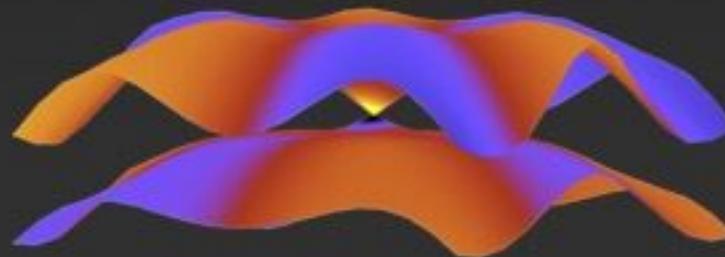


AB-INITIO SPIN-ORBITRONICS

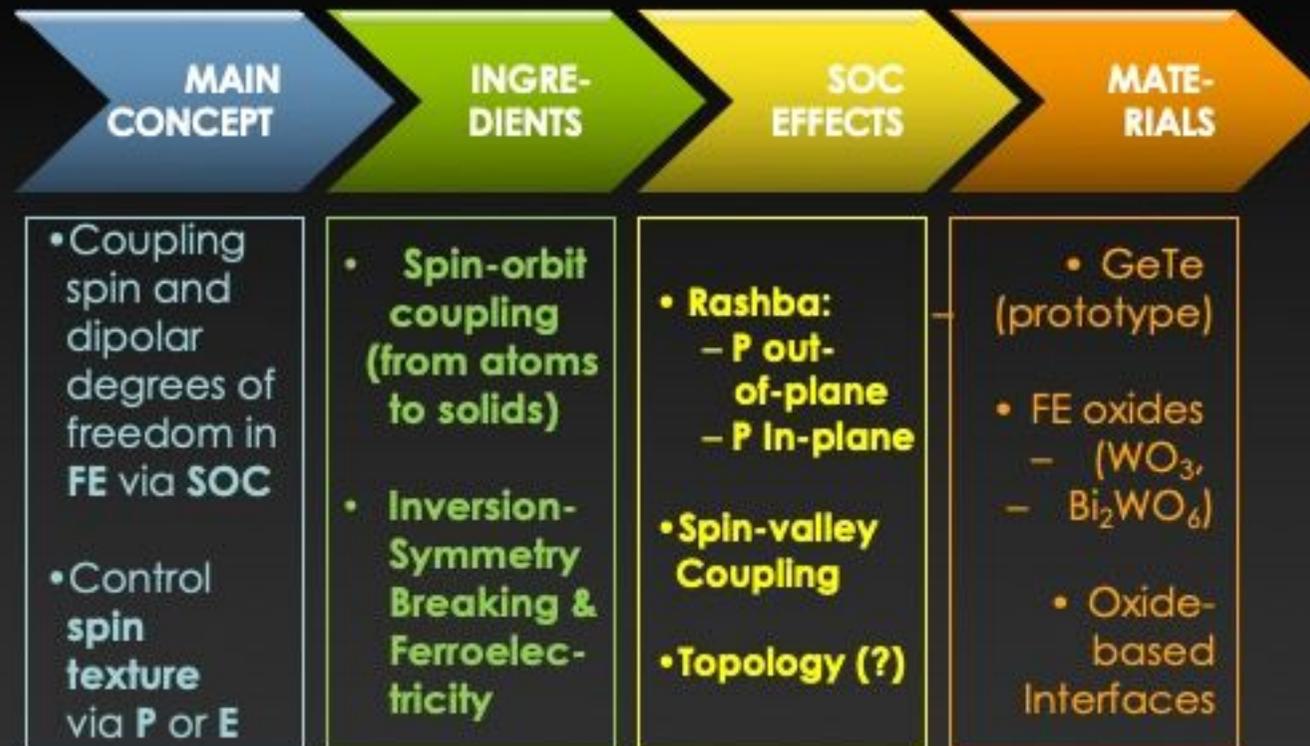


Dr. Silvia Picozzi

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Materiali e Dispositivi Innovativi
CNR-SPIN - 66100 Chieti, Italy*
<https://sites.google.com/site/silviapicozzi/>



OUTLINE

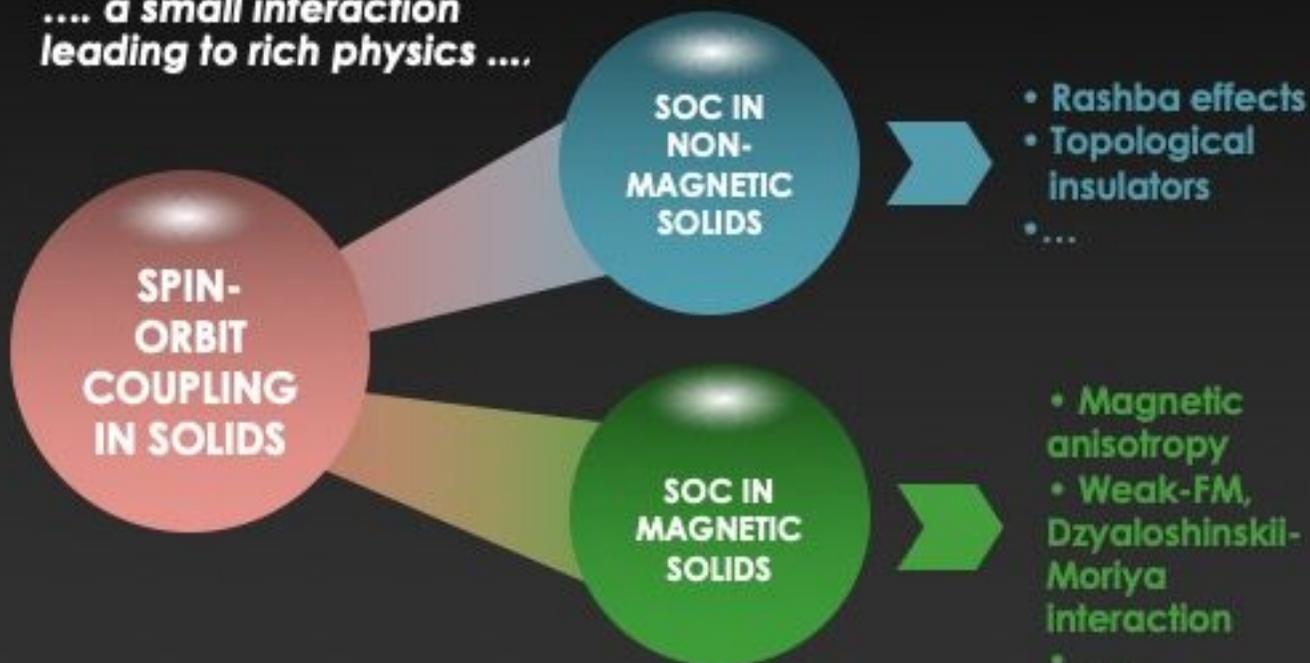


SPIN-ORBIT COUPLING

Relativistic interaction linking spin and orbitals (i.e. spin space and real space)

$$\underbrace{\frac{\mu_B}{2mcr}}_{\xi} \frac{dV(r)}{dr} \vec{\sigma} \cdot (\vec{r} \times \vec{p}) = \xi \vec{\sigma} \cdot \vec{L}$$

.... a small interaction leading to rich physics



SPIN-ORBIT INTERACTION: A RELATIVISTIC EFFECT

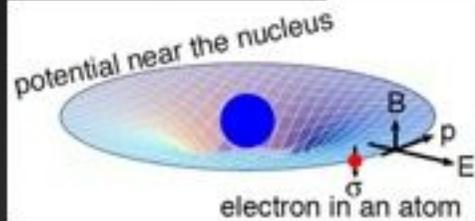
Relativity:

Switch the coordinate frame

Moving $E \rightarrow B = \frac{1}{c^2} E \times v.$

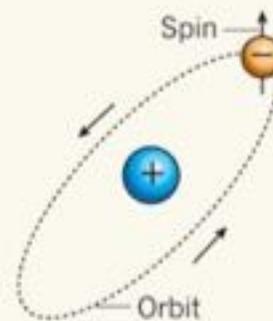
Quantum spin:

Electron has a spin $s \rightarrow h = -\mu_s \cdot B$

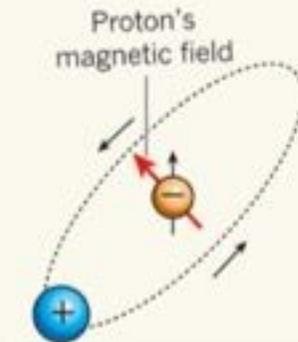


a Electron in an atom

Proton's point of view



Electron's point of view



$$B = \frac{1}{c^2} E \times v. \sim Z^4$$

$$E = -\frac{\mathbf{r}}{r} \phi'$$

$$B = -\frac{1}{rc^2} \phi' \mathbf{r} \times \mathbf{v}$$

As $\mathbf{l} = \mathbf{r} \times \mathbf{p}$, and hence $\mathbf{r} \times \mathbf{v} = \frac{1}{m_e} \mathbf{l}$, we get the Hamiltonian

$$b_{SO} = -\mu_e \cdot B = \frac{1}{m_e r c^2} \phi' \mu \cdot \mathbf{l} = \frac{e}{m_e^2 r c^2} \phi' \mathbf{s} \cdot \mathbf{l}.$$



RELATIVISTIC EFFECTS

$$H_0 = \frac{p^2}{2m} + V(r) = \frac{p^2}{2m} - \frac{Ze^2}{4\pi\epsilon_0} \frac{1}{r}$$

Unperturbed Hamiltonian

$$H = H_0 + H_1 + H_2 + H_3$$

Three Relativistic corrections:

$$H_1 = -\frac{p^4}{8m^3c^2}$$

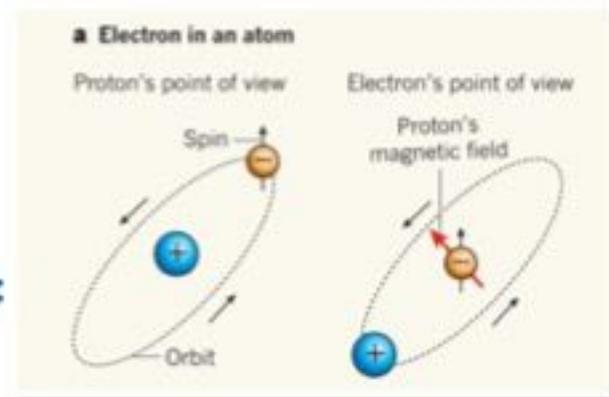
H_1 : Mass-velocity term (correction to Kinetic energy)

$$H_2 = \frac{1}{2m^2c^2} \frac{1}{r} \frac{dV}{dr} \mathbf{L} \cdot \mathbf{S}$$

H_2 : Spin-orbit term

$$H_3 = \frac{\hbar^2}{8m^2c^2} \nabla^2 V(r) = \frac{\pi\hbar^2}{2m^2c^2} \left(\frac{Ze^2}{4\pi\epsilon_0} \right) \delta(r)$$

H_3 : Darwin term



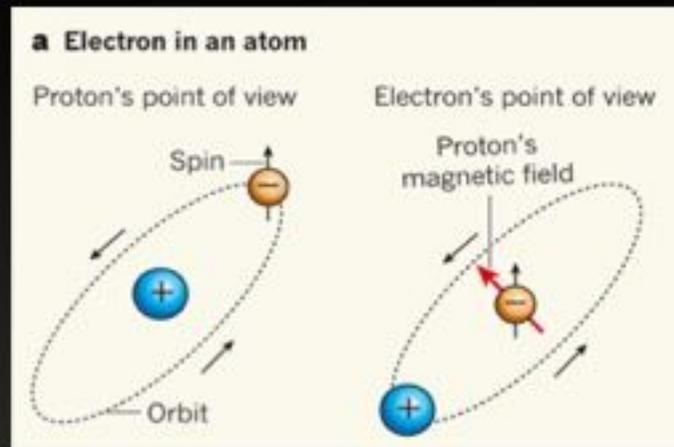
MAGNITUDE OF RELATIVISTIC CORRECTIONS

Tab. 1.1 Summary of relativistic and spin-dependent interaction terms.

Perturbation	Description	Magnitude
$\frac{\mathbf{p}^2}{2m} - \frac{Ze^2}{4\pi\epsilon_0} \frac{1}{r}$	The nonrelativistic motion,	$\gtrsim 10^5 \text{ cm}^{-1}$
$H_1 = -\frac{\mathbf{p}^4}{8m^3c^2}$	Relativistic mass correction term.	$\approx 0.1 \text{ cm}^{-1}$
$H_3 = \frac{\hbar^2}{8m^2c^2} \nabla^2 V(r)$	The "Darwin" term responsible for <i>s</i> -state shifts. It represents the relativistic nonlocalizability of the electron and is related to both the negative-energy sea and its rapid motion.	$< 0.1 \text{ cm}^{-1}$
$H_2 = \frac{1}{2m^2c^2} \frac{1}{r} \frac{dV}{dr} \mathbf{L} \cdot \mathbf{S}$	Spin-orbit interaction. As shown in Problem 1.4, this term can be written as $\frac{\hbar^2}{2m^2c^2} \frac{1}{r} \frac{d\phi}{dr} (\vec{l} \cdot \vec{s})$, where \vec{l} is the orbital angular momentum. In contrast to the previous term, this does not affect <i>s</i> states.	$10\text{--}10^3 \text{ cm}^{-1}$



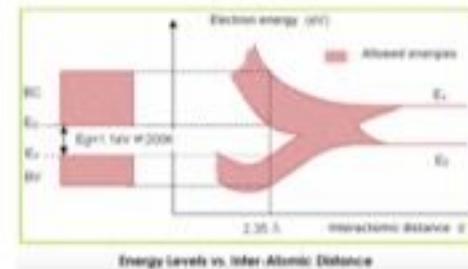
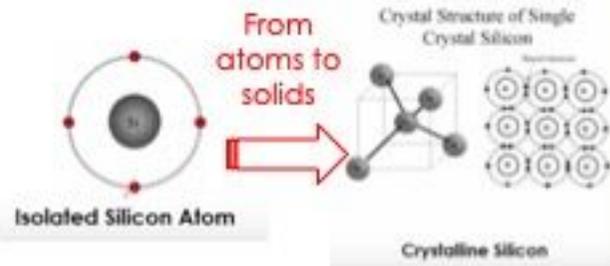
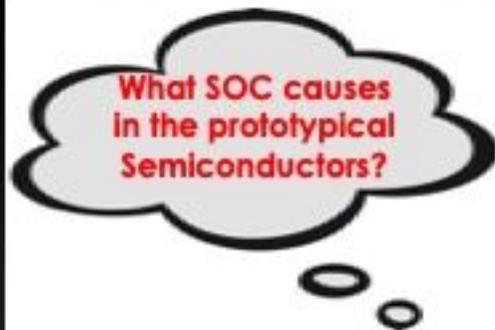
SPIN-ORBIT INTERACTION: FROM ATOMS TO SOLIDS



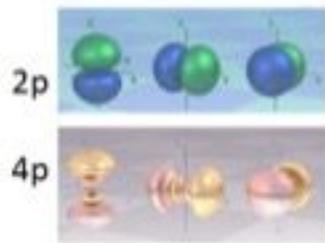
In an atom, an e^- (orange) orbits the nucleus (blue; here a single proton). From the e^- point of view, the proton orbits the e^- and produces a magnetic field that couples with the e^- spin and alters its orbit.



