

STEM EELS (Electron Energy-Loss Spectroscopy)

Odile Stéphan & Alexandre Gloter

STEM group

Laboratoire de Physique des Solides

Université Paris-Sud

<https://www.stem.lps.u-psud.fr/>

odile.stephan@u-psud.fr; alexandre.gloter@u-psud.fr



STEM imaging and EELS

Part I: introduction to Transmission Electron Microscopy / STEM imaging

Part II: spectroscopy at the atomic scale/ STEM EELS (Electron Energy-Loss Spectroscopy)

STEM imaging & EELS principles

A few generalities about transmission microscopy

Image formation principles in STEM and examples

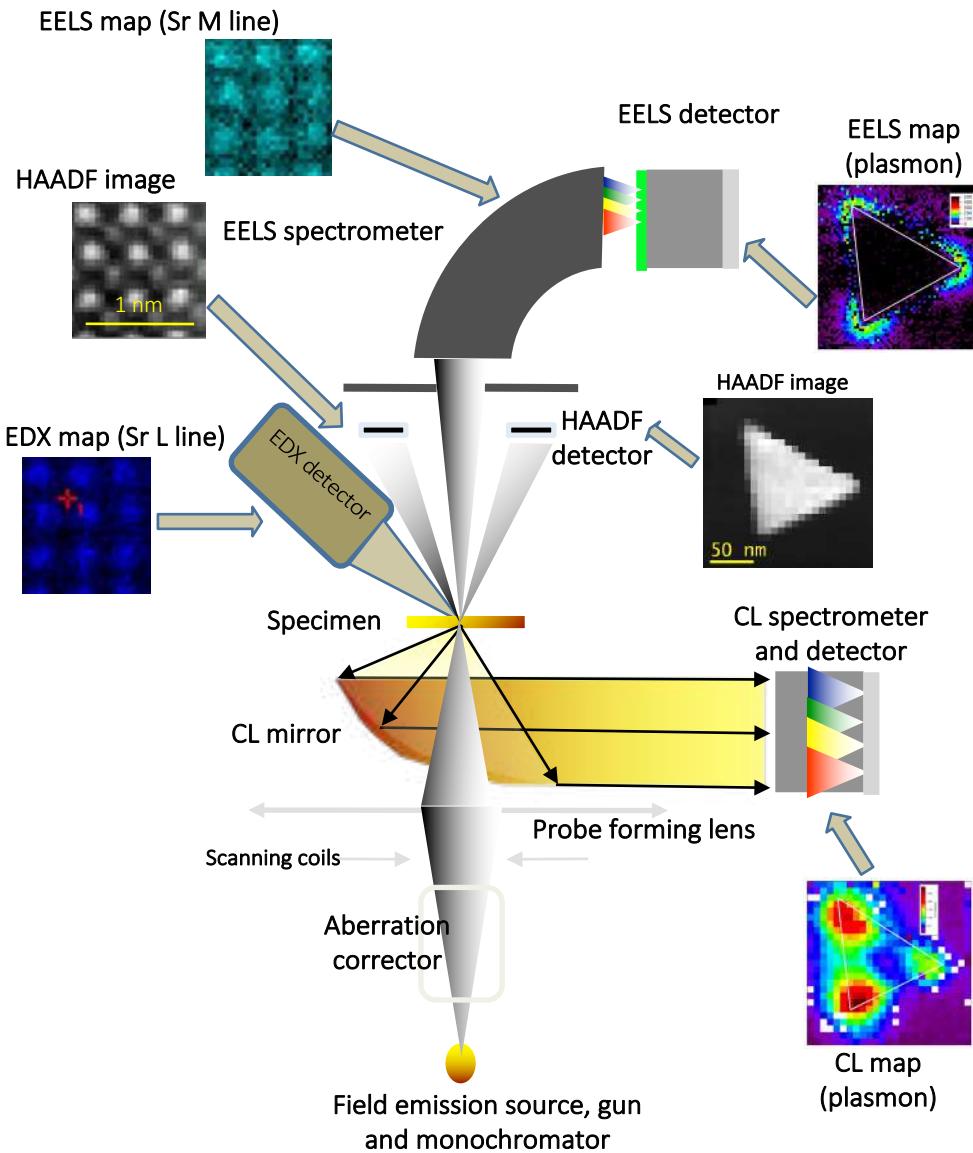
EELS principles and examples

Application to oxide thin films. Some examples below:

