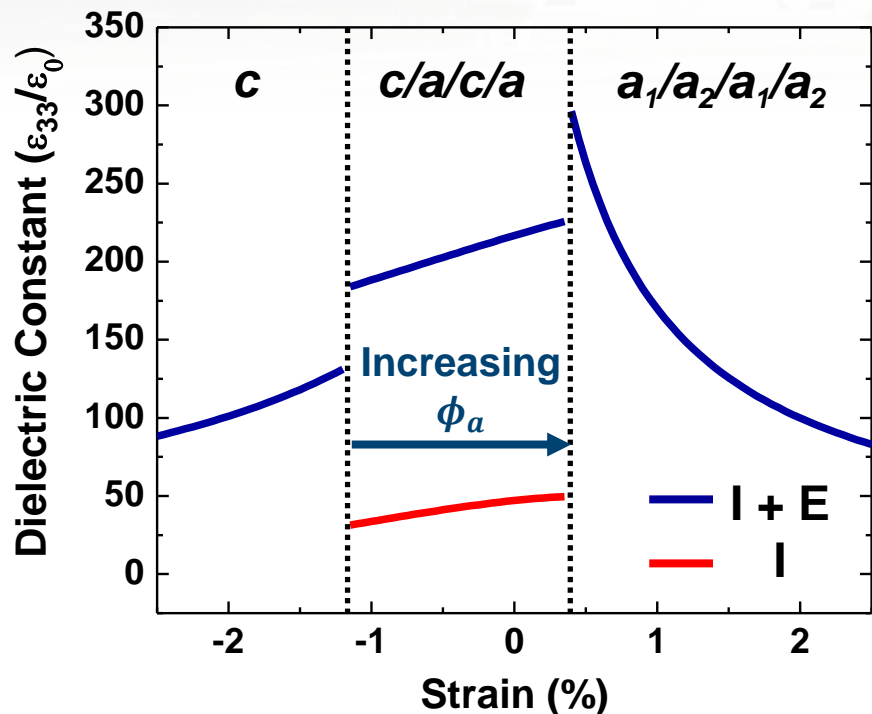
Very little work on extrinsic  
contributions to  $\pi$