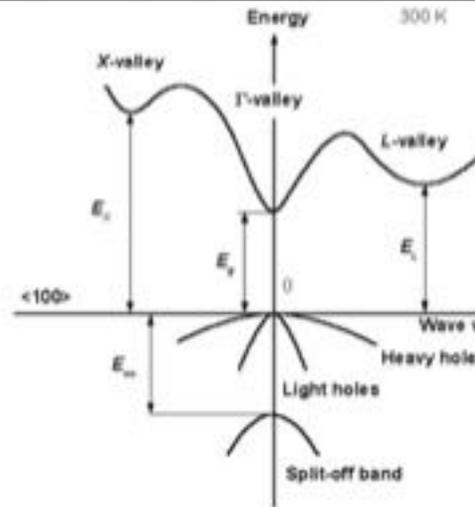
FROM ATOMS TO SOLIDS: SYMMETRY-INDEPENDENT SOC



SYMMETRY-INDEPENDENT SOC



$$\Delta \sim Z^2 \text{ (outer p shell)}$$



Compound	Δ_0^{sp} (eV)
C	0.006
Si	0.044
Ge	0.29
α -Sn	
AlN	
AlP	
AlAs	
AlSb	0.75
GaN	0.011
GaP	0.127
GaAs	0.34
GaSb	0.80
InN	
InP	0.11
InAs	0.38
InSb	0.82
ZnO	-0.005
ZnS	0.07
ZnSe	0.43
ZnTe	0.93
CdS	0.066
CdSe	
CdTe	0.92
HgS	
HgSe	
HgTe	1.08

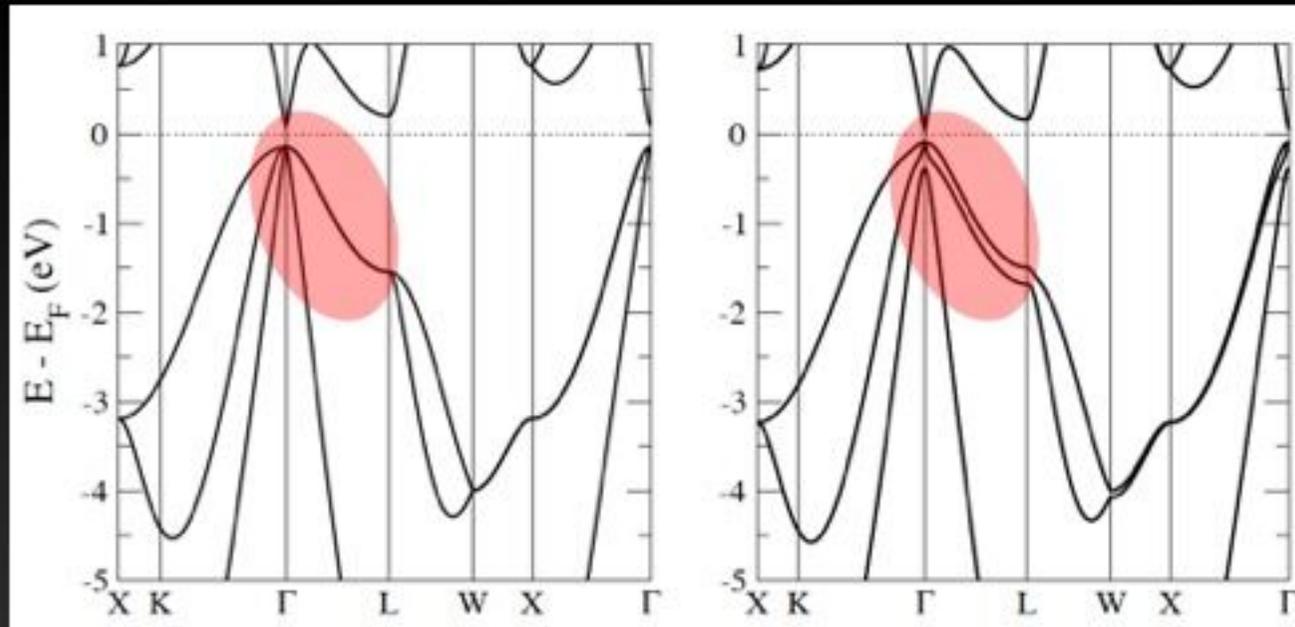
$$\left(\frac{Z_{Ge}}{Z_{Si}}\right)^2 = \left(\frac{(Z_{Ga} + Z_{As})/2}{Z_{Si}}\right)^2 = \left(\frac{32}{14}\right)^2 = 5.2$$

$$\frac{\Delta_{Ge}}{\Delta_{Si}} = 6.6 \quad \frac{\Delta_{GaAs}}{\Delta_{Si}} = 7.7$$

SOC splits bands.



SOC-DERIVED BAND SPLITTING



Bandstructure of Ge around the Fermi level without SOC (left) and with SOC (right). Three-fold degeneracy of the highest occupied state at the Γ point is split by SOC, as well as the doubly degenerate band along the lines ΓL and ΓX .



SYMMETRY PROPERTIES IN BAND STRUCTURE : SPACE/TIME INVERSION

What happens to
spin-orbit coupling
upon time-reversal
symmetry?

$$\begin{aligned}t &\rightarrow -t \\l &\rightarrow -l \\ \sigma &\rightarrow -\sigma\end{aligned}$$



$$V_{so} \sim l \cdot \sigma \rightarrow V_{so}$$

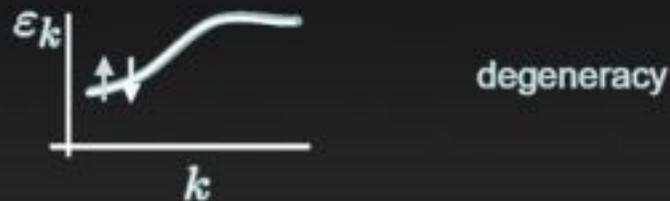
SOC preserves time reversal symmetry!



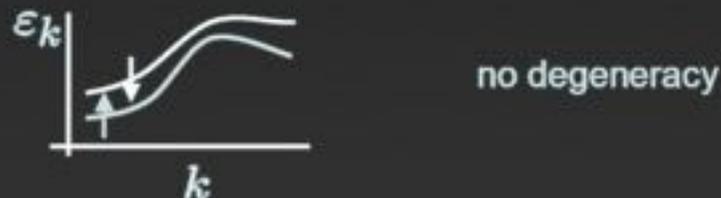
SYMMETRY PROPERTIES IN BAND STRUCTURE : SPACE/TIME INVERSION

Two cases need to be distinguished:

- solids with space inversion symmetry



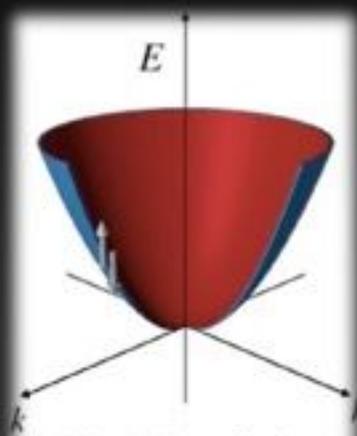
- solids without space inversion symmetry



SYMMETRY PROPERTIES IN BAND STRUCTURE : SPACE/TIME INVERSION

Time reversal symmetry: $\epsilon_{k\uparrow} = \epsilon_{-k\downarrow}$

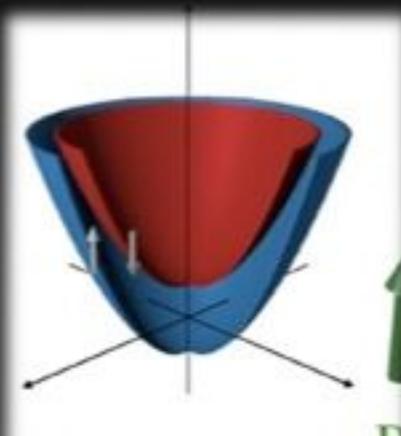
(Space) inversion symmetry: $\epsilon_{k\uparrow} = \epsilon_{-k\uparrow}$



a) No Spin-splitting

WITH time-reversal
and **WITH** Inversion
symmetry

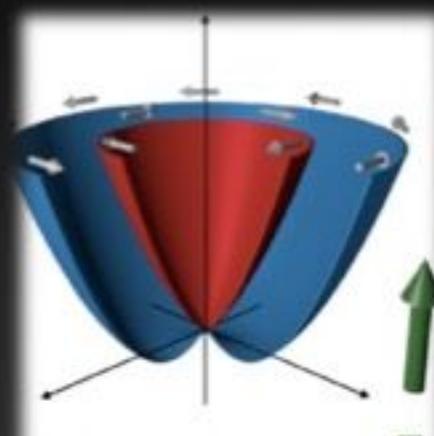
$$\epsilon_{k\uparrow} = \epsilon_{-k\uparrow} = \epsilon_{-k\downarrow}$$



b) Zeeman Spin-splitting **B**

WITHOUT time-
reversal symmetry
WITH Inversion

$$\epsilon_{k\uparrow} = \epsilon_{-k\uparrow} \neq \epsilon_{k\downarrow}$$



c) Rashba Spin-splitting **E**

WITH time-reversal
WITHOUT Inversion
symmetry:

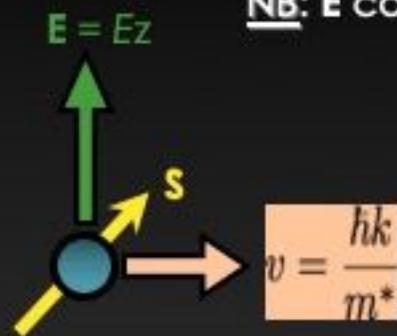
$$\epsilon_{k\uparrow} = \epsilon_{-k\downarrow} \neq \epsilon_{k\downarrow}$$



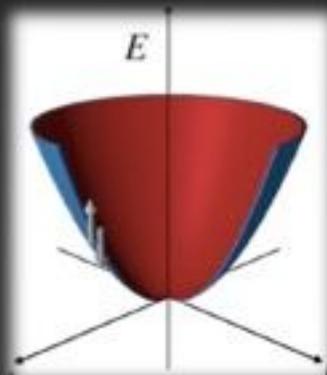
WHAT IS THE RASHBA EFFECT ?

SOC effect: particle in **electric field E** experiences an internal **effective magnetic field** $\mathbf{B}_{\text{eff}} \propto \mathbf{v} \times \mathbf{E}$ in its moving frame

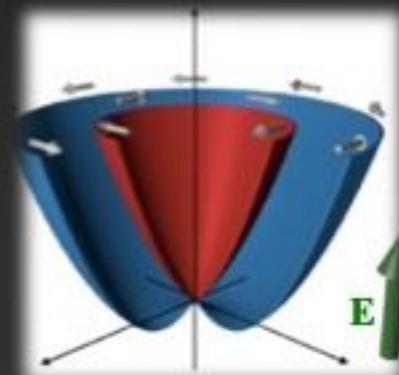
NB: **E** could be external (i.e. 2DEG), or “effective” (as in FE)



Without Rashba

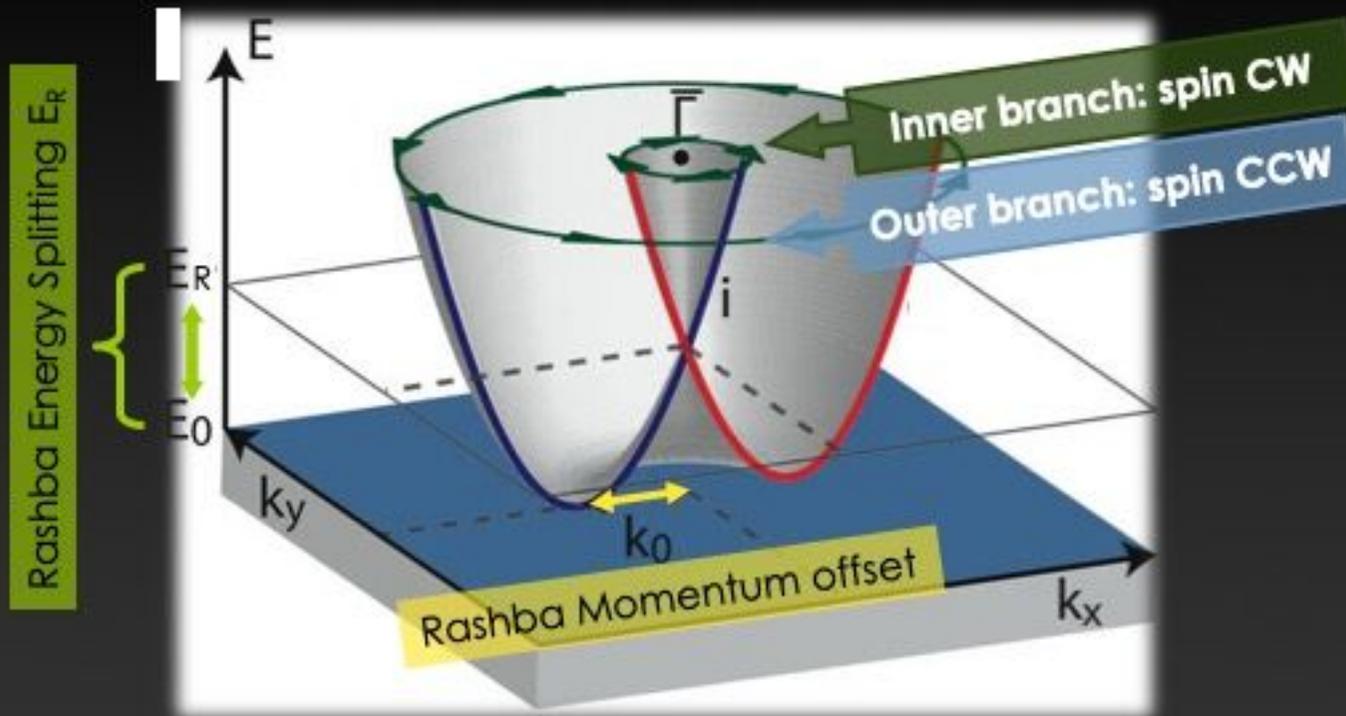


With Rashba



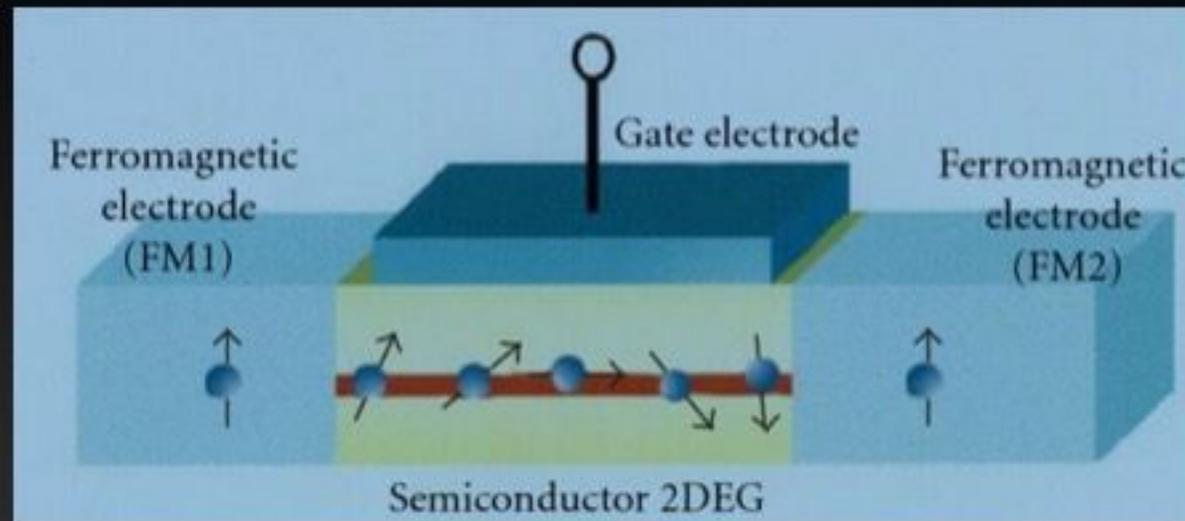
WHAT IS THE RASHBA EFFECT ?

Relevant properties/parameters:



Rashba coupling parameter (in eV Å): $\alpha_R = 2E_R/k_0$

USING RASHBA SPIN-ORBIT COUPLING: SEMICONDUCTOR-BASED SPIN-FET



- Das-Datta proposal (*Appl. Phys. Lett.* 1990 , 56 , 665)
- Animation by Bernevig and Sinova.