- 1) Octahedra rotation in $(\text{La},\text{Sr})\text{MnO}_3/\text{SrTiO}_3$ superlattices
- 2) Charge control in manganite by ferroelectric switching in $\text{CaMnO}_3/\text{BiFeO}_3$ based Mott transistor
- 3) Orbital ordering in $\text{LaAlO}_3/\text{SrTiO}_3$ bilayers

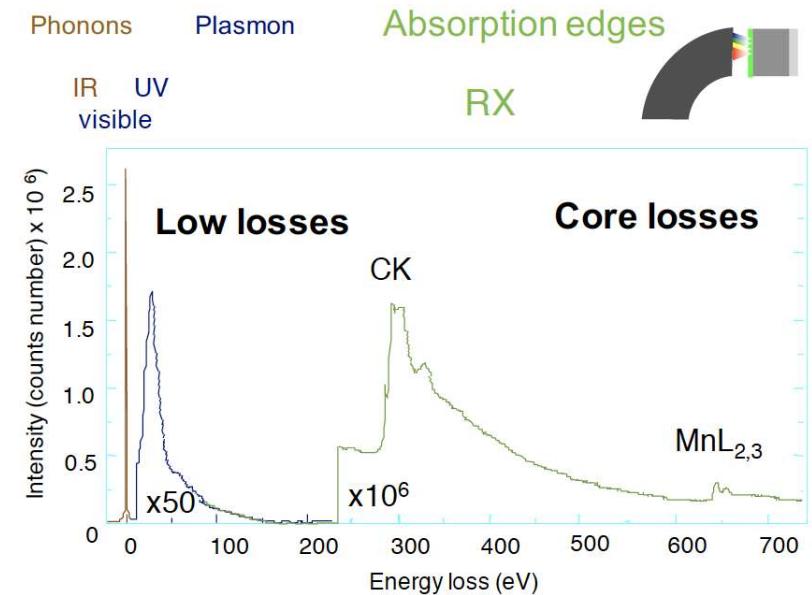
Few perspectives on STEM-EELS for oxides / interfaces studies