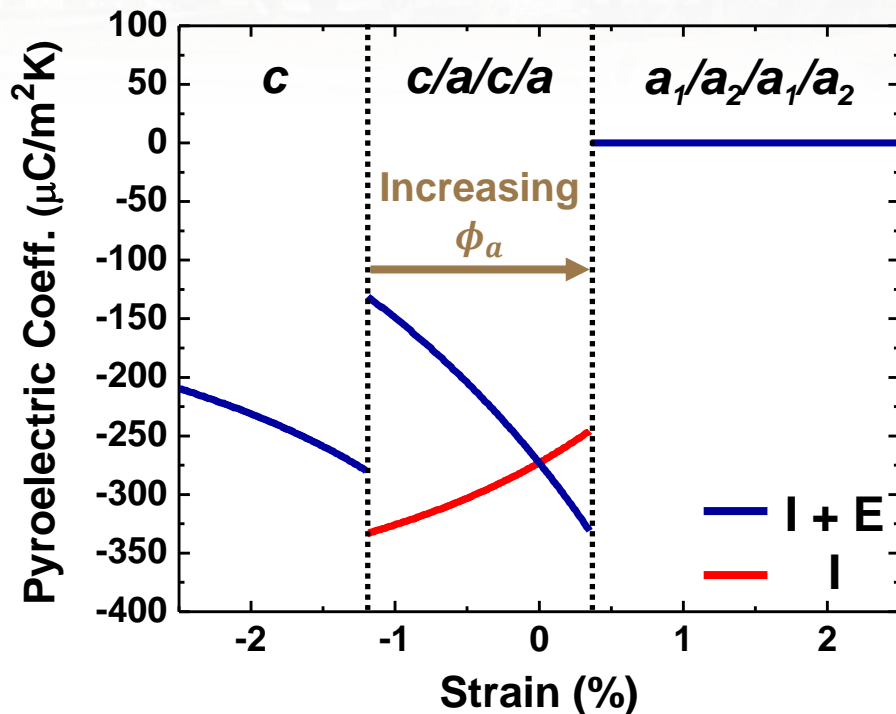


# DW Effects in $\text{PbZr}_{0.2}\text{Ti}_{0.8}\text{O}_3$ Thin Films

Dielectric permittivity  $\epsilon = \left(\frac{dP}{dE}\right)_T$



Pyroelectric coefficient  $\pi = \left(\frac{dP}{dT}\right)_E$



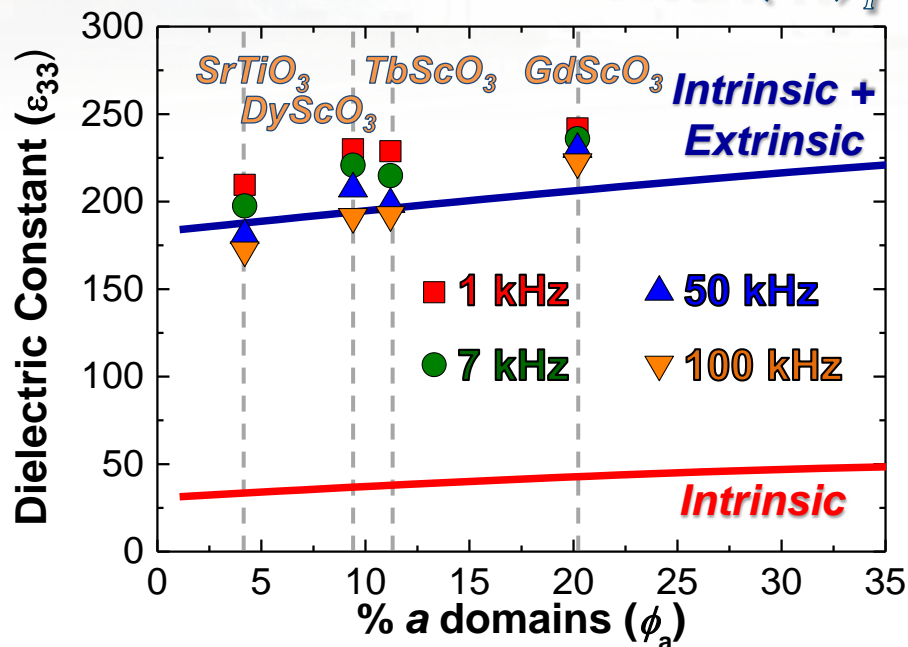
Models suggest that 90° domain walls...

- Increase dielectric and pyroelectric response
- Can dramatically change nature of response

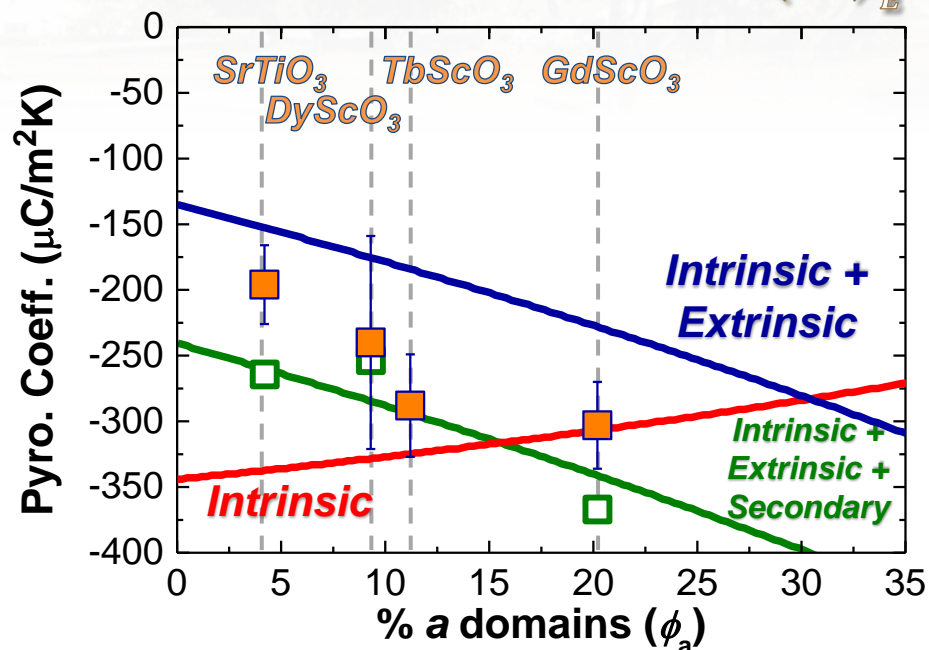


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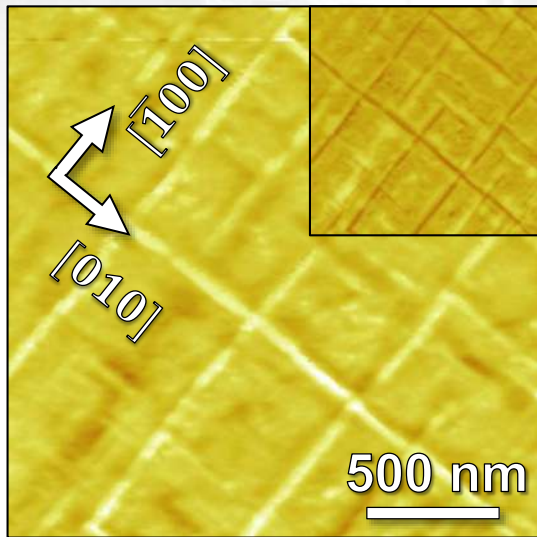


- Optimization of  $\chi_\alpha$  requires knowledge of domain structure
- $90^\circ$  domain walls  $\rightarrow$  impact the evolution of properties
- Subtle (but important) differences btw  $E$  &  $T$  susceptibilities