WHAT IS THE SPIN TEXTURE ?

Calculate the SPIN expectation value for each wavefunction

$$\langle \sigma_x \rangle_R = \langle \psi_n^*(k) | \sigma_x | \psi_n(k) \rangle / \langle \psi_n^*(k) | \psi_n(k) \rangle$$

$$\langle \sigma_y \rangle_R = \langle \psi_n^*(k) | \sigma_y | \psi_n(k) \rangle / \langle \psi_n^*(k) | \psi_n(k) \rangle$$

$$\langle \sigma_z \rangle_R = \langle \psi_n^*(k) | \sigma_z | \psi_n(k) \rangle / \langle \psi_n^*(k) | \psi_n(k) \rangle$$

Pauli spin matrices

$$\sigma_x = \frac{1}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$\sigma_y = \frac{1}{2} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

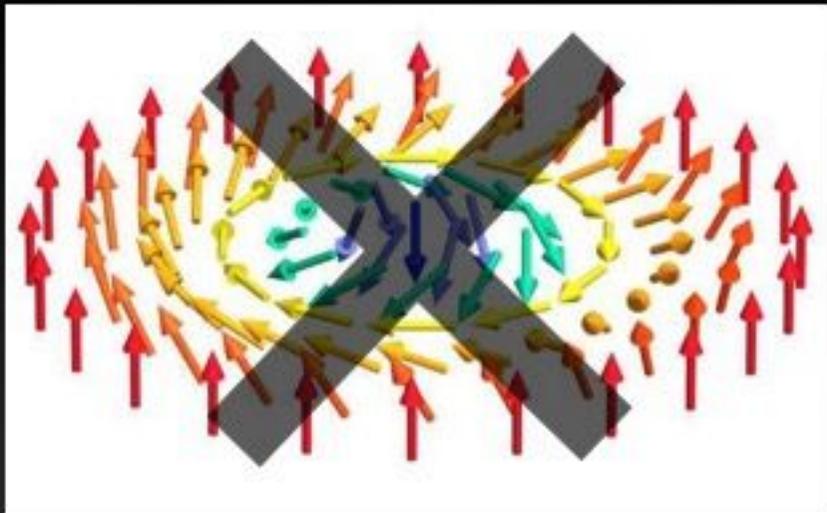
$$\sigma_z = \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

- Can one measure it?
Yes, with **Spin-resolved ARPES**
- What does it mean?
Intuitively: "spin" of an electron in a certain (n,k) state

Spin texture of Bi_2Se_3 Dirac cone



WHAT IS THE SPIN TEXTURE ?



- NB: we're in **k-space** (not real space!)

*Arrows do not represent magnetic moments,
Vortex-like spin textures have nothing to
do with Skyrmions, antiskyrmions, etc*

Spin texture of
 Bi_2Se_3 Dirac cone



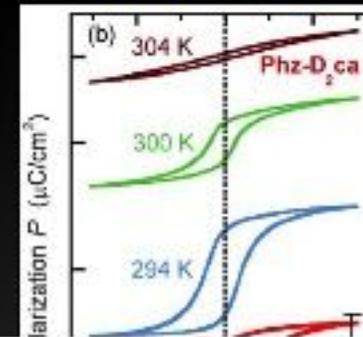
FERROELECTRICS: BASICS

- Polar materials, in which a spontaneous electric polarization can be switched via an external electric field

Ferroelectric Curie
Temperature for PE-FE
Transition

- 
- Electric "bistability":
✓ for non-volatile memory elements
 - By symmetry, all Ferroelectrics are non centrosymmetric (CS)

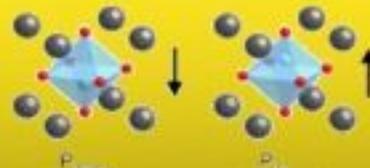
P-E hysteresis



- «All Insulator spin-orbitronics» (not based on spin-charge current conversion)
 - voltage-control (not current-control)
 - no Joule heating

KEY CONCEPTS & MAIN MESSAGE

FERROELECTRICITY



Permanent and switchable P, controllable via E

SPIN-ORBIT COUPLING



Relativistic interaction linking spin and lattice

- Via SOC in non-CS mat. → Spin-polarization in non-magnetic Materials
- Spin-texture switched via E fields in non-volatile way

✓ (Bulk)-RASHBA effects



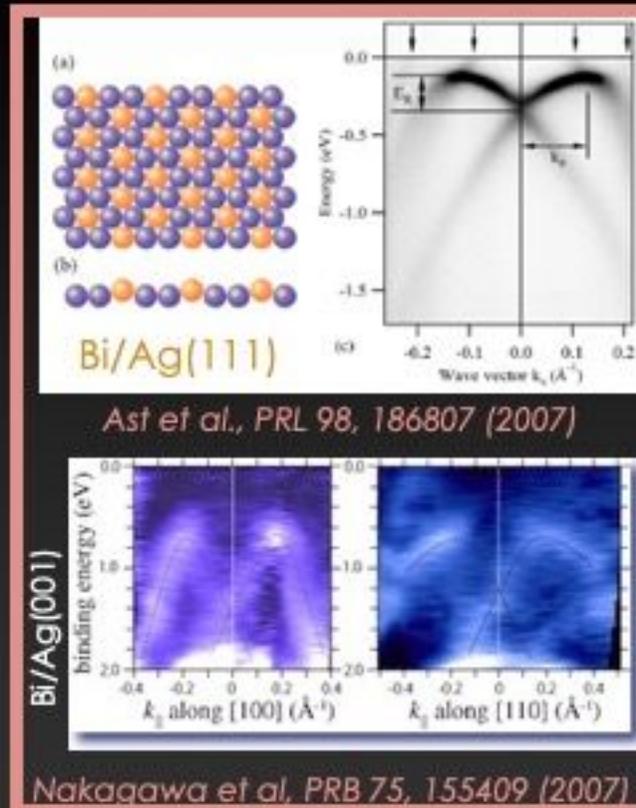
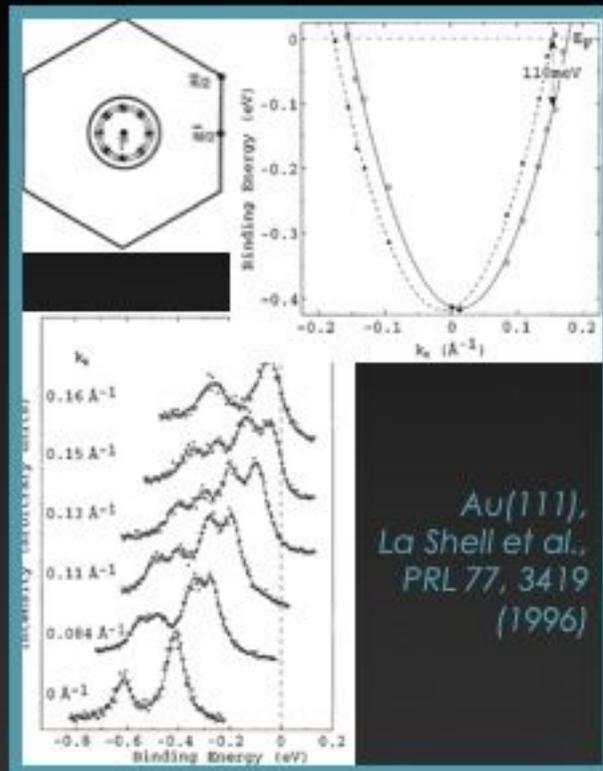
✓ SPIN-VALLEY coupling

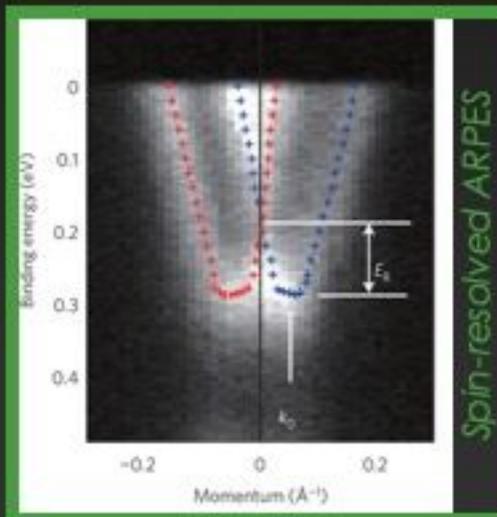
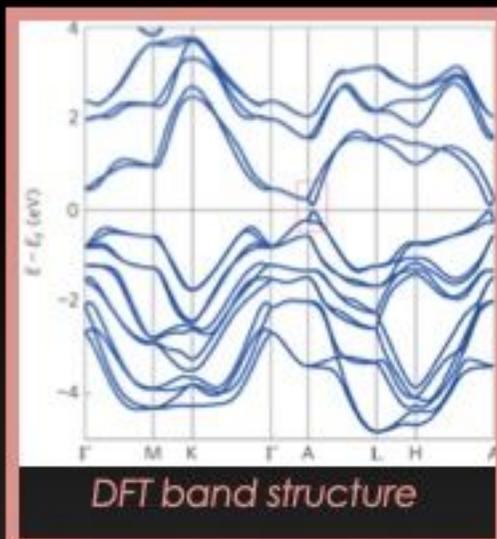
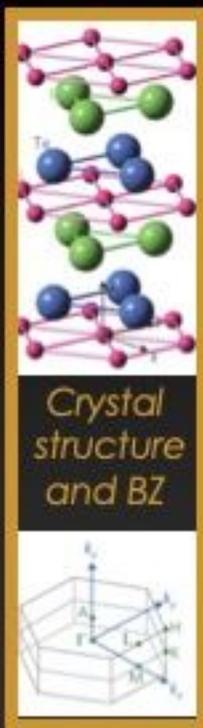


✓ TOPOLOGY



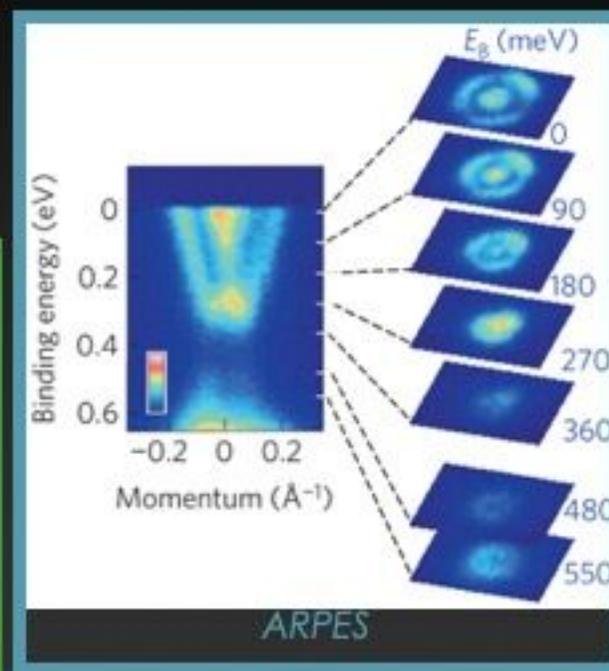
RASHBA EFFECT AT SURFACES





GIANT RASHBA EFFECT IN BULK BiTeI

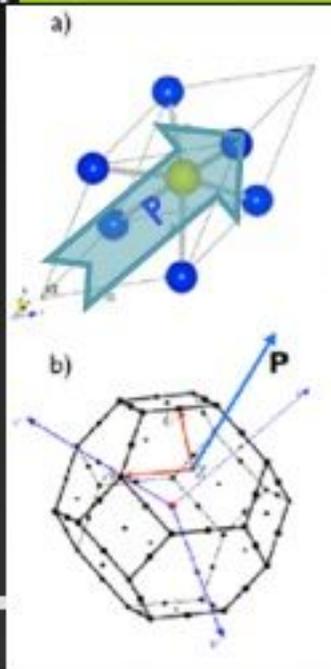
Tokura's group, *Nature Materials* 10, 521-526 (2011)



FERROELECTRIC GeTe: “BULK” RASHBA SPIN-SPLITTING

$T_C^{FE} \sim 720$ K
Space group:
R3m (no. 160)
P along [111]

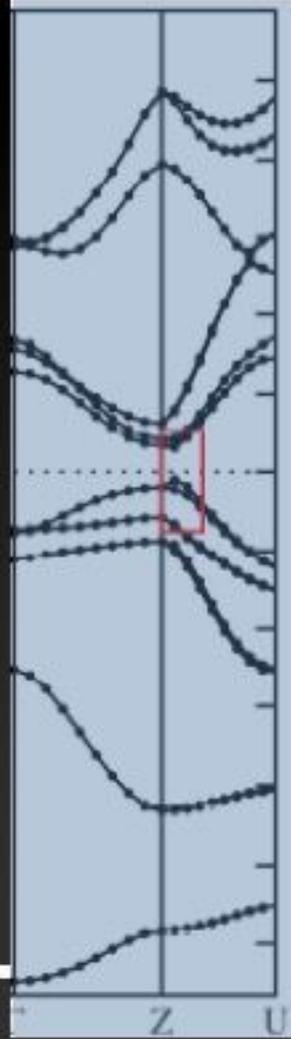
	Ω (Å ³)	a (Å)	a_0 (Å)	α (°)	θ (°)	τ	l (Å)	s (Å)	E_g (eV)	P ($\mu\text{C}/\text{cm}^2$)
PBE	56.08	4.373	6.080	57.76	88.02	0.0305	3.25	2.85	0.33	64.56
HSE	54.58	4.333	6.018	57.61	87.88	0.0331	3.24	2.81	0.65	70.27
Exp	53.31 ^a		5.98 ^a		88.35 ^a	0.0248 ^c	3.13 ^b	2.80 ^b	0.61 ^{c,d}	