Spectromicroscopy in a STEM

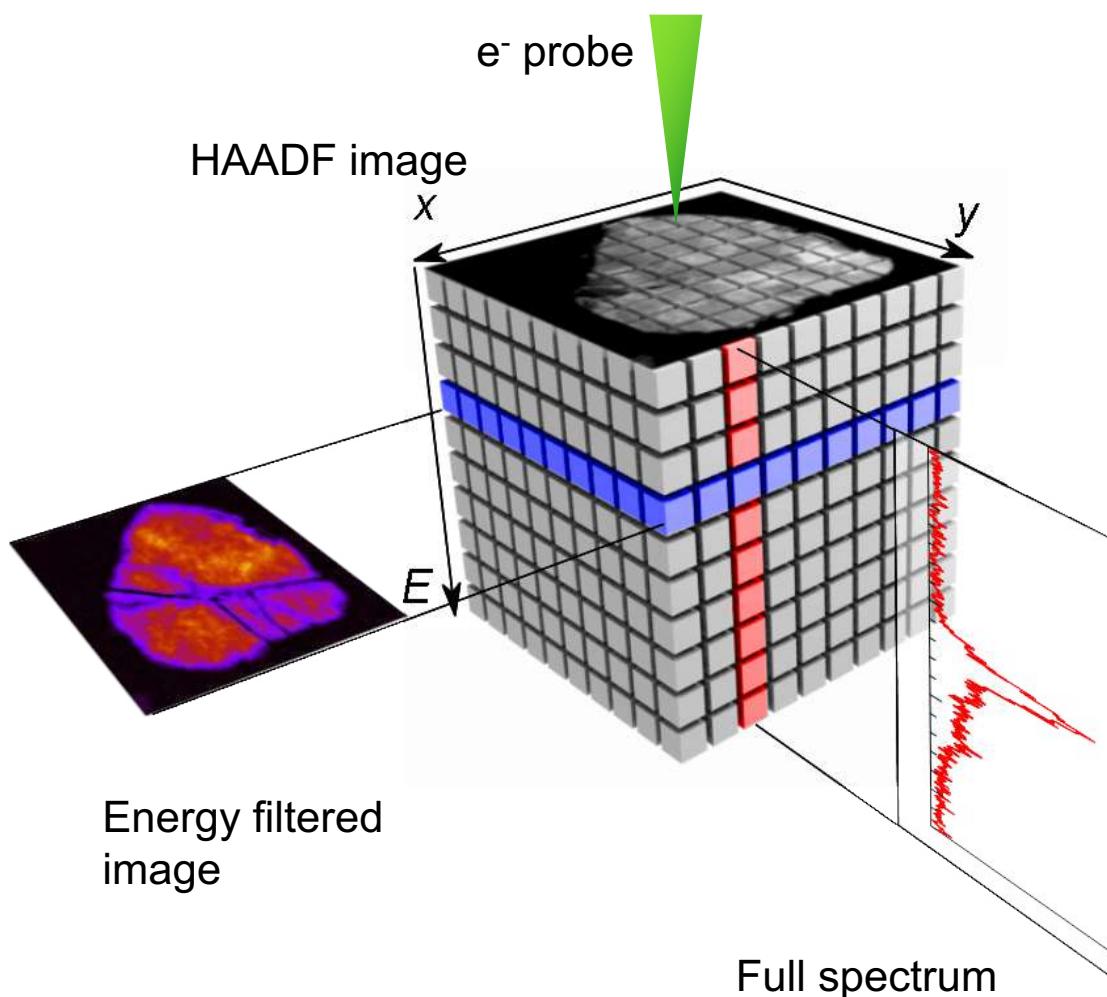


The multisignal approach

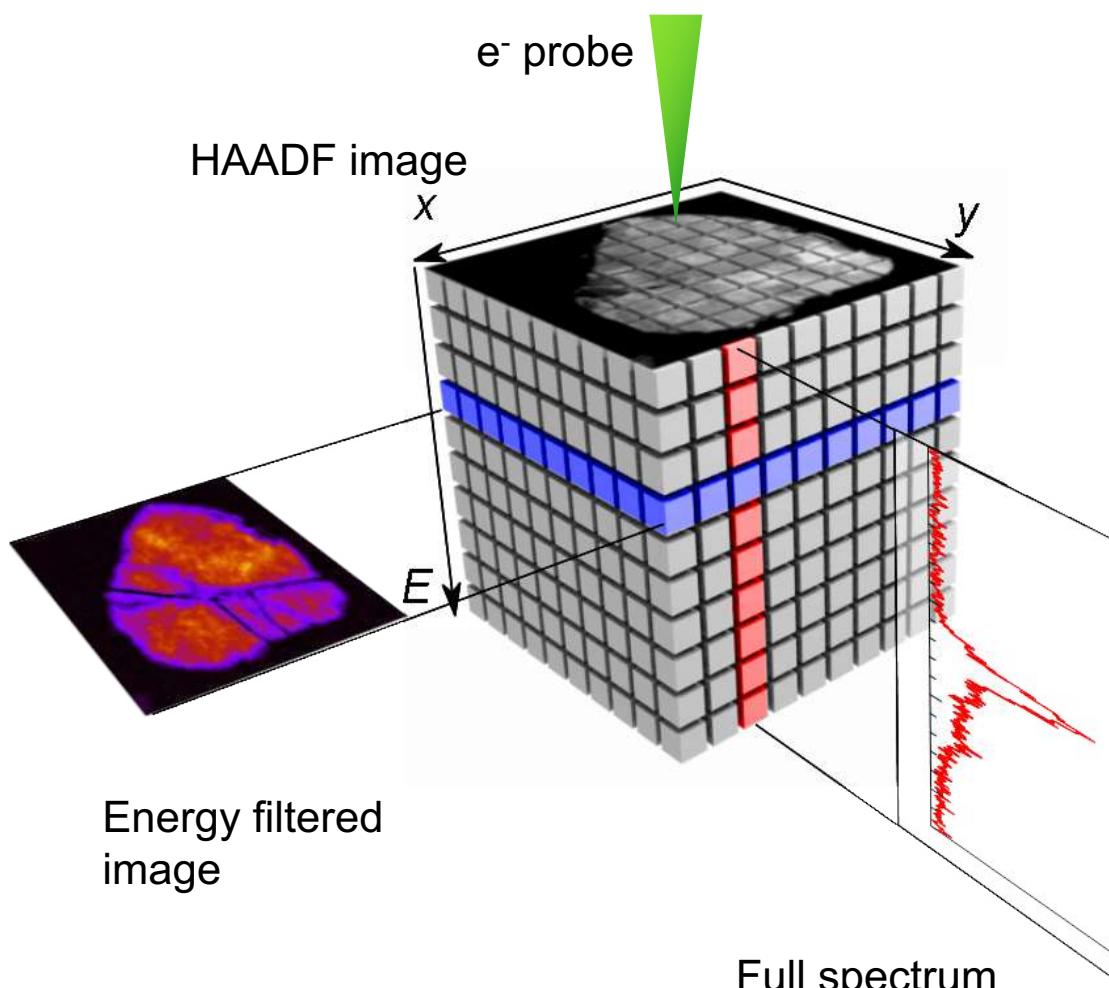
EELS spectroscopy : spectral domains