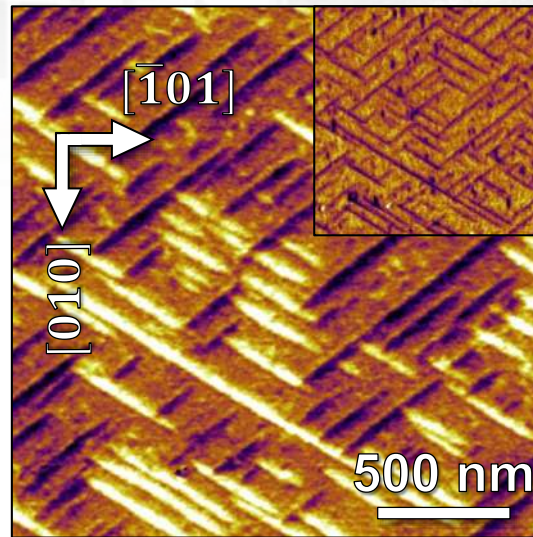


# Orientation & Domains: $\text{PbZr}_{0.2}\text{Ti}_{0.8}\text{O}_3$

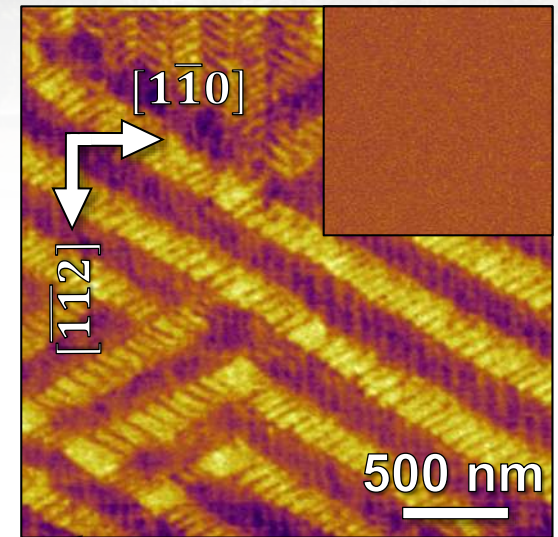
(001)



(101)



(111)



- Orientation provides a knob by which we can control the domain structure of materials
- All samples possess  $90^\circ$  domain walls  $\rightarrow$  controlled structures
- Advantage of thin films  $\rightarrow$  direct observation and quantification of domain structure features

**Volume fraction of minority domains**

**Line density of domain walls**

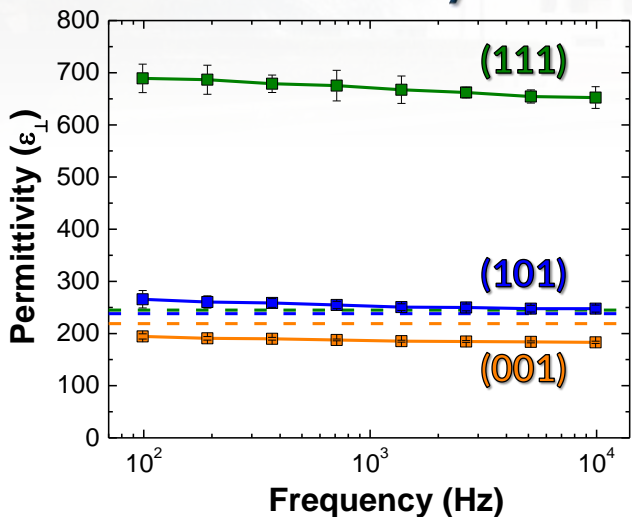
Orientation	$\lambda$ ( $\mu\text{m}^{-1}$ )	$\phi$ (%)
(001)	8.91	15.3
(101)	16.3	19.9
(111)	48.9	33.3





# Exotic Low-/High-Field Effects

## Stationary (Frozen) Permittivity



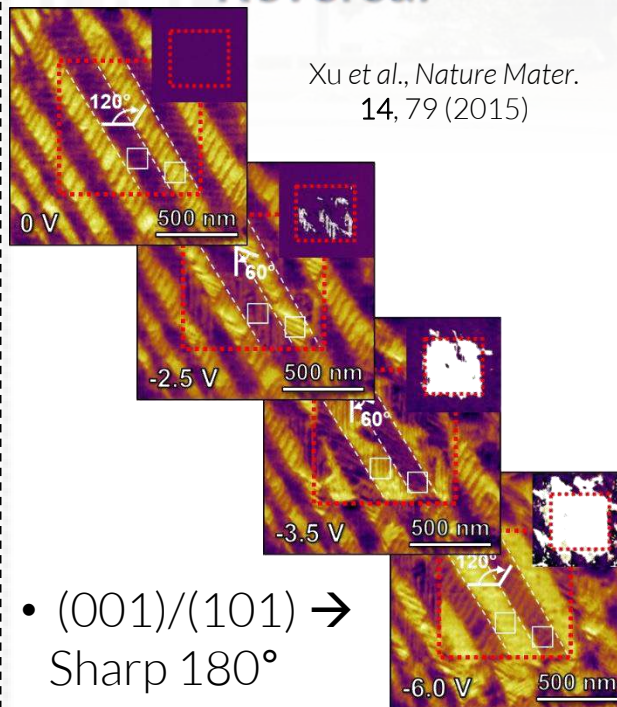
Xu et al., *Nature Commun.* 5, 3120 (2014)