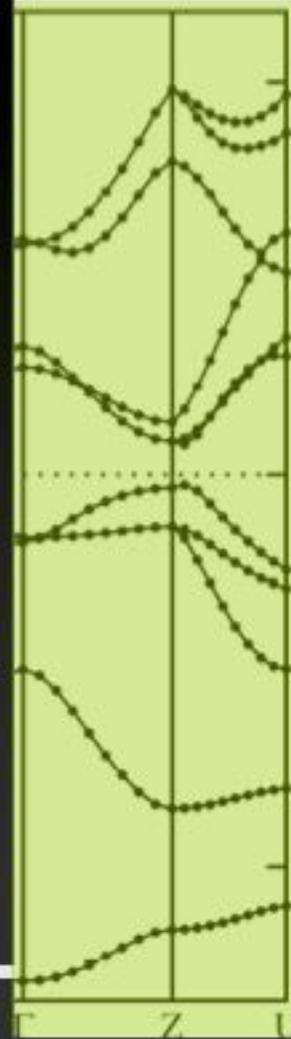
ORIGIN OF RASHBA SPIN SPLITTING

- ✓ Presence of spin-orbit coupling

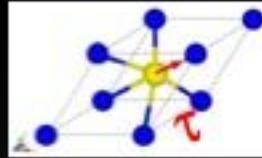
BAND STRUCTURE WITH SOC



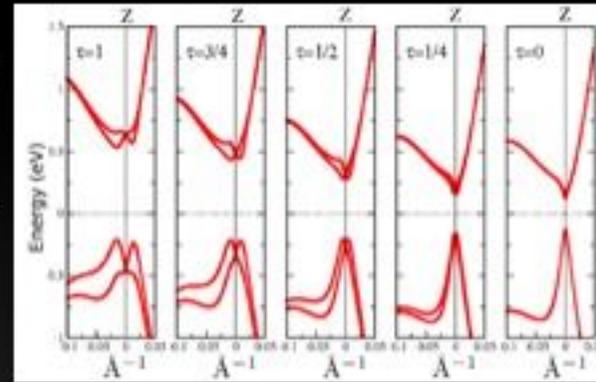
BAND STRUCTURE WITHOUT SOC



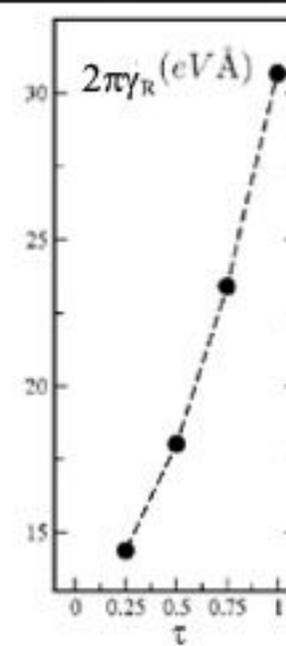
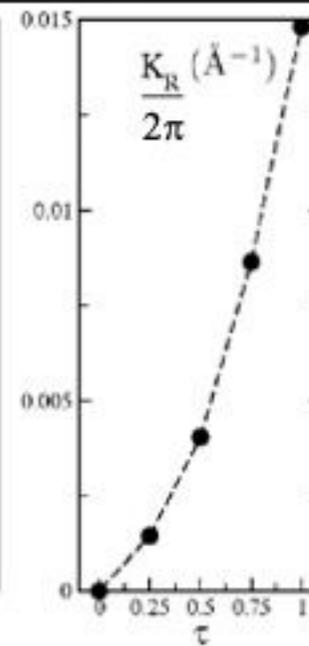
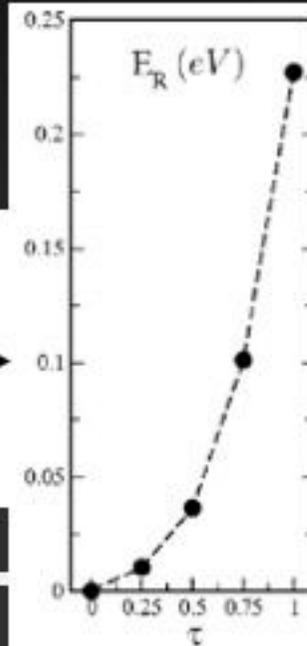
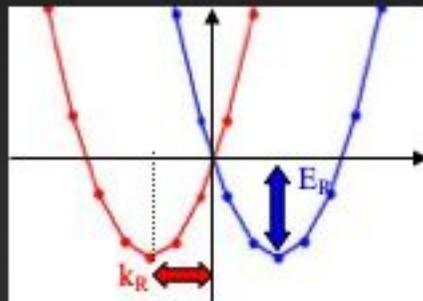
EVALUATION OF RASHBA PARAMETERS



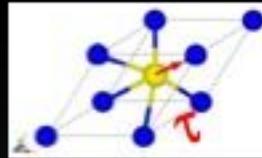
$$E_{\pm}(k_{\parallel}) = \frac{\hbar^2 k_{\parallel}^2}{2m^*} \pm \gamma_R |k_{\parallel}| + E_0, \quad k_{\parallel} = (k_x, k_y)$$



$$k_R = |m^*| \gamma_R / \hbar^2$$



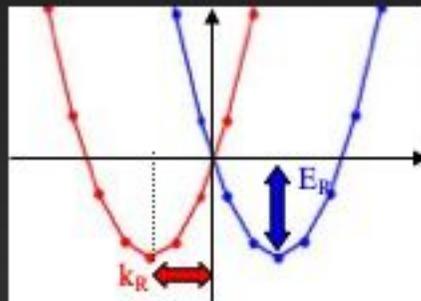
EVALUATION OF RASHBA PARAMETERS



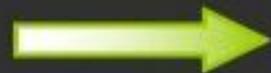
The effect is "GIANT" in GeTe!

$$E_{\pm}(k_{\parallel}) = \frac{\hbar^2 k_{\parallel}^2}{2m^*} \pm \gamma_R |k_{\parallel}| + E_0, \quad k_{\parallel} = (k_x, k_y)$$

$$k_R = |m^*| \gamma_R / \hbar^2$$



Sample	$k_0 (\text{\AA}^{-1})$	$E_R (\text{meV})$	γ_R (eV\AA)
Surface state			
Au(111)	0.012	2.1	0.33
Bi(111)	0.05	14	0.55
1/3 ML Bi on Ag surface alloy	0.13	200	3.05
Interface			
InGaAs/InAlAs	0.028	<1	0.07
QW state			
Pb thin film (6-22 ML)	0.035	$\lesssim 10$	0.04
Bi thin film (7-40 BL)	-	-	-
1 ML Bi on Cu	N/A	N/A	2.5
Bulk			
BiTeI	0.052	100	3.8

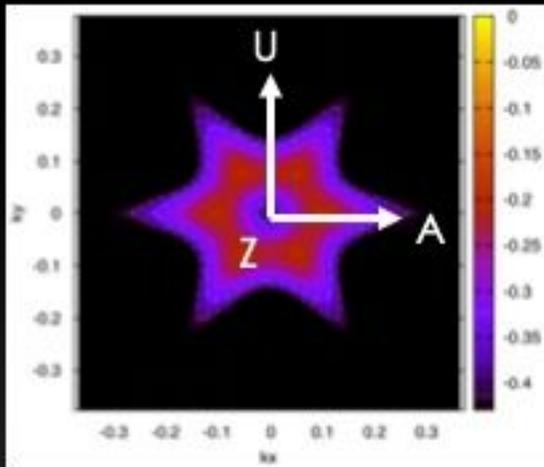


Ferroelectric GeTe

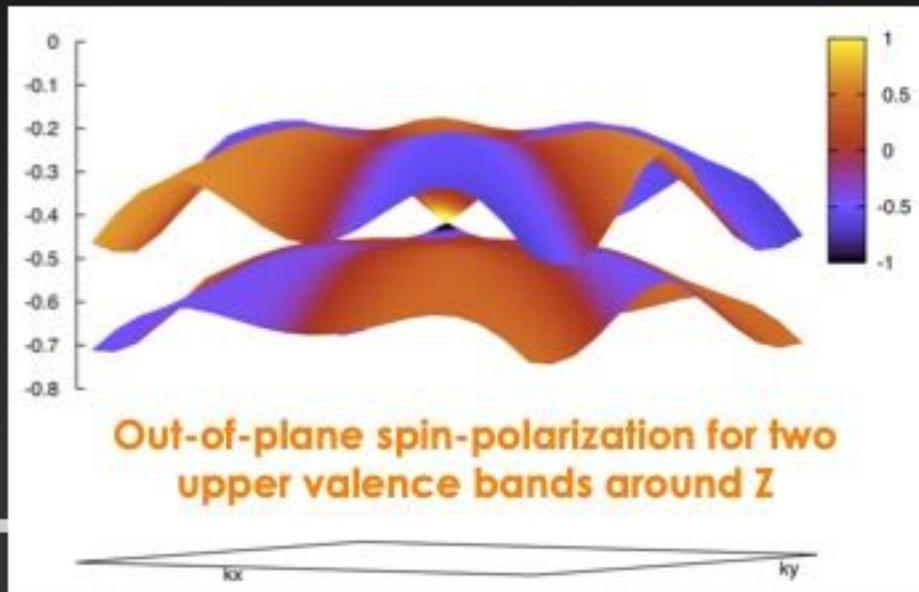
0.09 227 4.8



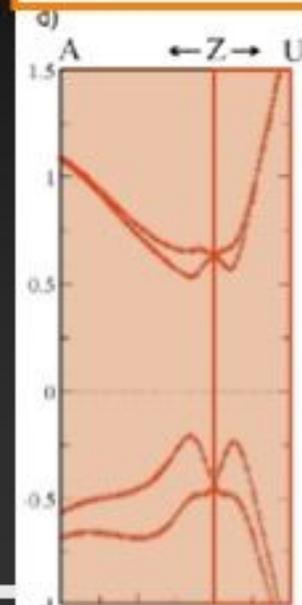
RASHBA EFFECT AT TOP VALENCE BANDS



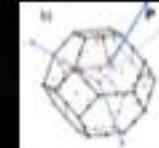
Isoenergy cuts off valence bands around Z



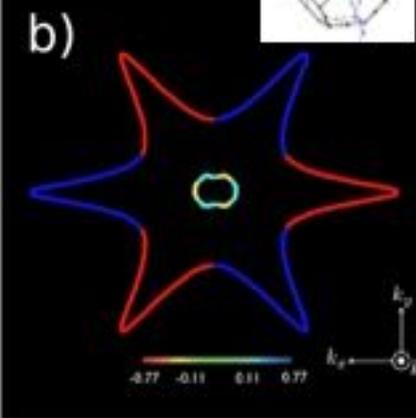
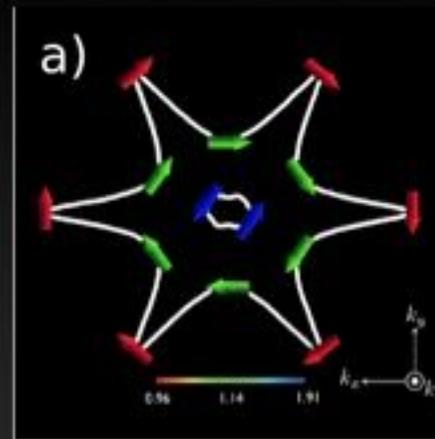
Out-of-plane spin-polarization for two upper valence bands around Z



SPIN POLARIZATION ALONG CONSTANT ENERGY CONTOURS AROUND Z



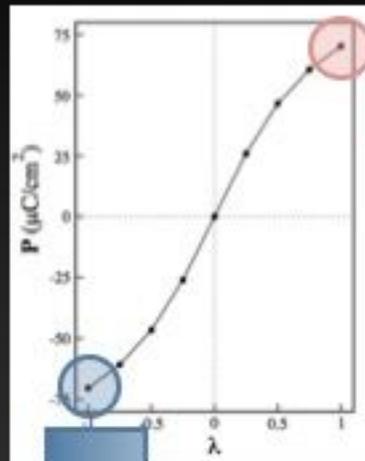
- a) **In-plane spin texture** (constant-energy cut at -0.47 eV). Arrows refer to the in-plane orientation of spin (colors indicate the modulus of the spin polarization).



- b) **Out-of-plane spin component distribution**



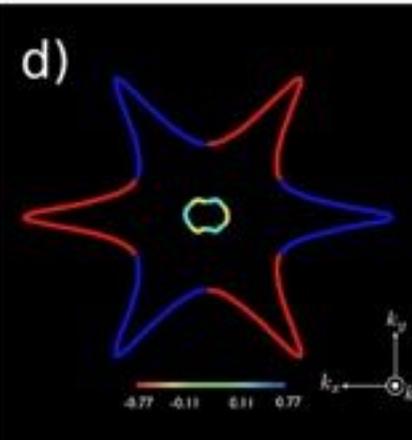
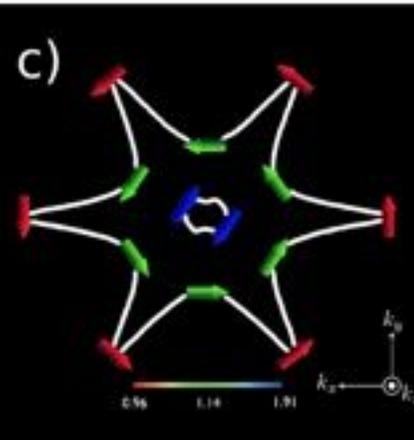
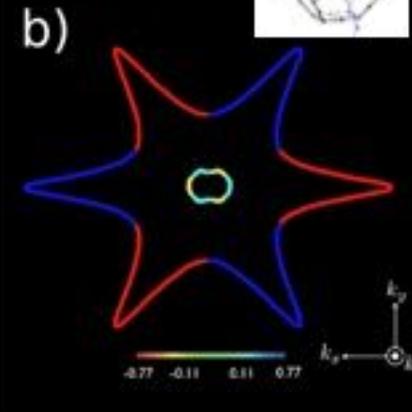
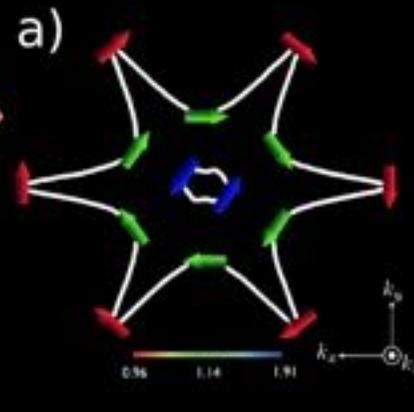
SPIN POLARIZATION ALONG CONSTANT ENERGY CONTOURS AROUND Z



+P

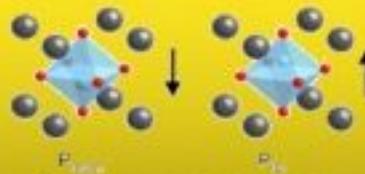
-P

c) and d) in-plane and out-of-plane spin components after switching of P



KEY CONCEPT & MAIN MESSAGE (I)

FERROELECTRICITY



Permanent and switchable P , controllable via E

SPIN-ORBIT COUPLING



Relativistic interaction linking **spin** and **lattice**

RASHBA EFFECT

SPIN TEXTURE: LINKED to P (via SOC) → CONTROL and SWITCH via E in a permanent (non volat.) way

a) Rashba (P_{up})



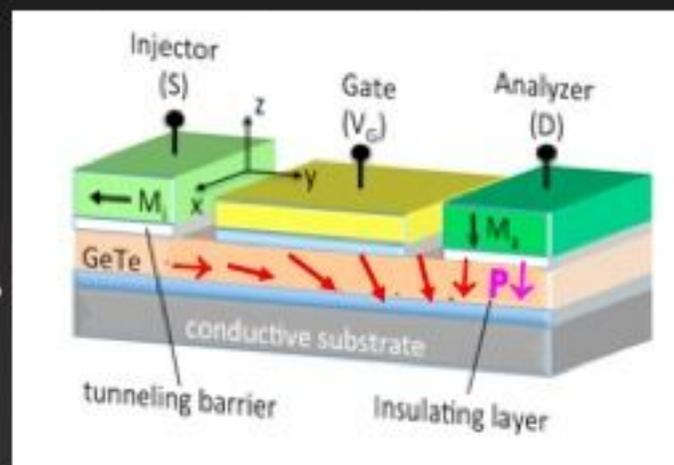
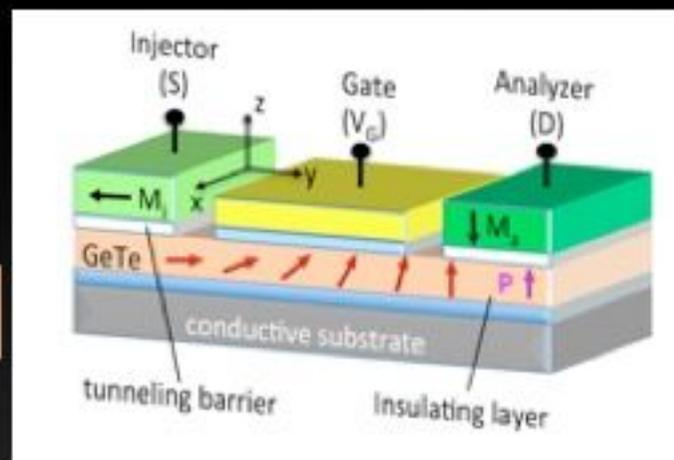
b) Rashba (P_{dw})