Hyperspectral imaging



Hyperspectral imaging

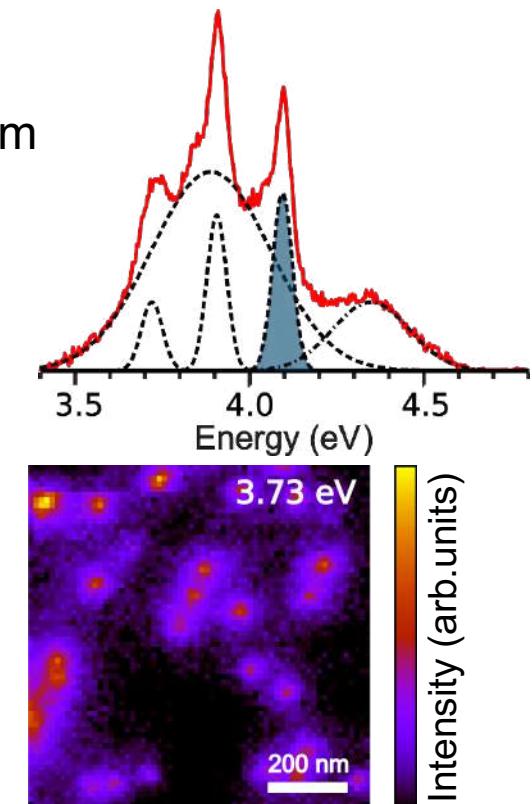


Energy filtered image

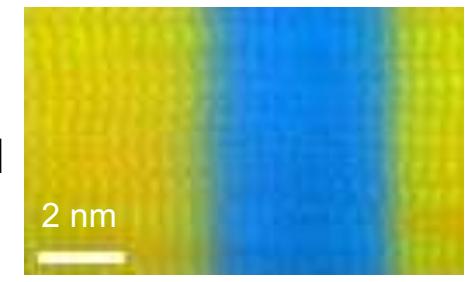