- Response of the volume of the ferroelectric material within the finite width of the domain walls (non-motional)

$$\epsilon_{dw} \approx 1,500-19,000 \text{ (for 1-10 nm)}$$

$$\rightarrow 6-78 \times \epsilon_b$$

## Multi-Step Polarization Reversal

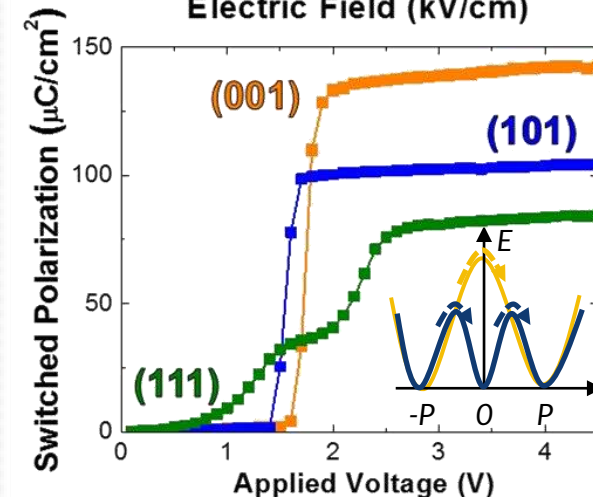
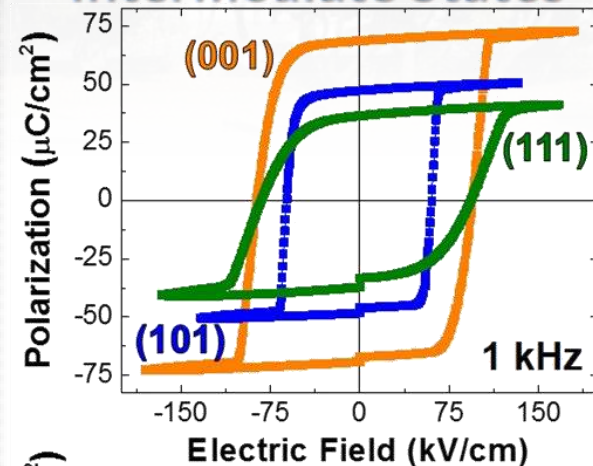


Xu et al., *Nature Mater.* 14, 79 (2015)

- (001)/(101)  $\rightarrow$  Sharp 180° switching process
- (111)  $\rightarrow$  Broad switching, 90° switching events that match models

Observation of multi-step 90° switching process  $\rightarrow$  intermediate states?

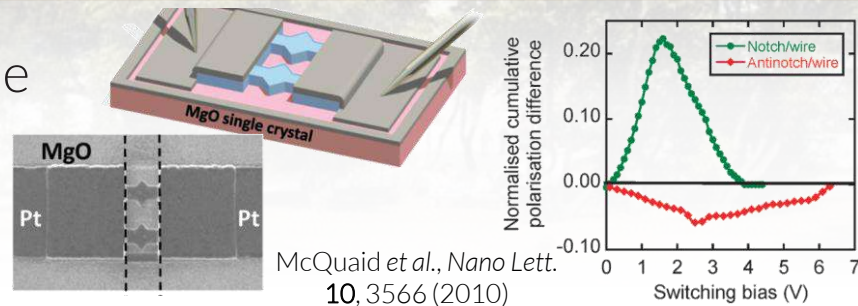
## Indications of Intermediate States



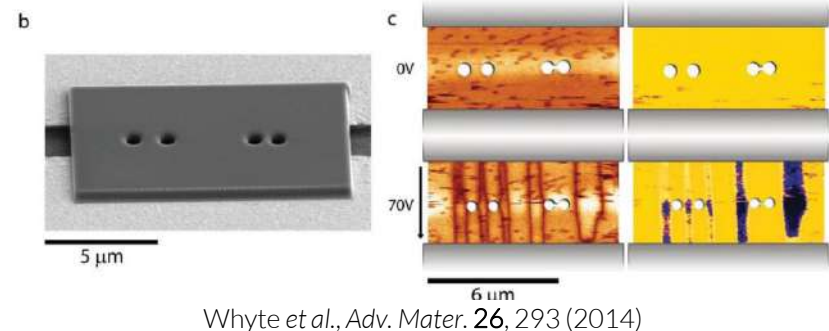
Capacitor-based studies suggest potential for intermediate states...