A POSSIBLE DEVICE

Datta-Das SPIN-FET

Polarization
switching



RECONFIGURABLE LOGIC FUNCTIONS USING FERSC-BASED SPIN-FET

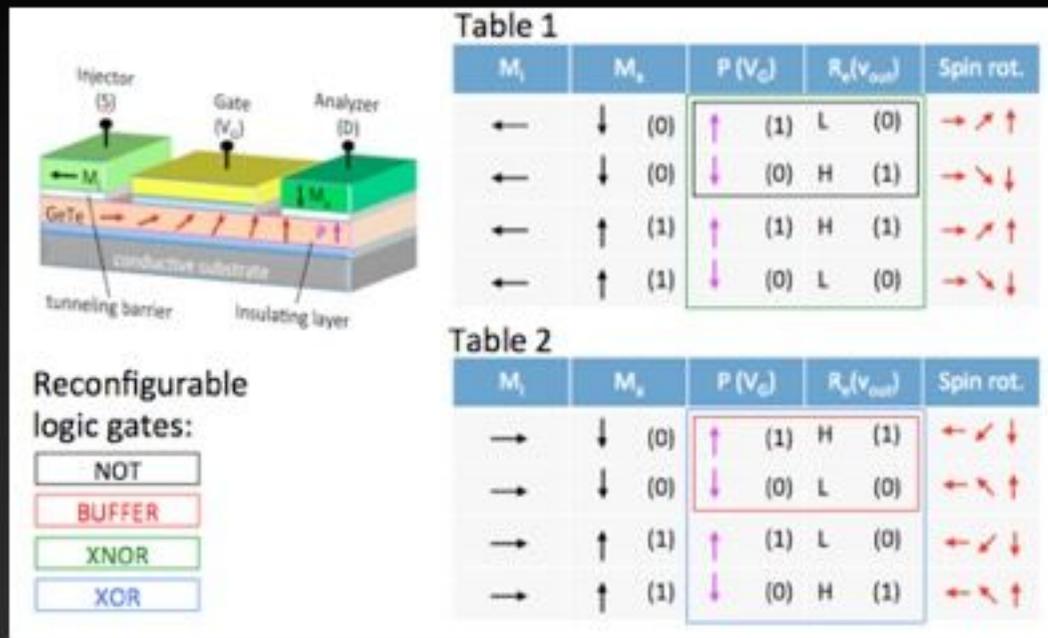


Table of truth for a spin-FET based on GeTe, operated in different configurations of the magnetization (M_i , M_o) in the two FM electrodes (injector, analyser) and of the dielectric polarization (P) in the GeTe channel. The high (low) device resistance (R_e) state H (L) is associated to a logic state 1 (0)



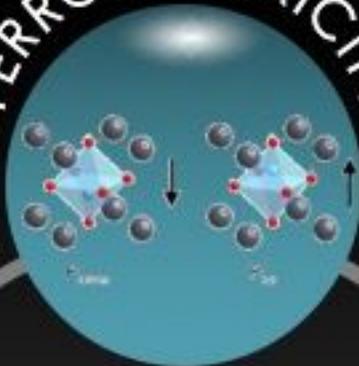
COEXISTENCE OF DIFFERENT PHYSICS PHENOMENA

- Adds the spin degree of freedom
- Traditionally exploited for **logic** elements

RASHBA EFFECTS



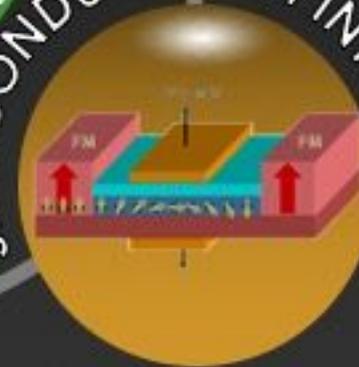
FERROELECTRICITY



- Non-volatility
- Switchability
- Control via **E**
- **Memory** element

FERROELECTRIC RELATIVISTIC SEMICONDUCTORS (FERSC)

SEMICONDUCTOR SPINTRONICS

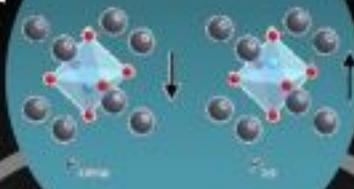


- Integration with device technology



THE QUEST FOR MULTI-FUNCTIONALITY

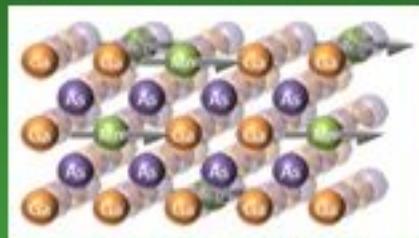
FERROELECTRICITY



- Non-volatility
- Switchability
- Control via **E**
- **Memory** element

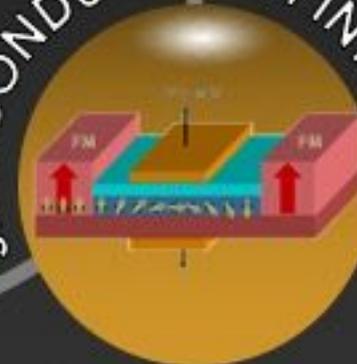
- Integration with device technology

Past attempt to use semiconductors for non-volatile applications:
DILUTED MAGNETIC SEMICONDUCTORS



However, operating temperatures are too low !

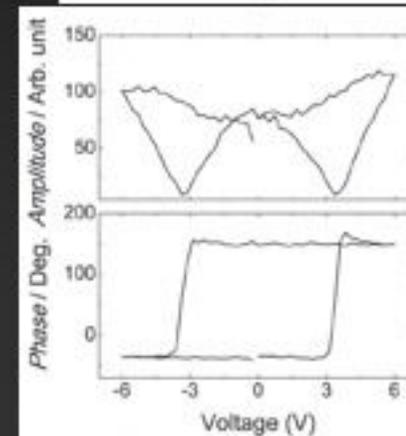
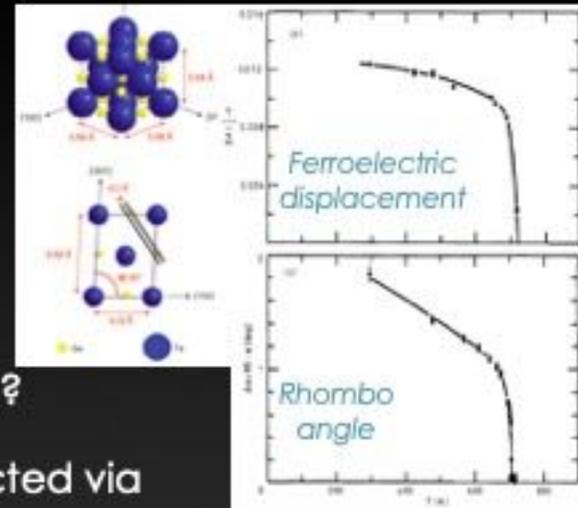
SEMICONDUCTOR SPINTRONICS



GeTe WORKS “IN THE COMPUTER”, BUT

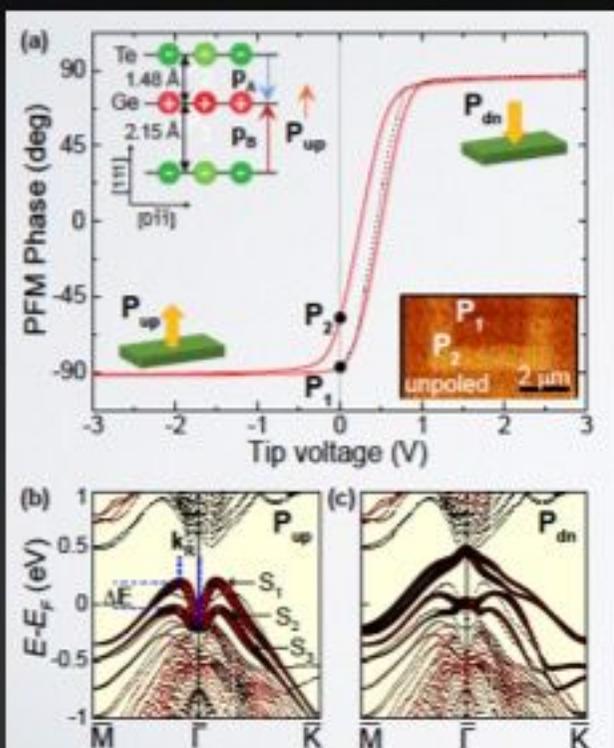
Shortcomings:

- low electrical resistivity (high density of Ge vacancies, leading to a density of free holes of up to $\sim 10^{20} \text{ cm}^{-3}$)
- a too small band gap (0.1 eV ? 0.6 eV ? indirect)
- structural ferroelectric distortions detected via neutron diffraction but no ferroelectric polarization ever measured
- recently, **FE switching via PFM** (*Gruverman et al, APL Mater (2014)*) on MBE films grown on highly lattice-mismatched **Si(111)** substrates
- **Rhombohedrally** distorted rocksalt structure α -GeTe

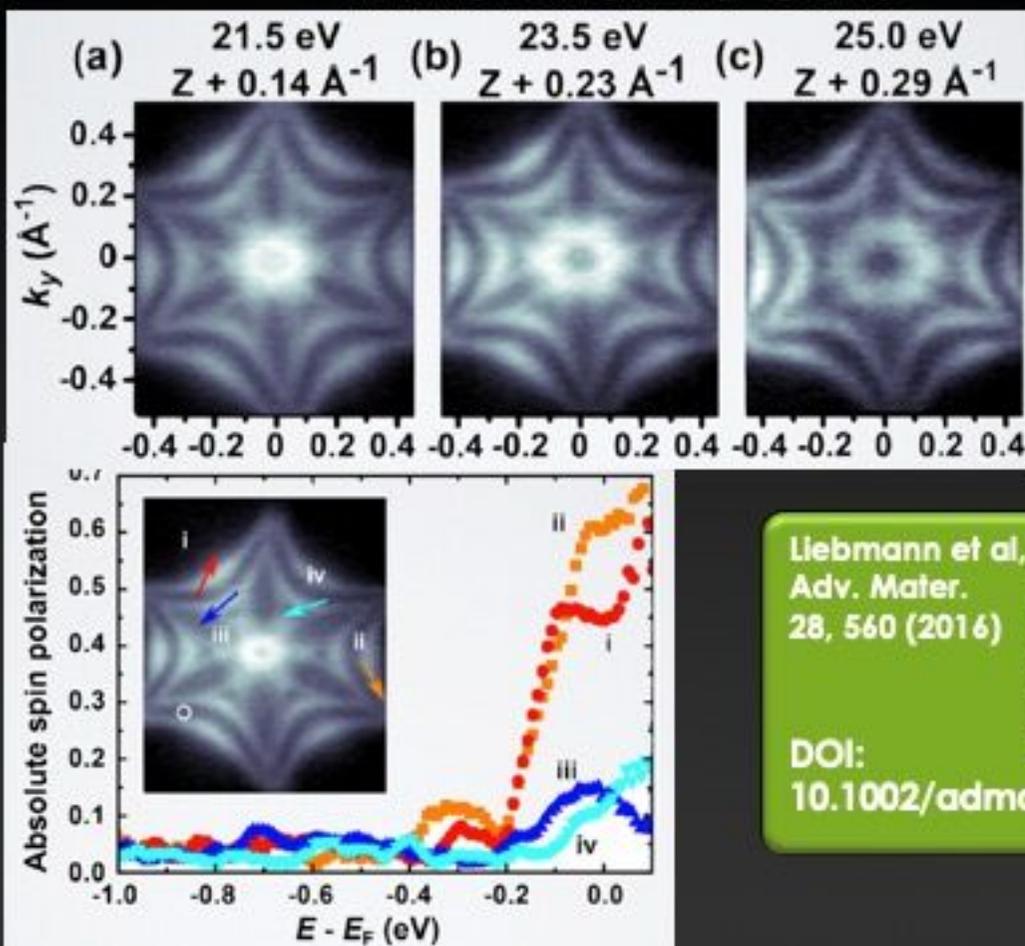


EXPERIMENTAL CONFIRMATION: SURFACE AND BULK RASHBA BANDS IN FE GeTe

M. Liebmann, C. Rinaldi, D. Di Sante, A. Giussani, R. Wang, S. Bertoli, M. Cantoni, L. Baldrati, I. Vobornik, G. Panaccione, R. Calarco, S. Picozzi, R. Bertacco and M. Morgenstern

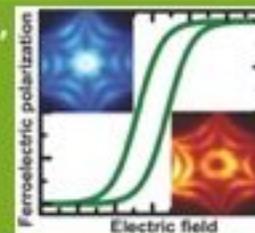


SURFACE AND BULK RASHBA BANDS IN FERROELECTRIC GeTe



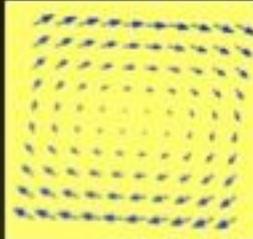
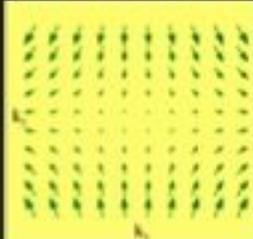
Liebmann et al,
Adv. Mater.
28, 560 (2016)

DOI:
10.1002/adma.201503459



SPIN SPLITTING: RASHBA VS DRESSELHAUS

Historically, two kind of SOC-induced spin-splitting effects:

Rashba Effect	Dresselhaus effect
<i>E. I. Rashba, Sov. Phys. Solid State 2, 1109 (1960).</i>	<i>G. Dresselhaus, Phys. Rev. B 100, 580 (1955)</i>
	
Structural inversion Asymmetry (SIA) <i>i.e. Interfaces, surfaces, 2DEG</i>	Bulk Inversion Asymmetry (BIA) <i>i.e. Zincblende, wurtzite crystal</i>

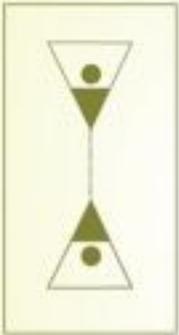
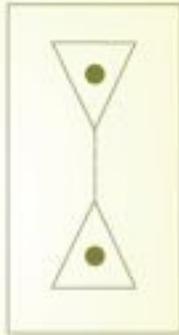
... but Rashba-like spin-textures
in bulk 3D crystals have been reported....



SPIN SPLITTING: RASHBA VS DRESSELHAUS

Unified description of spin-polarization effects:

"Because SOC is a relativistic effect, anchored on particular nuclear sites in the solid, it is the symmetry of such individual atomic sites in the solid that forms a good starting point to describe the SOC-induced spin polarization effect, rather than the global symmetry of the unit cell (as in SIA or BIA)"

	a Centrosymmetric		b Non-centrosymmetric (bulk inversion asymmetry)		c Centrosymmetric	
Bulk symmetry:	Inversion symmetry		Dipole field	Inversion asymmetry	Dipole field	Inversion asymmetry
Site symmetry:	Inversion symmetry					
Symmetry schematic:						
Effect/consequence:	Absence of spin splitting and spin polarization		Site dipole field induced net spin polarization	Site inversion asymmetry induced net spin polarization	Site dipole field induced spin polarization compensated by its inversion counterpart	Site inversion asymmetry induced spin polarization compensated by its inversion counterpart
Name:			R-1	D-1	R-2	D-2

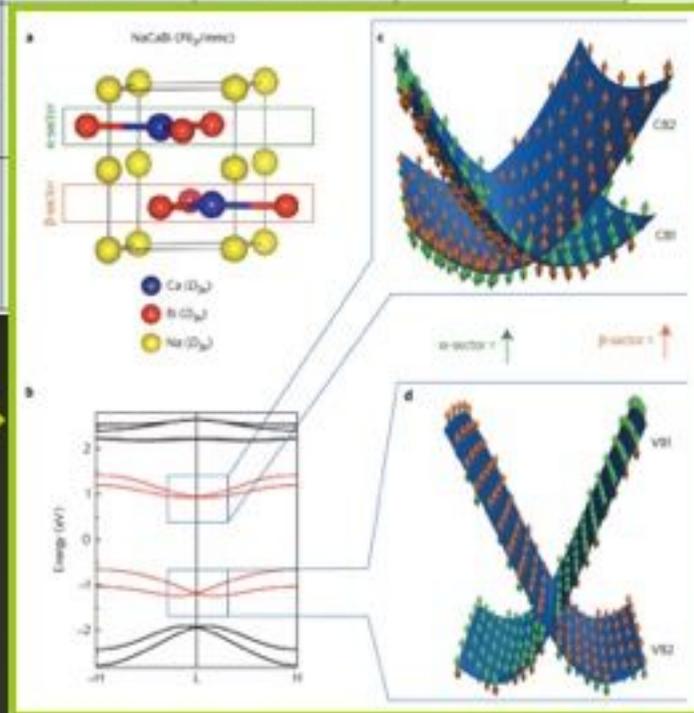
Zhang et al, Nature Physics 10, 387 (2014)



SPIN SPLITTING: RASHBA VS DRESSELHAUS

Table 1 | Classification of spin polarization in nonmagnetic bulk materials on the basis of bulk space group and site point group.