Full spectrum

CL spectrum

CL Intensity map

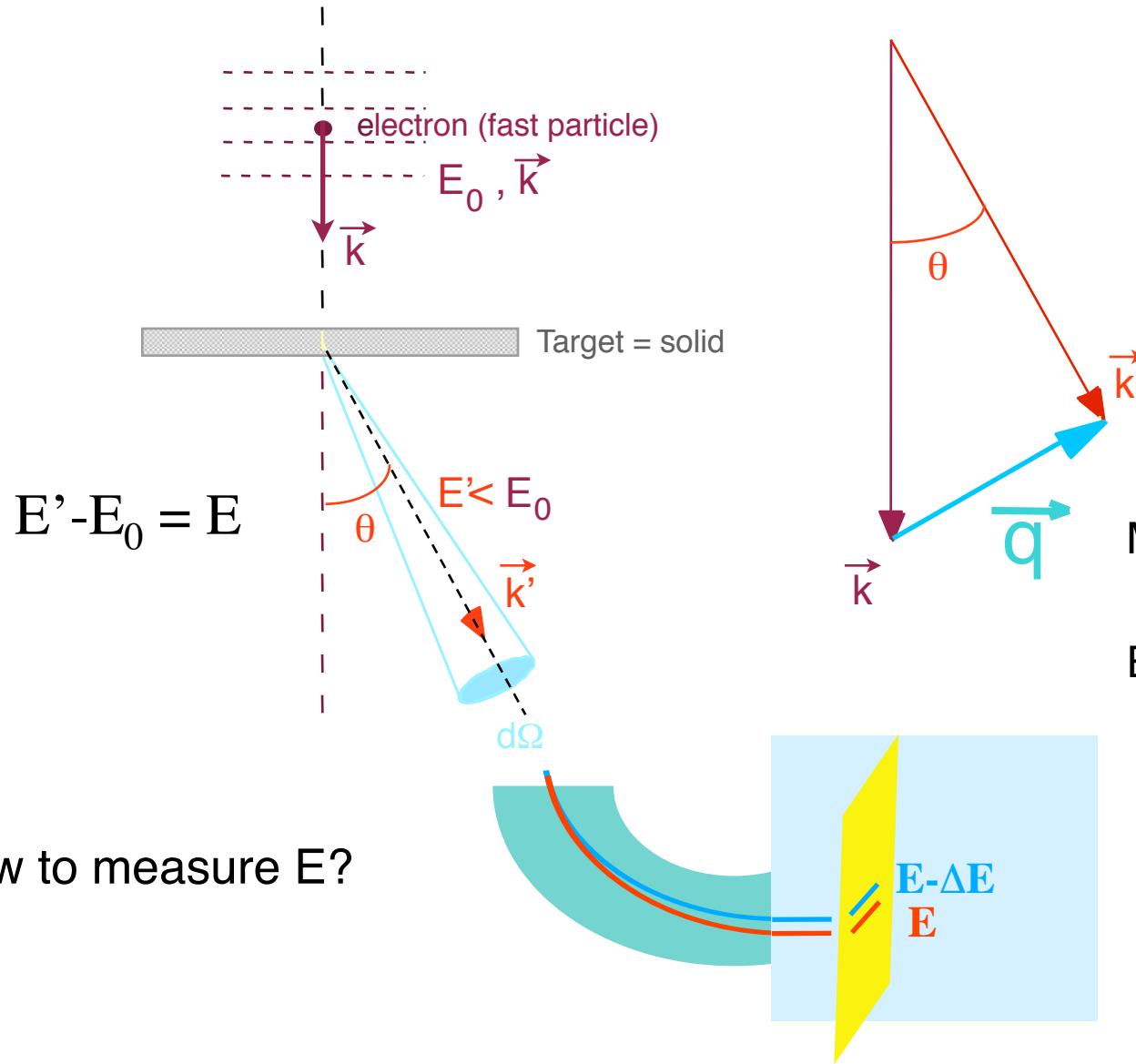


EELS Elemental map



Ti ■ Sr ■ La ■

EELS Scattering geometry



How to measure E?