# Controlling DW Formation/Motion

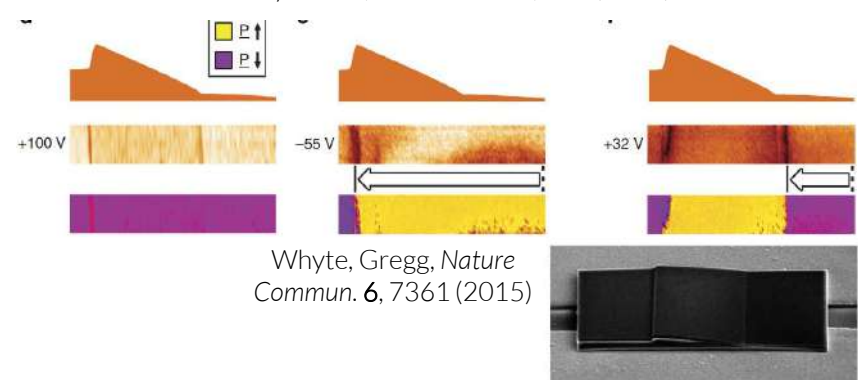
- Micro- and nano-machining → Controlling the boundary conditions controls how domains respond, tune mobility with “notches” and “antinotches”



- Making something from nothing → Researchers used “air holes” w/ FIB to create hot-spots to manipulate domain nucleation/growth



- Ferroelectric domain-wall diode → Single direction of motion for domain walls, irrespective of polarity, under a series of alternating electric field pulses → Sawtooth morphology is central to its function

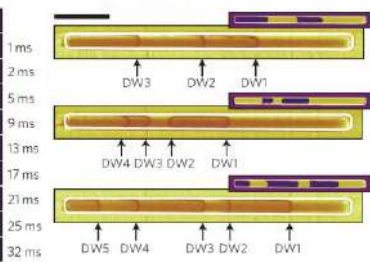
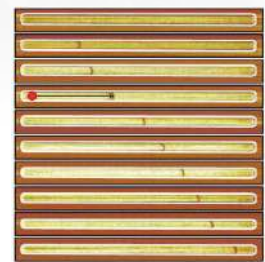


**Understanding the energies at play enables one to grab deterministic control over domain production and motion under applied fields**



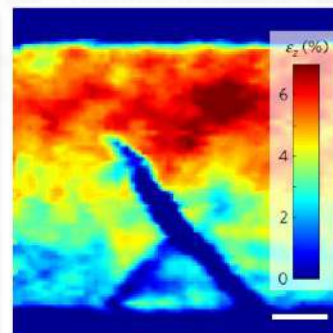
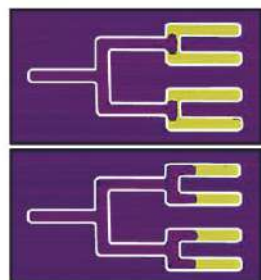
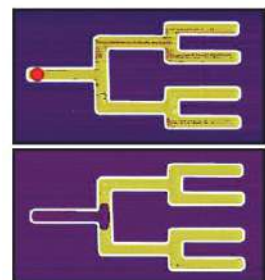
# Controlling DW Formation/Motion

To fully harness potential of DWs functional entities → Essential to achieve reliable, precise control of their nucleation, location, number, and velocity

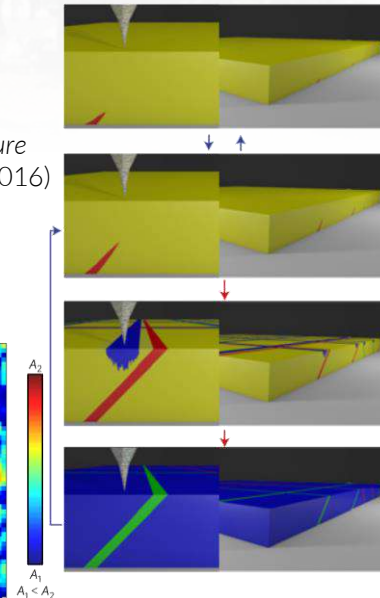
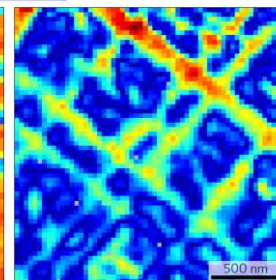
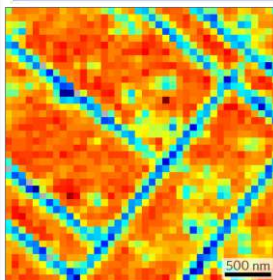


McGilly et al., *Nature Nanotechnol.* **10**, 145 (2015)

**Emulates magnetic logic devices...**



Agar et al., *Nature Mater.* **15**, 549 (2016)



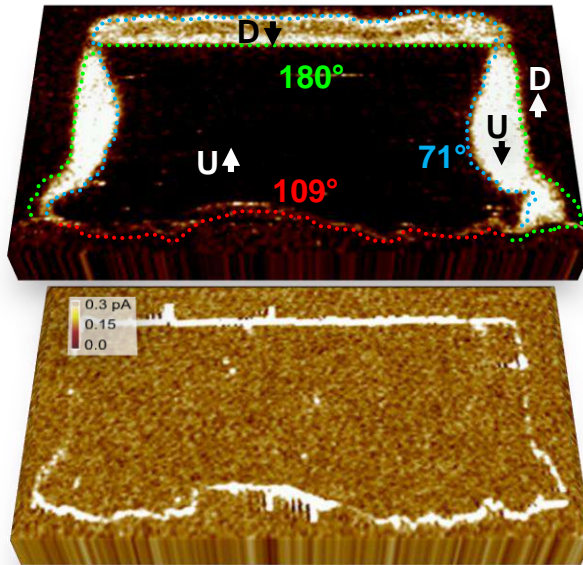
- Ability to “pick” a nucleation site
- Position controlled by tuning of voltage pulses; multiple domain walls can be nucleated and handled reproducibly
- A step towards the realization of domain-wall nanoelectronics utilizing ferroelectric thin films