Bulk space group	Site point group	Non-centrosymmetric (at least one site)		Centrosymmetric (all sites) ($C_1, C_2, D_2, C_{2v}, C_{3v}, C_{4v}, C_{6v}, C_{2h}, C_{3h}, C_{6h}, C_{\infty v}, C_{\infty h}$)	
		Non-polar (all sites) ($D_2, D_2d, D_4, D_4d, S_6, D_{2d}, C_{2v}, D_{3h}, T, T_d, O$)	Polar (at least one site) ($C_1, C_2, C_3, C_4, C_6, C_{2v}, C_{3v}, C_{4v}, C_{6v}$)		
			Dipoles add up to zero		Dipoles add up to non-zero
Non-centrosymmetric (for example, $F\bar{4}3m$)	a D-1 Example: GaAs, ZrCoBi			Centrosymmetric (all sites) ($C_1, C_2, D_2, C_{2v}, C_{3v}, C_{4v}, C_{6v}, C_{2h}, C_{3h}, C_{6h}, C_{\infty v}, C_{\infty h}$)	
Centrosymmetric (for example, $R\bar{3}m$)	d D-2 Example: Si, NaClB				



Zhang et al, Nature Physics 10, 387 (2014)



**ARE THERE
OTHER
FERSC?**

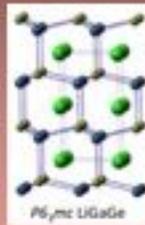
**CANDIDATE
MATERIALS**

SnTe: TCI or FERSC?



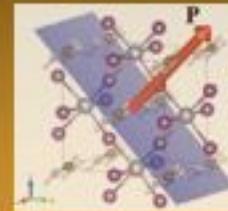
**FERROELECTRIC
RASHBA
SEMICONDUCTORS
(FERSC)**

ABC SEMICONDUCTORS



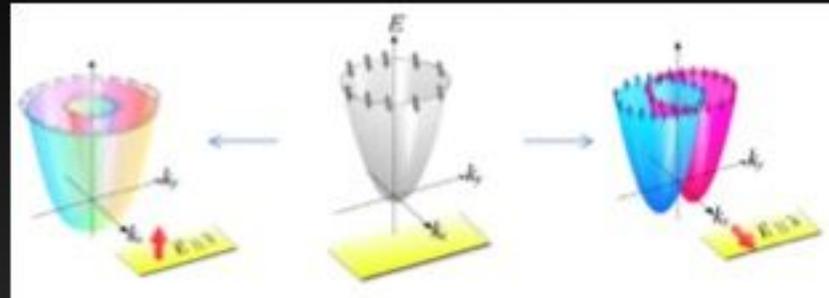
KMGbI

"SOLAR" HALIDE-PEROVSKITES

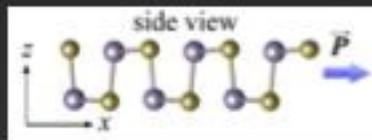


CAN WE USE OXIDES ?

- **Motivations:**
 - Better **insulators** (i.e. *ferroelectrics*) wrt GeTe (quite "leaky" semicond.)
 - **Different directions for polarization** (... P is often in-plane for thin-films due to depolarization fields)



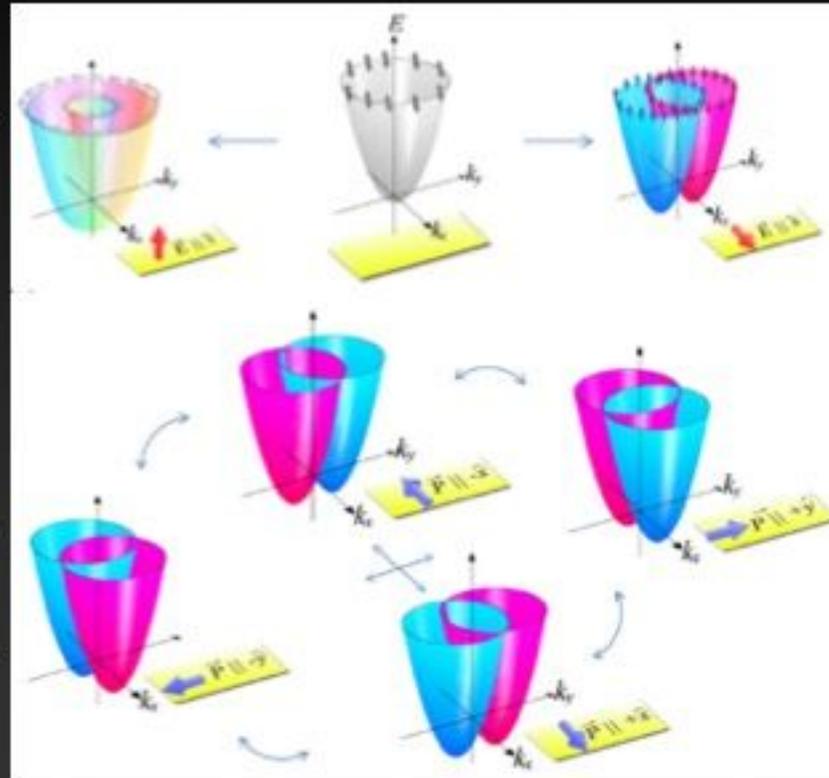
What happens if P is in-plane?



H. Lee et al., arxiv:1712.06112v3
on FE SnTe thin-films

CAN WE USE OXIDES ?

- **Motivations:**
 - Better **insulators** (i.e. *ferroelectrics*) wrt GeTe (quite "leaky" semicond.)
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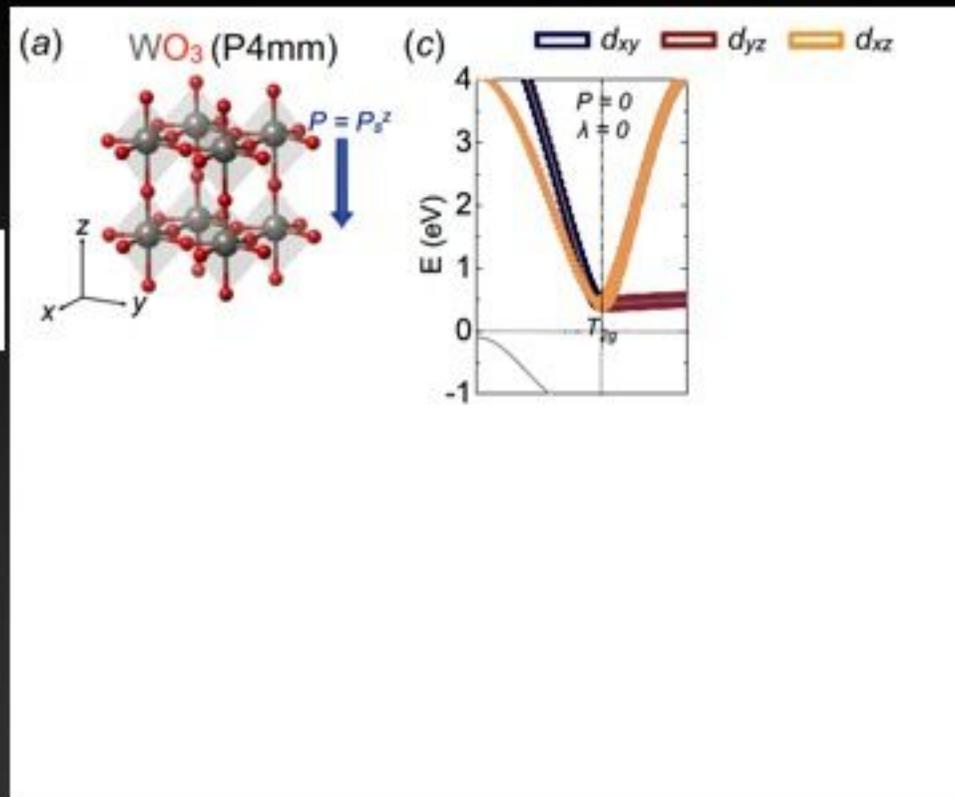


TUNGSTEN-BASED OXIDES (I)

- WO_3 : ground state (non-polar) $P2_1/c$, but (polar) orthorhombic $Amm2$ only 20 meV/f.u. higher in energy.
- (Simpler) tetragonal polar $P4mm$ also close (74 meV/f.u.)

✓ $P=0$, no SOC:

Triply degenerate t_{2g}



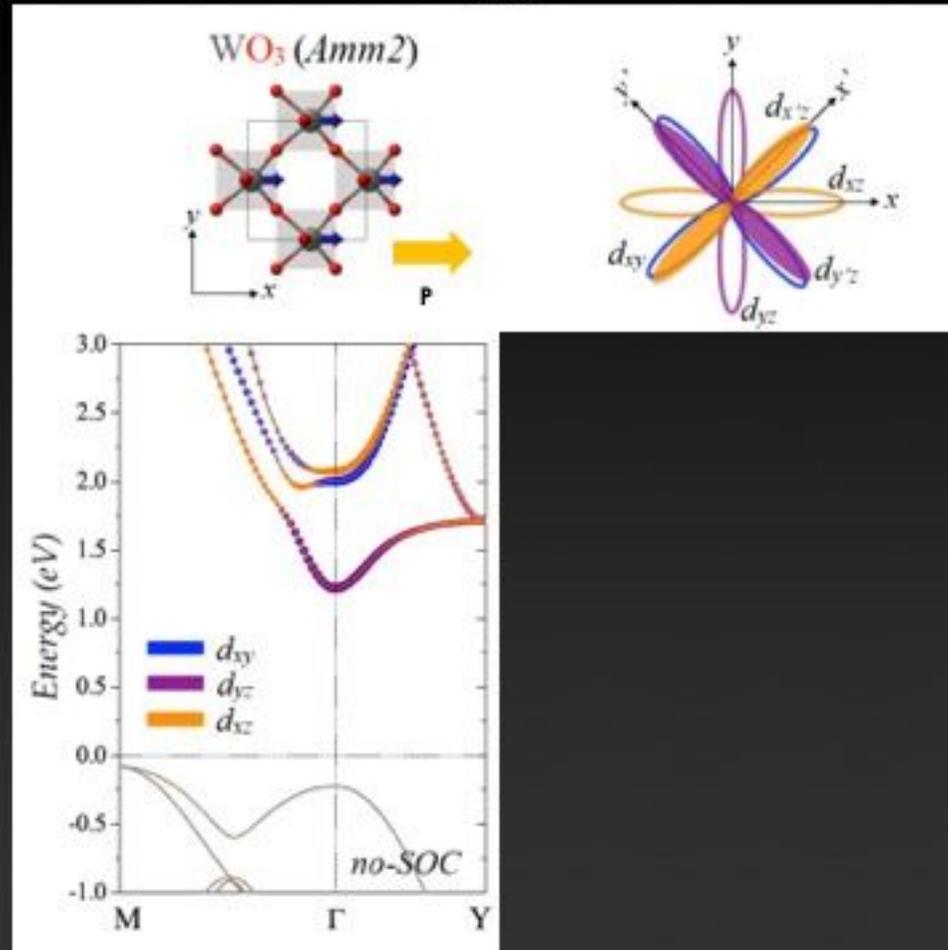
TUNGSTEN-BASED OXIDES (II)

- WO_3 : ground state (non-polar) $P2_1/c$, but (polar) orthorhombic $Amm2$ only 20 meV/f.u. higher in energy.
- (Simpler) tetragonal polar $P4mm$ also close (74 meV/f.u.)

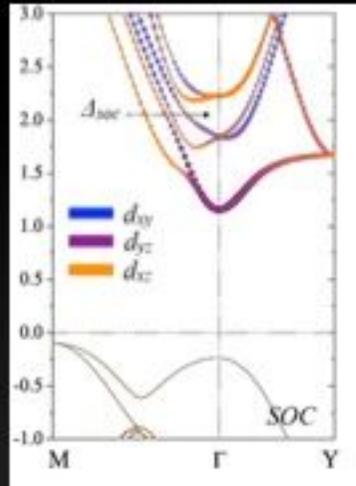
d_{xz} : Highest energy band, points towards P

d_{xy} : medium energy band, in the plane (also containing P)

d_{yz} : lowest energy band, perpendic. to P

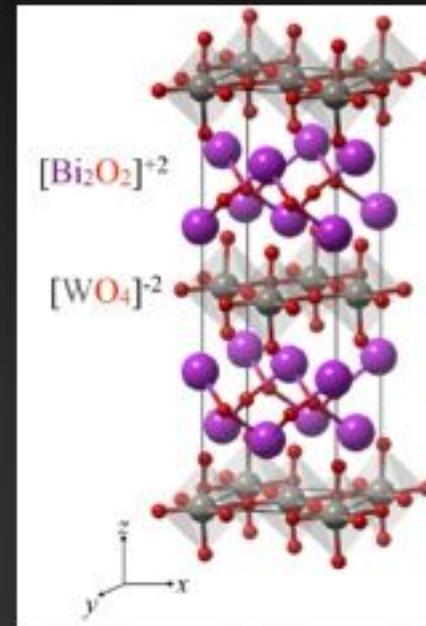


TUNGSTEN-BASED OXIDES (III)



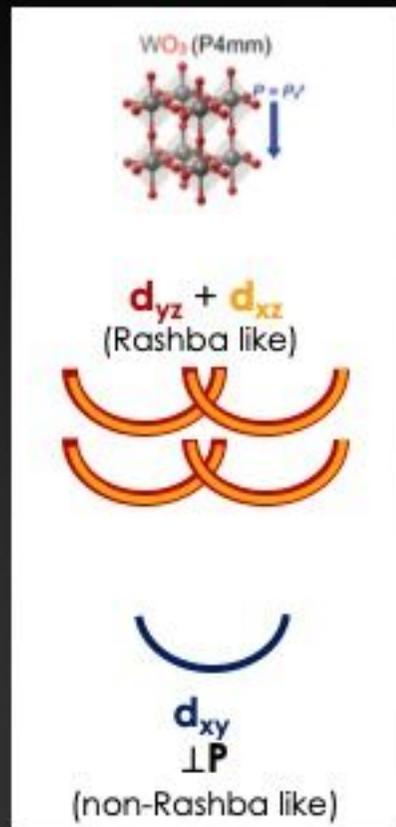
- **STRATEGY:** raise the energy of (non-Rashba-like) d_{yz} (perpendic. to P)
- **HOW?** Destabilize out-of-plane orbitals by changing the W-O-W out-of-plane coordination and keep d_{xy} as lowest (Rashba-like) state

- **POSSIBLE SOLUTION:** Aurivillius phase Bi_2WO_6
- **Layered structure:** $(\text{Bi}_2\text{O}_2)^{+2} (\text{A}_{m-1}\text{B}_m\text{O}_{3m+1})^{-2}$
 - Bi_2O_2 = fluorite layer
 - $\text{A}_{m-1}\text{B}_m\text{O}_{3m+1}$ = perovskite blocks
- Here $m=1$ (smallest member of Aurivillius phases, A-site deficient)
- Space group $F2mm$
- P along x $\sim 50 \mu\text{C}/\text{cm}^2$
- Quite wide gap (about 2.7 eV)