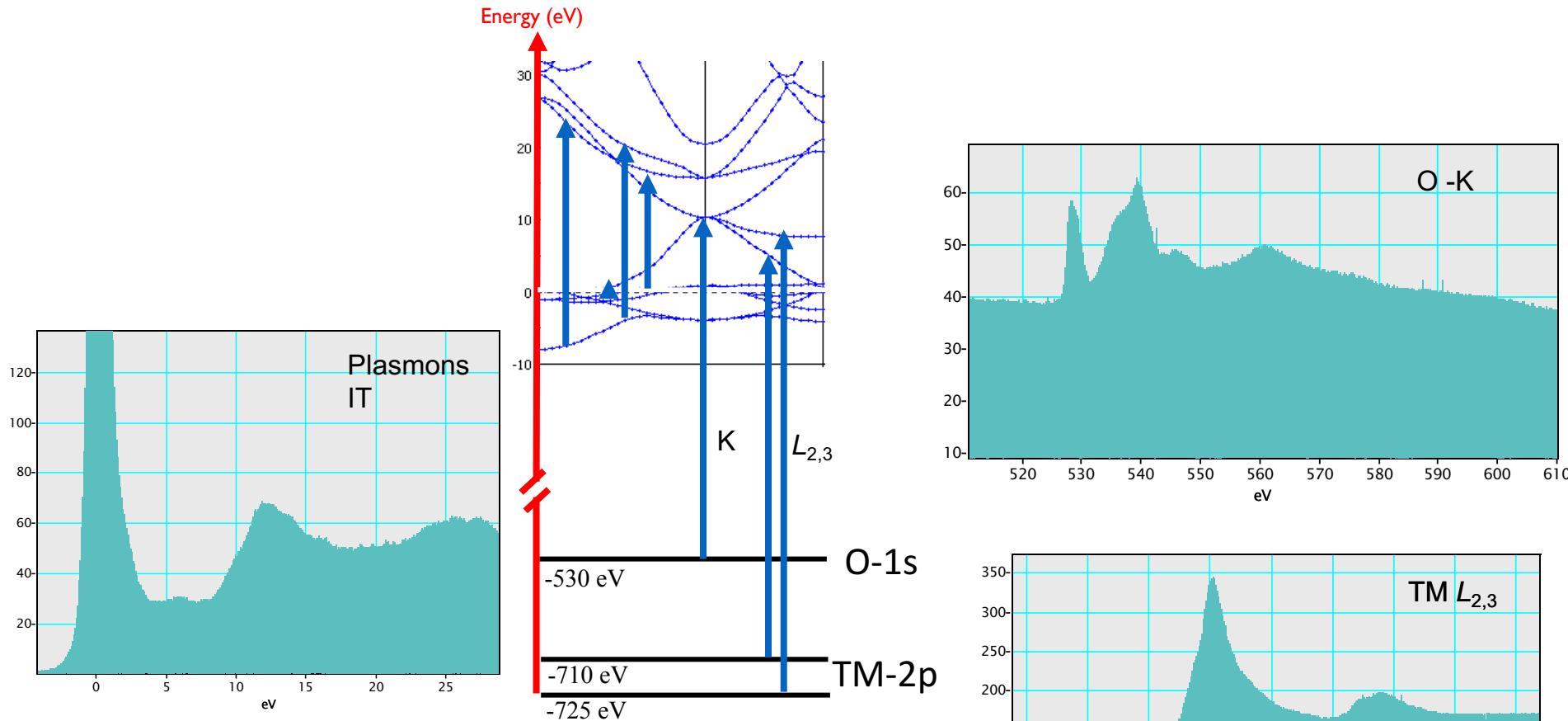
Measured quantities:

$$\frac{d^2S(E, q)}{dEd\Omega}$$

Momentum transfer \vec{q}

Energy loss E

EELS: involved electron populations and associated transitions



EELS gives informations on the elemental composition of a material (down to the atomic scale)

EELS gives informations on the electronic structure but it is usually not a direct probe of the Ground State

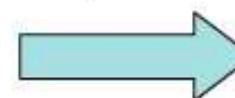
Comparison between the XAS and EEELS x-sections

$$\sigma(\hbar\omega) = 4\pi^2 \epsilon \hbar \omega \sum_F \left| \langle \Psi_F | \sum_i \vec{\epsilon} \cdot \vec{r}_i | \Psi_I \rangle \right|^2 \delta(E_i - E_f - \hbar\omega)$$