- Strain-gradients → manipulate domain structure, “needle” domains
- Such domains are found to be highly mobile → out-of-plane fields drive reduction/growth of the  $a$  domain
- Produces large piezoelectric responses at domain walls → not seen in classic

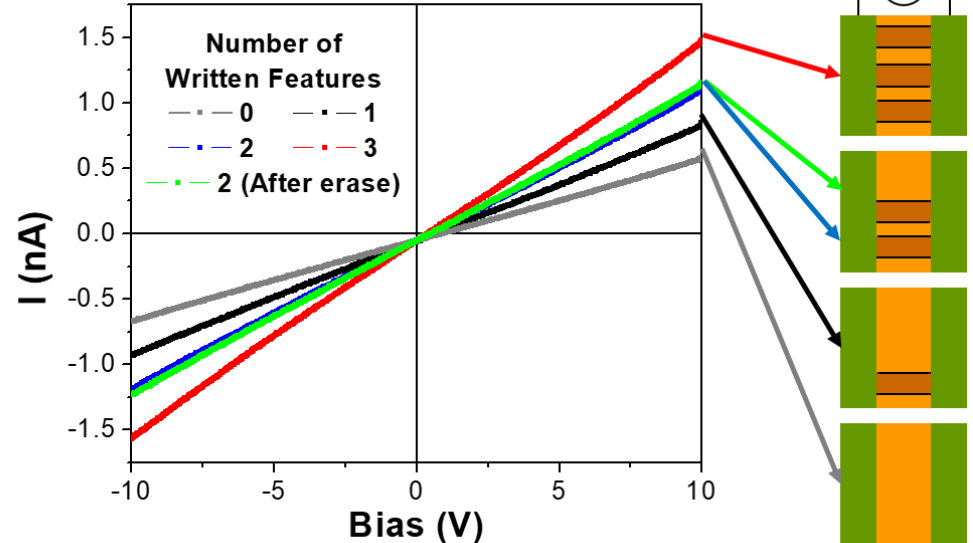


# Domain Wall as the Device

- Domain walls can contribute to properties via *extrinsic* contributions...
- But what about making the domain wall itself the functional unit?
- Ferroelectric domain walls as device elements Catalan et al., Rev. Mod. Phys. **84**, 119 (2012); Vasudevan et al., Adv. Funct. Mater. **23**, 2592 (2013)
- Domain walls can exhibit properties that differ from those of the bulk
  - Domain-wall conductivity in insulating ferroelectrics → BiFeO<sub>3</sub>



Seidel, Martin et al., Nature Mater. **8**, 229 (2009)



- Resistance of a single wall  $\sim 8 \times 10^{10} \Omega$  → Estimated resistivity  $\sim 2.5 \Omega \cdot m$ 
  - Bulk resistivity  $\sim 10^7 \Omega \cdot m$
- Conductivity was ascribed to an increased carrier density resulting from an electrostatic potential step and a decrease in the bandgap within the wall

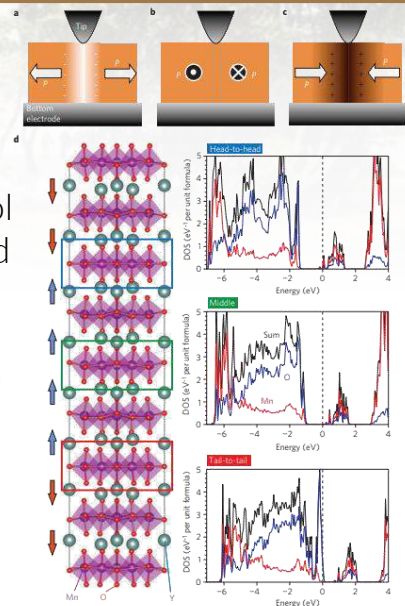
Eliseev et al., Phys. Rev. B **83**, 235313 (2012); Eliseev et al., Phys. Rev. B **85**, 045312 (2012); Morozovska et al., Phys. Rev. B **86**, 085315 (2012)

# Function at Domain Walls

- Conduction at domain walls in...
  - BiFeO<sub>3</sub> (all types) Farokhipoor *et al.*, *Phys. Rev. Lett.* **107**, 127601 (2011)
  - PbZr<sub>1-x</sub>Ti<sub>x</sub>O<sub>3</sub> Guyonnet *et al.*, *Adv. Mater.* **23**, 5377 (2011)
  - BaTiO<sub>3</sub> Sluka *et al.*, *Nature Commun.* **4**, 1808 (2013)
  - Er(Y)MnO<sub>3</sub> Meier *et al.*, *Nature Mater.* **11**, 284 (2012)
  - HoMnO<sub>3</sub> Wu *et al.*, *Phys. Rev. Lett.* **108**, 077203 (2012)
  - LiNbO<sub>3</sub> Schroder *et al.*, *Adv. Funct. Mater.* **22**, 3936 (2012)