TUNGSTEN-BASED OXIDES (IV)

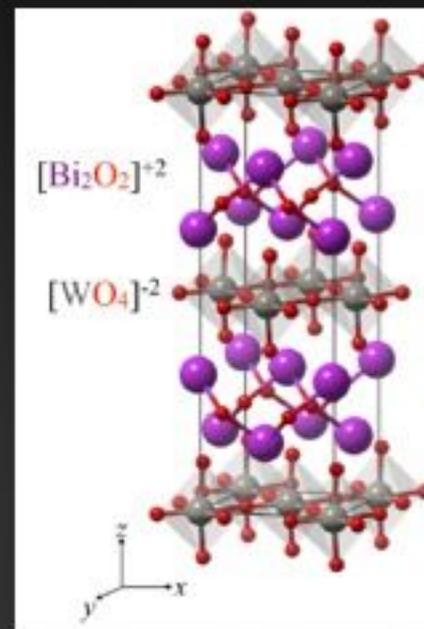
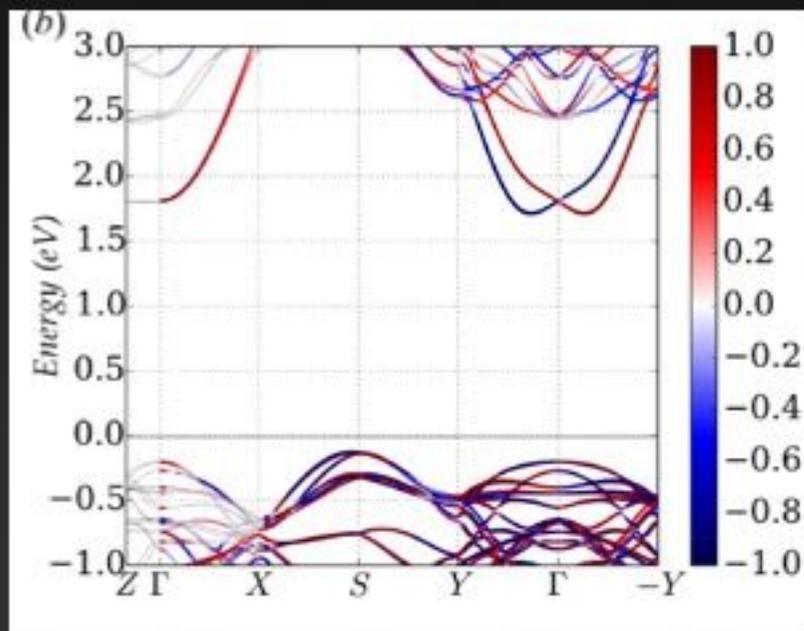
" t_{2g} Band engineering" in Ferroelectric oxides



TUNGSTEN-BASED OXIDES (V)

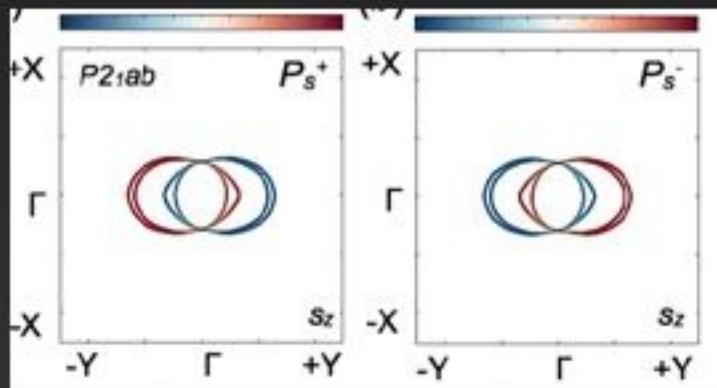
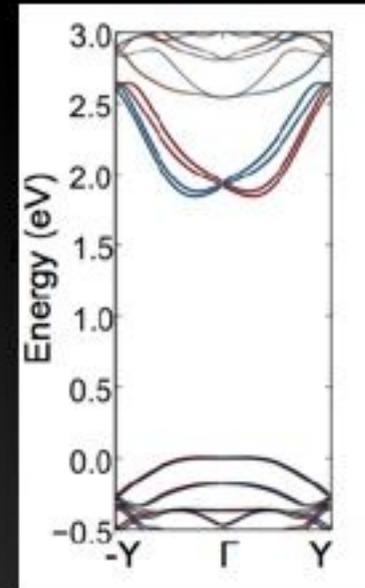
- $F2mm$ Bi_2WO_6 (Γ_5^- from tetragonal $I4/mmm$ PE)
- Small dispersion along Γ -Z: roughly "2D"
- Lowest band Rashba-like: $\Delta E = 104$ meV, $\Delta k = 0.14 \text{ \AA}^{-1}$, $\alpha_R = 1.45 \text{ eV\AA}$
 - Mainly d_{xy} , whereas d_{xz} and d_{yz} upper in energy

Spin-polarization s_z : red/blue color



TUNGSTEN-BASED OXIDES (VI)

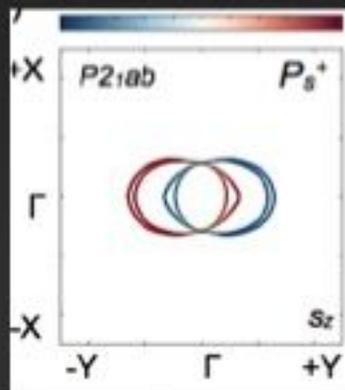
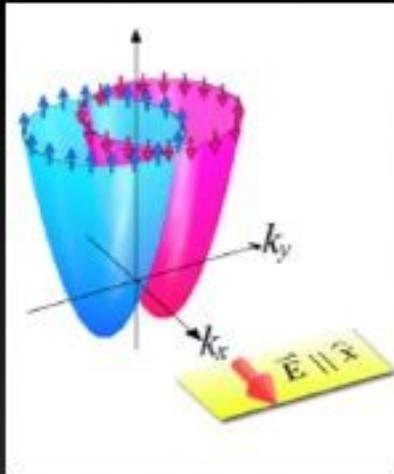
- Complex distortions/tilting in ground-state structure of Bi_2WO_6 :
 - X_2^+ : rotation of O octahedra around z
 - X_3^+ : tilt of O octahedra around x
 - Γ_5^- : polar zone center along x
- $F2mm$ (Γ_5^- from tetragonal $I4/mmm$ PE)
- $P2_1ab$ ($X_2^+ + X_3^+ + \Gamma_5^-$)



Polarization switching implies s_z switching:

Electric control of spin-texture !

TUNGSTEN-BASED OXIDES (VI)



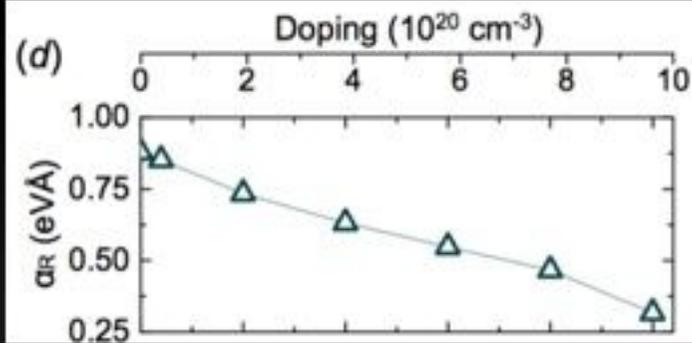
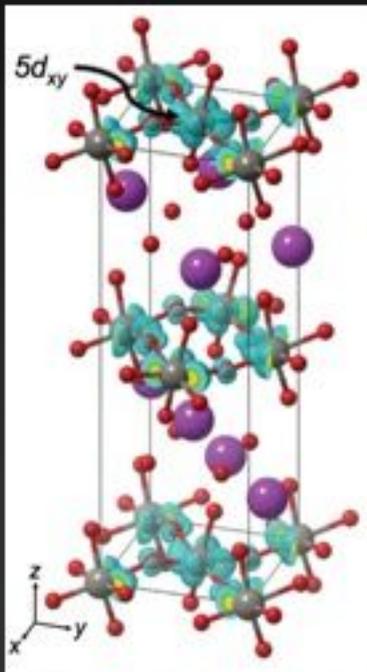
PERSISTENT SPIN HELIX

(well known in Semiconducting Quantum wells)

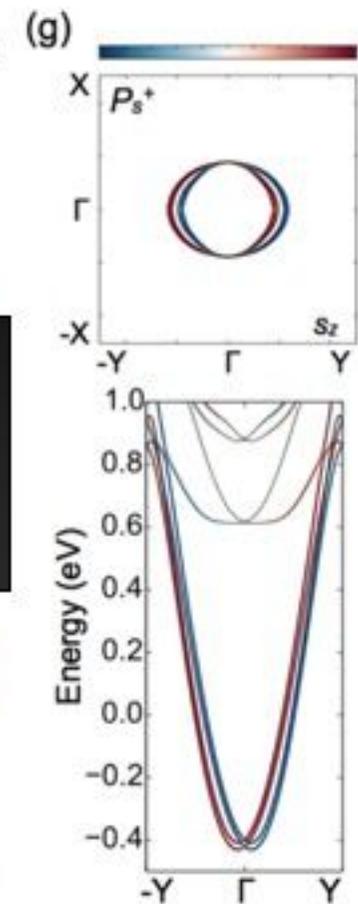
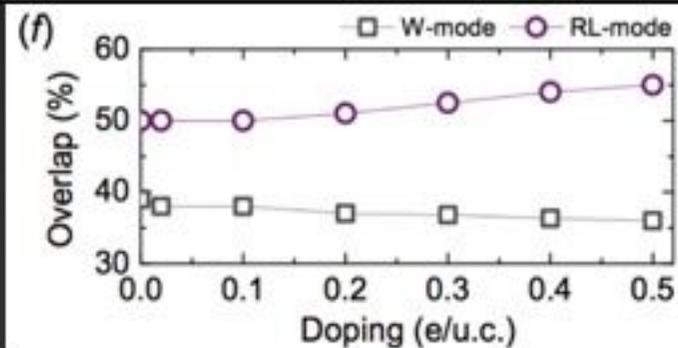
- ✓ Protection against spin scattering
- ✓ Long spin lifetime !

TUNGSTEN-BASED OXIDES (V)

- What happens upon electron doping?

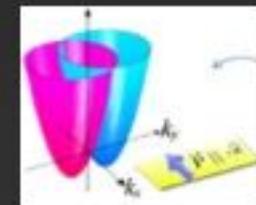
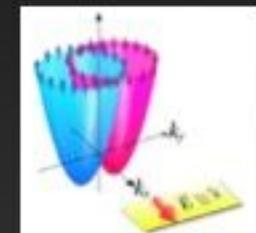
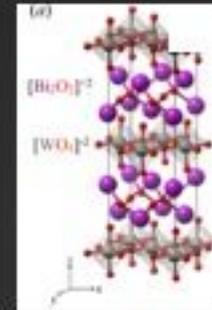
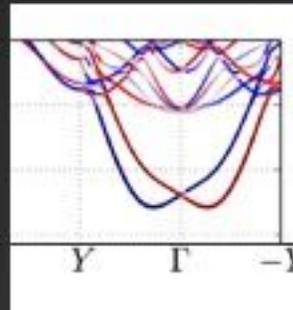
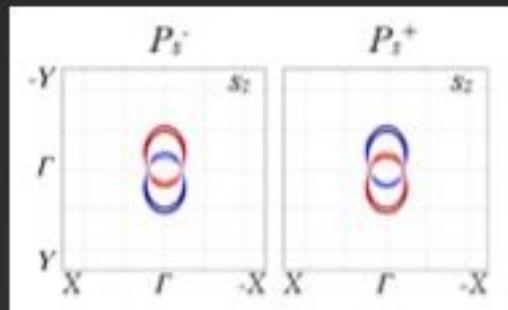
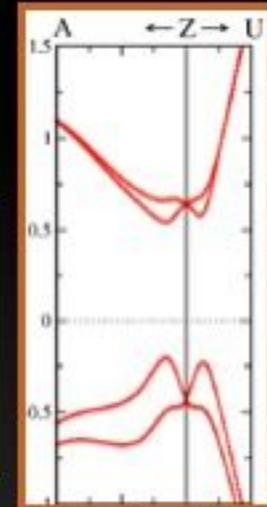
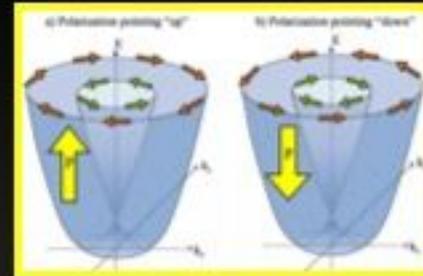


- Doping of 0.5 e/unit cell
- Spin-texture into $k_x k_y$ plane at $E = 2.0 \text{ eV}$



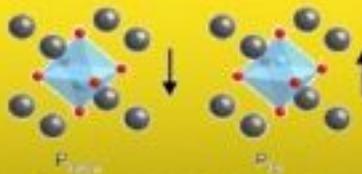
TAKE HOME MESSAGES (I)

- **GeTe**: first example of **FERSC** (**F**erroelectric **R**ashba **S**emi**C**onductor)
- Sense of spin-rotation: **opposite when P is switched**
- **Large Bulk Rashba** effect in **oxides**: Aurivillius-phase **Bi₂WO₆**: **PSH**
- Oxides (strong Ferroelectrics) offer possibilities to "**engineer**" SOC effects in band structures



KEY CONCEPT & MAIN MESSAGE (II)

FERROELECTRICITY



Permanent and switchable P , controllable via E

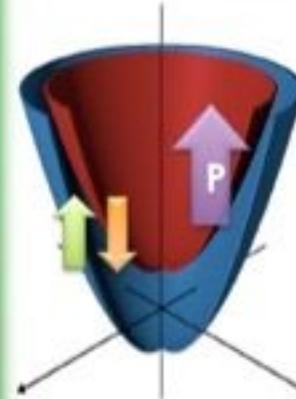
SPIN-ORBIT COUPLING



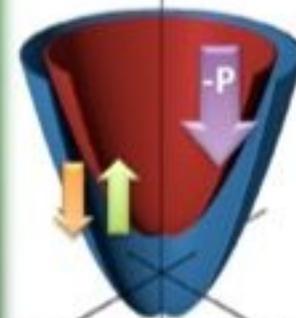
Relativistic interaction linking **spin** and **lattice**

SPIN-VALLEY

c) Zeeman (P_{up})



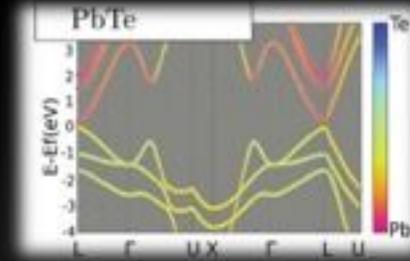
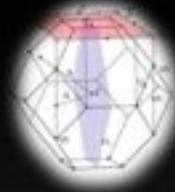
d) Zeeman (P_{dw})



INTRODUCING THE VALLEY CONCEPT

- **Valleys:** energy extrema in momentum space

Band structure of Rock-salt Chalcogenides - Valleys at TRIM L points



- Valley two-valued **pseudospin** : defined when valleys occur at points which are not time-reversal invariant momenta (TRIM)

ALL-OXIDE VALLEYTRONICS

Starting point: $(ABO_3)_2(AB'O_3)_n$ [111] heterostructure

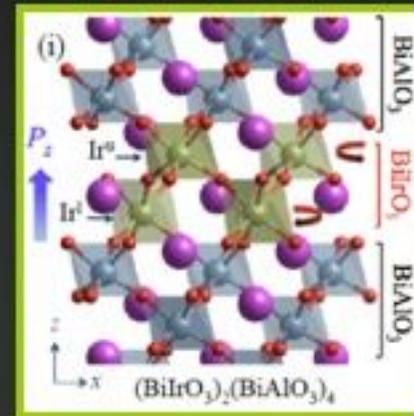
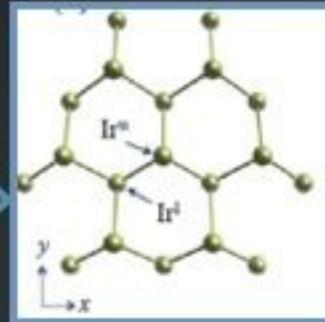
- ✓ (111) TM bilayer with honeycomb lattice
- ✓ 2 trigonal sublattices on different layers
- ✓ Lattice geometry supports Dirac points*