Probe-based control of head-to-head and tail-to-tail domain wall configuration and band structure in YMnO<sub>3</sub>



Seidel *et al.*, *Phys. Rev. Lett.* **105**, 197603 (2010);  
Gaponenko *et al.*, *Appl. Phys. Lett.* **106**, 162902 (2015)

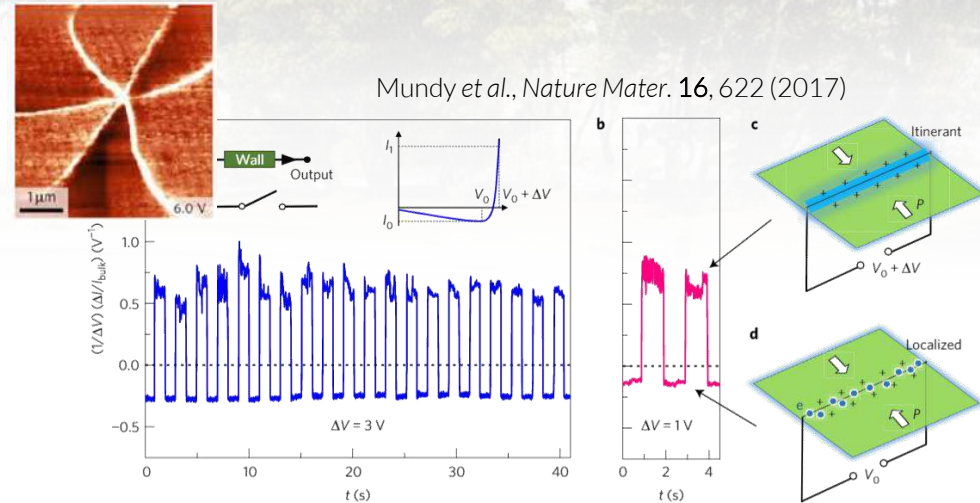
- Studies have probed...
  - Role of oxygen vacancies in the conduction behavior Seidel *et al.*, *Phys. Rev. Lett.* **105**, 197603 (2010); Gaponenko *et al.*, *Appl. Phys. Lett.* **106**, 162902 (2015)
  - How domain-wall curvature can give rise to large conductivity changes (~500%) Vasudevan *et al.*, *Nano Lett.* **12**, 5524 (2012)
  - How certain conducting domain walls in multiferroics can exhibit large magnetoresistance (~60%) He *et al.*, *Phys. Rev. Lett.* **108**, 067203 (2012)
  - Topological defect states and exotic properties in vortex domain structures Balke *et al.*, *Nature Phys.* **8**, 81 (2012)
  - Memristor-like functionality Maksymovych *et al.*, *Nano Lett.* **11**, 1906 (2011)

**Advances in synthesis/characterization have exposed new phenomena...  
Was the ultimate nanoscale device there all along?**

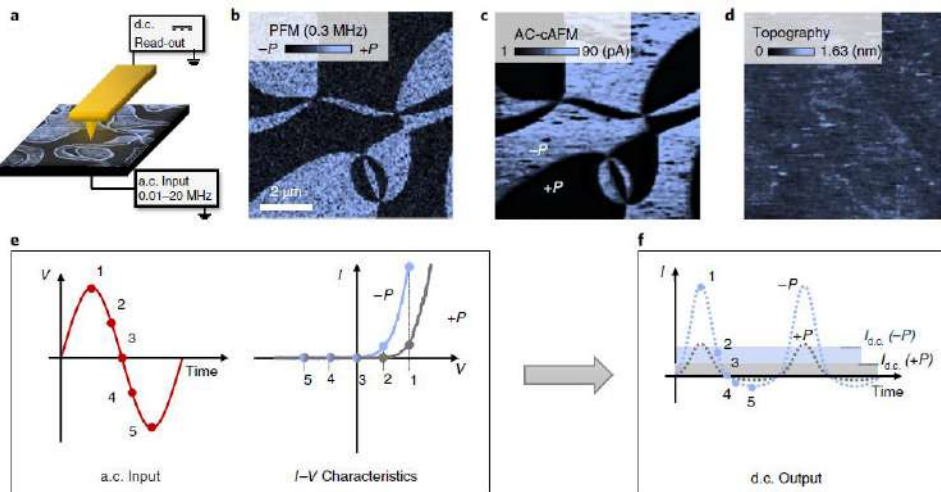


# Domain-Wall Devices

- *E*-field control of the electronic transport at FE domain walls
- Reversibly switch resistive/conductive behavior at charged walls in  $\text{ErMnO}_3$
- Transition  $\rightarrow$  formation (activation) of inversion layer that acts as the channel for the charge transport
- “Foreshadow the possibility to design elementary digital devices for all-domain-wall circuitry”