* Nagaosa, Okamoto, Nat. Comms (2011)

- Our **Receipt**: Embed **5d-oxide bilayer*** (large SOC) in a **(111) FE host**
- **FE $BiAlO_3$** ($P//[111]$, $R3c$)
- **Materials design:**
 $BiIrO_3$ in **$BiAlO_3$**

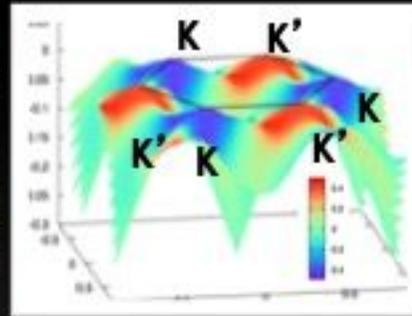
Top view:
Honeycomb
lattice



Side view:
Unit cell: $Bi_6Ir_2Al_4O_{18}$

BiIrO₃/BiAlO₃ BAND STRUCTURE

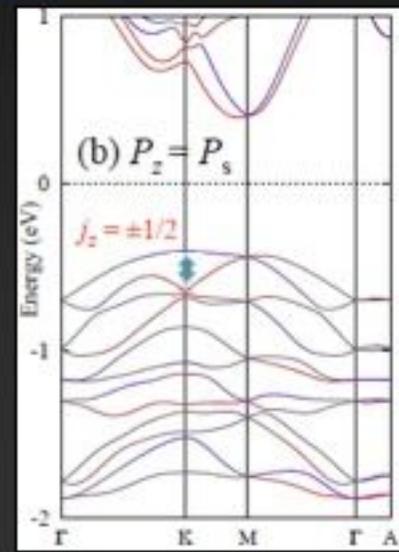
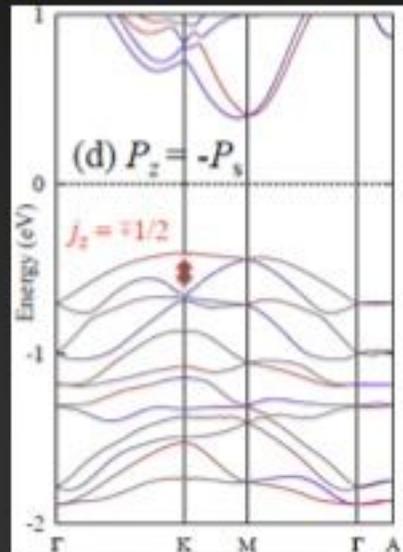
Exactly as in MoS₂:
spin-polarized
valleys!



Valence Band
Maximum
at K/K' with
opposite
spin-polarization

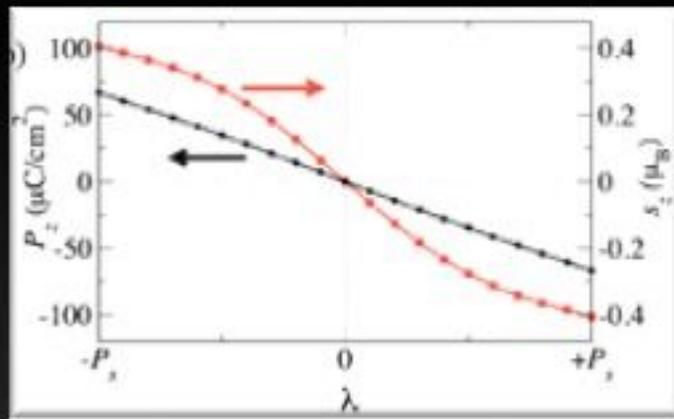
WHAT HAPPENS UPON POLARIZATION SWITCHING?

-P_s (w SOC)
* Spin-polarization switching



P_s (w SOC)
* VBM @K
(hybridization
A_{1g} and E_g')

COUPLING BETWEEN P AND SPIN-SPLITTING



—————→
polar ionic distortion

Electric polarization P_z and s_z polarization are strongly coupled
i.e. one can tune s_z by applying electric field !

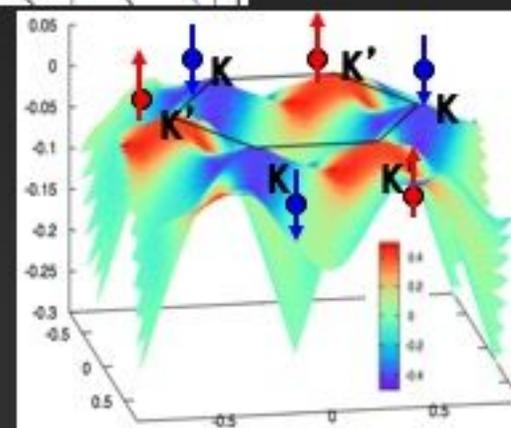
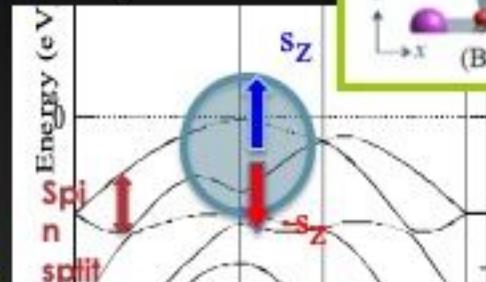
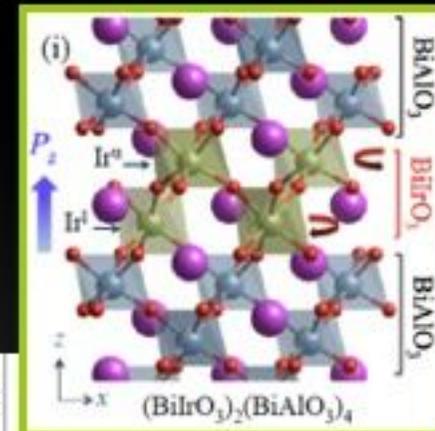
- ☛ Layer degree of freedom related to FE polarization
- ☛ **Spin-valley-layer coupling :**
Sort of "magnetoelectric" coupling
(Gong et al Nat Commun 4, 2053 (2013))



TAKE HOME MESSAGES (II)

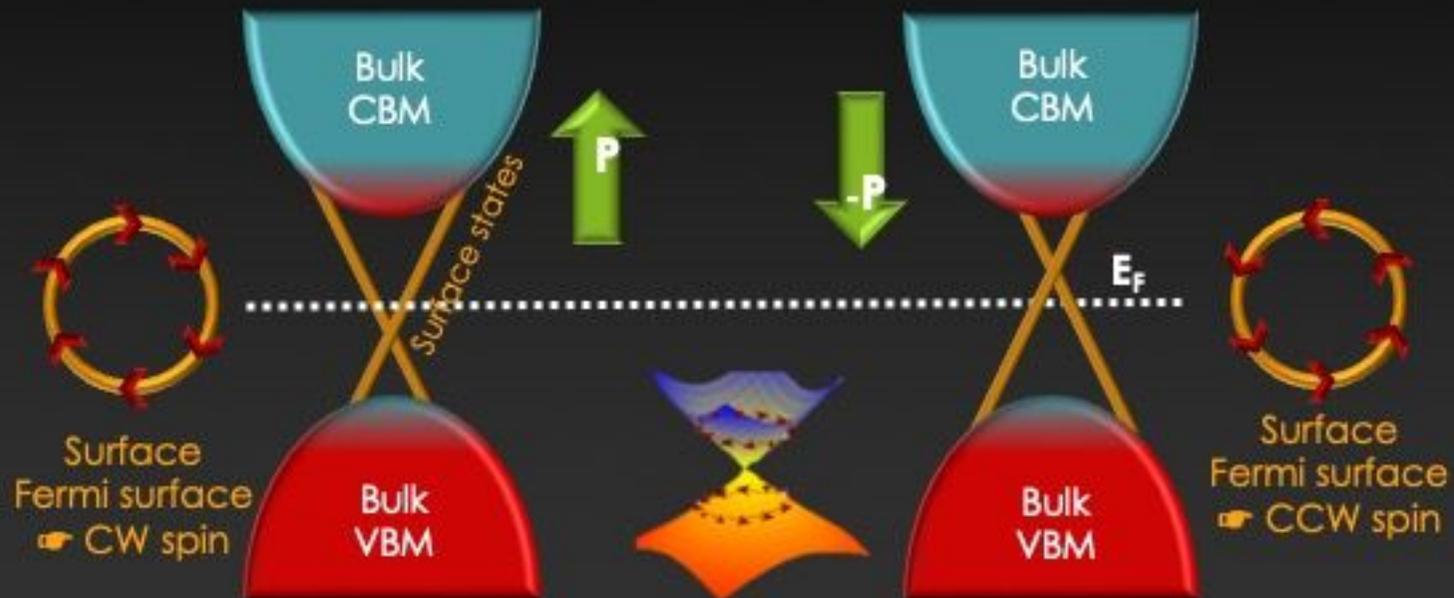
- Combination of **Ferroelectricity** with **SOC** leads to **Zeeman-like spin-splitting**
- **Receipt:** Embed a **5d-oxide bilayer** (large SOC) in **ferroelectric host**
- **BirO₃ bilayer in BiAlO₃:**
 - **Spin-polarized valley at K** (useful for **valleytronics**)
 - **S_z spin-polarization** can be **switched** by flipping **P**

✓ **ALL-OXIDE VALLEYTRONICS
IN FERROELECTRIC HETEROSTRUCTURE**

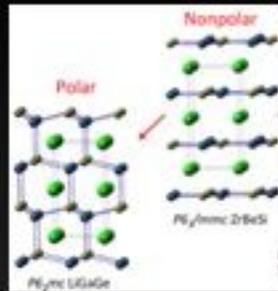


KEY CONCEPT & MAIN MESSAGE (III)

Topological Ferroelectric Semiconductor (SOC-induced band inversion)



"ABC" HYPER-FERROELECTRICS

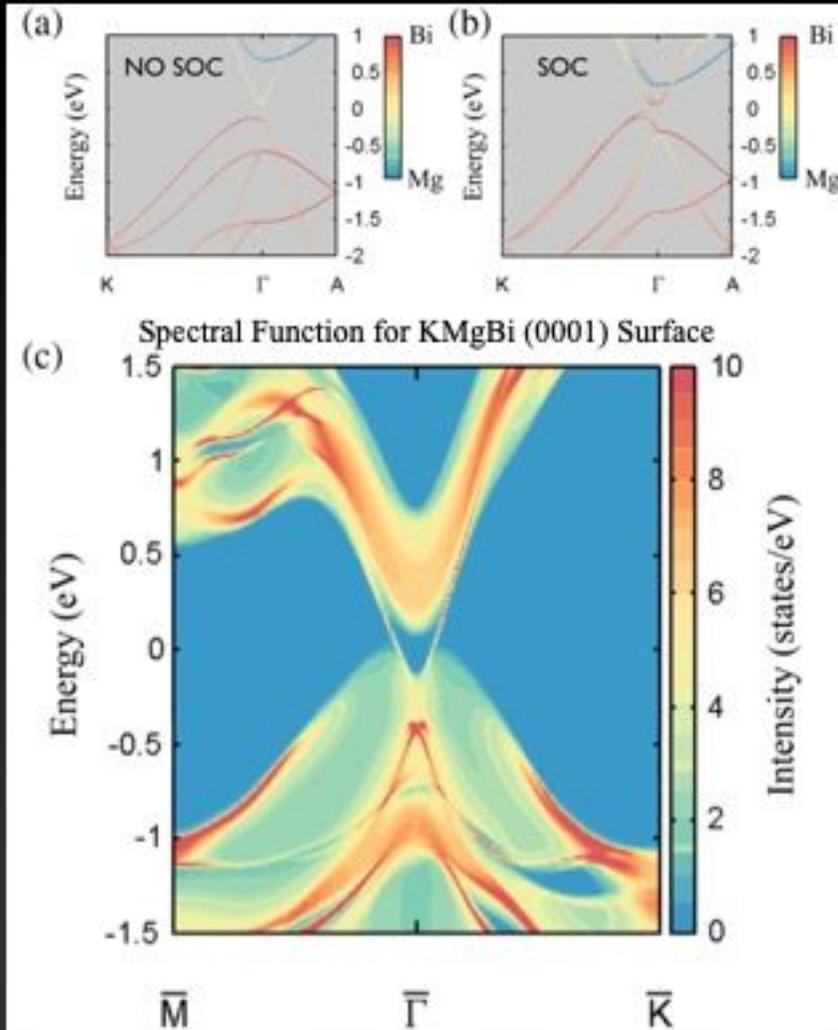


☞ FE that should maintain their **P** even in thin film limit

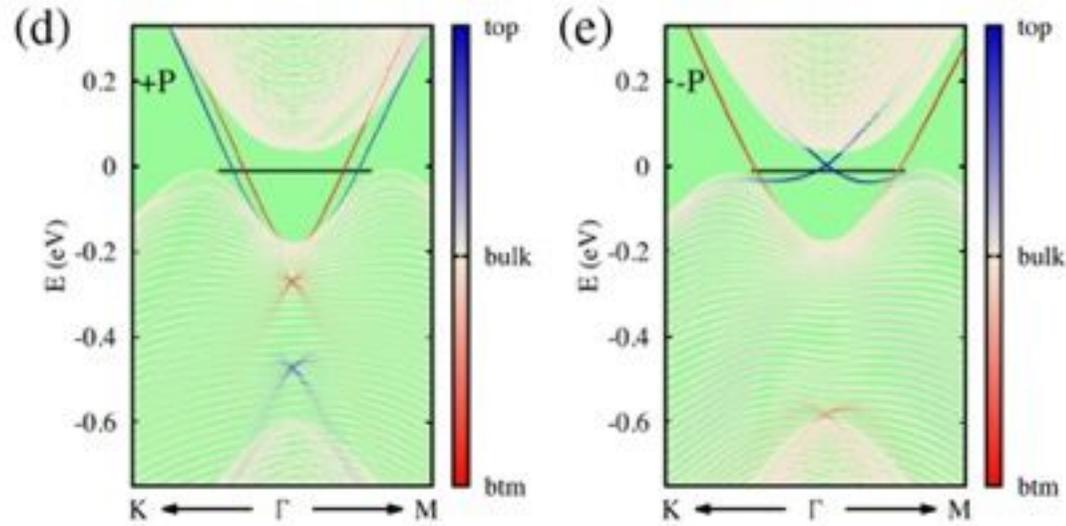
*J. Bennett et al.,
PRL 109, 167602 (2012)*

☞ Some of these (i.e. **KMgBi**) are (FE) Top. Ins.!

*D. Di Sante et al.,
PRL 117, 076401 (2014)*



"ABC" HYPER-FERROELECTRICS



- 4 term:
- K & MgBi;
 - P_{In} & P_{out}



Summary:
**WHY ARE RELATIVISTIC FERROELECTRICS
USEFUL AND INTERESTING?**

- From the **Fundamental** research point of view:
New mechanism for **cross-coupled spin-electric effect** =>
interplay between ferroelectric order and spin degree of freedom
- From the **Materials Science** point of view:
New class of **multifunctional Materials** (... and, as such, open
to understanding and optimization from theory, synthesis,
characterization, ...)
- From the **Device/Technology** point of view:
New class of **electrically controlled spintronic devices**
combining logic functions and storage

**THANK YOU FOR YOUR
ATTENTION !!**