- *Diode-like alternating-to-direct current conversion* based on neutral FE domain walls in  $\text{ErMnO}_3$
- Showed rectification at the tip-wall contact for frequencies at which the walls are effectively pinned
- Mechanism  $\rightarrow$  Transport behavior at the walls arises from oxygen defects
- Frequency regime/magnitude of the direct current output controlled by bulk conductivity  $\rightarrow$  electrode-wall junctions



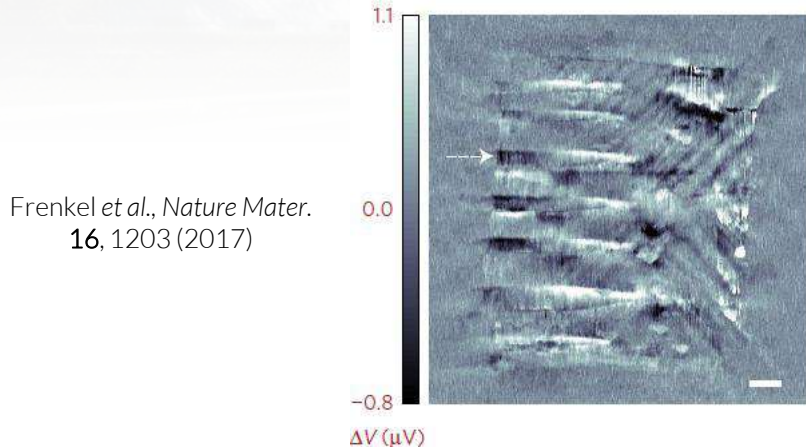
Schaab et al., *Nature Nanotechnol.* **13**, 1028 (2018)



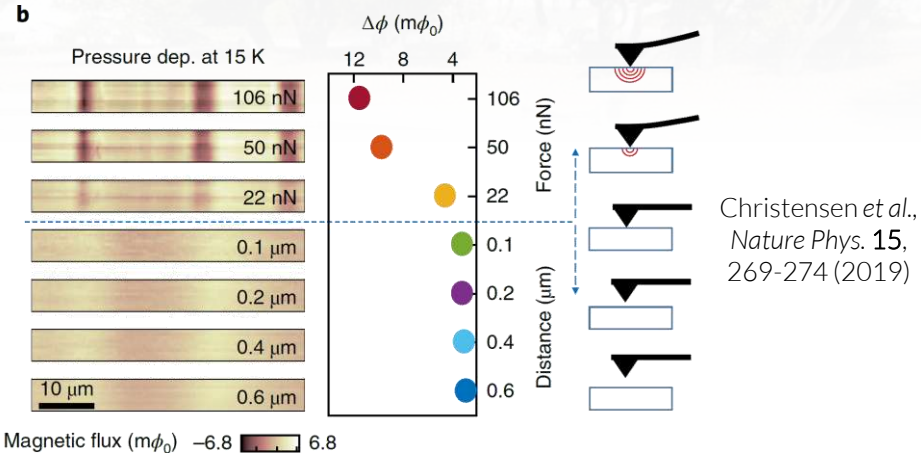
# Domains in Ferroelastic SrTiO<sub>3</sub>

## Domain-wall Nanoelectronics Using Ferroelastic Materials

### Local strain tuning of conductivity



### Local strain tuning of magnetism



- Twin boundaries, with properties that are different from their surrounding bulk, tune LaAlO<sub>3</sub>/SrTiO<sub>3</sub> interface 2DEG
- SrTiO<sub>3</sub> ferroelastic domain boundaries → Remaining highly mobile @ low-T, recently suggested to be polar
- Localized pressure to a twin boundary, detect a change in LaAlO<sub>3</sub>/SrTiO<sub>3</sub> interface
- Polarity at the twin boundaries, control of conducting state

- Ferroelastic domain walls display strain-tunable polarity and enhanced conductivity
- Long-range magnetic order with modulations along the ferroelastic domain walls in SrTiO<sub>3</sub>/heterointerfaces, manifests as a striped pattern in scanning SQUID
- Conducting interfaces → clear signatures in magnetotransport measurements
- Magnetic state coupled dynamically to the lattice, reversibly tuned by local forces

## New Multiferroic (Ferroelastic + Ferroelectric + Ferromagnetic)



# Where do we go from here?

High-throughput discovery and design of next-generation functional materials → Can we design new domain-wall function?

Non-traditional ferroelectrics → Domain walls in polymer, 'stuffed wurtzite' structures, 2D materials, 'hybrid perovskite' crystals, thiophosphates,  $\text{HfO}_2$ ,...

Real-time studies → Simulation/experiment can do exciting things; data analysis/big-data/machine learning will extract new understanding of domain walls/structures

Ferroelectrics for energy → Domain walls can play a roll in impacting photovoltaic response, catalysis, physisorption rates of reactants,...

True domain-wall devices → emulate real device components with domain walls? Can we deterministically place/control them?

**Domains & domain walls result from energy minimization in ferroic materials...**

**Their configuration can have profound impact on the macroscale properties...**

**Domain walls themselves can contribute properties unique from the bulk...**