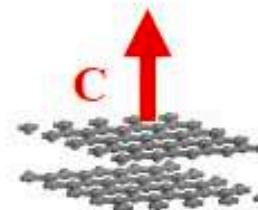
where \mathbf{e} is the polarisation vector of the electromagnetic field

$$\frac{\partial^2 \sigma}{\partial E \partial \Omega} = \frac{4\gamma^2}{a_0^2 q^4} \frac{k_f}{k_i} \sum_F \left| \langle \Psi_F | \sum_i \vec{q} \cdot \vec{r}_i | \Psi_I \rangle \right|^2 \delta(E_i - E_f - E)$$

Where $\mathbf{q} = \mathbf{k}_f - \mathbf{k}_i$ is the transferred momentum called the diffusion vector or the scattering vector



\mathbf{q} and ϵ may play a similar role



In case of an uniaxial symmetry

$$1s \rightarrow \pi^*$$

If \mathbf{q} or ϵ are // to the z axis, only $\Delta m = 0$ that means for a transition from a 1s core state, only the transition to a p_z state are allowed.

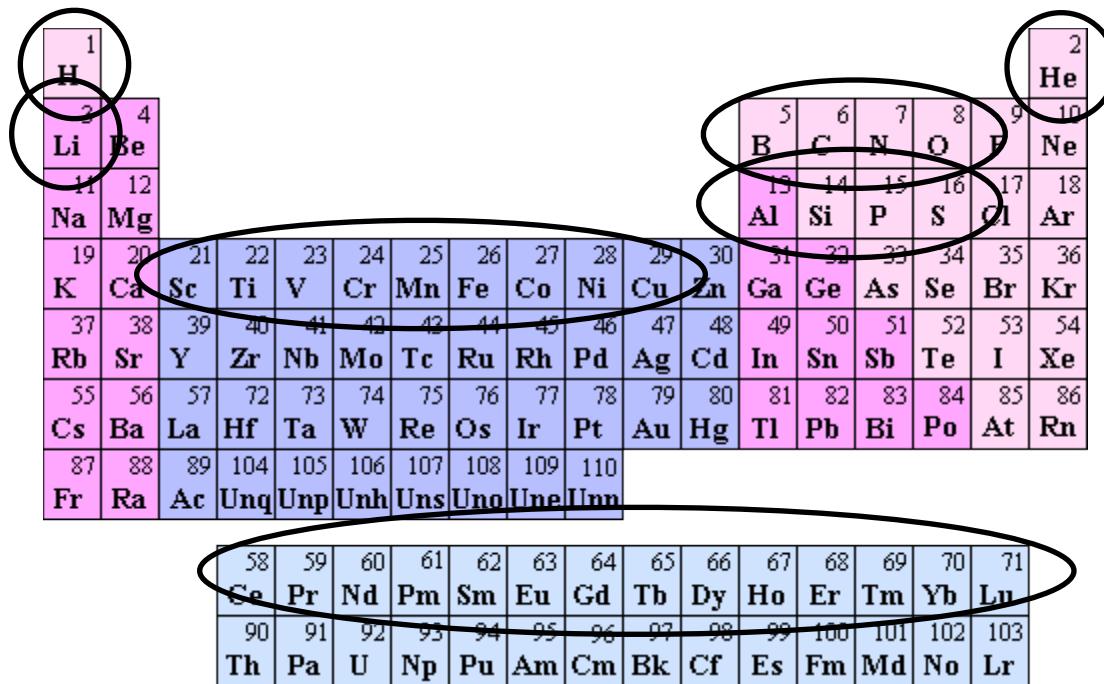
If \mathbf{q} or ϵ are within the x,y plane, $\Delta m = +1, -1$ then transition to state with σ^* hybridization are allowed for graphite .

$$1s \rightarrow \sigma^*$$

$I(\Delta m = 0) - I(\Delta m = +1, -1)$ is usually called the Linear Dichroic signal

Caution: just like in XAS, for TM, this picture is not truly valid
There is a trade off between momentum and spatial resolution

Absorption edges accessible with EELS



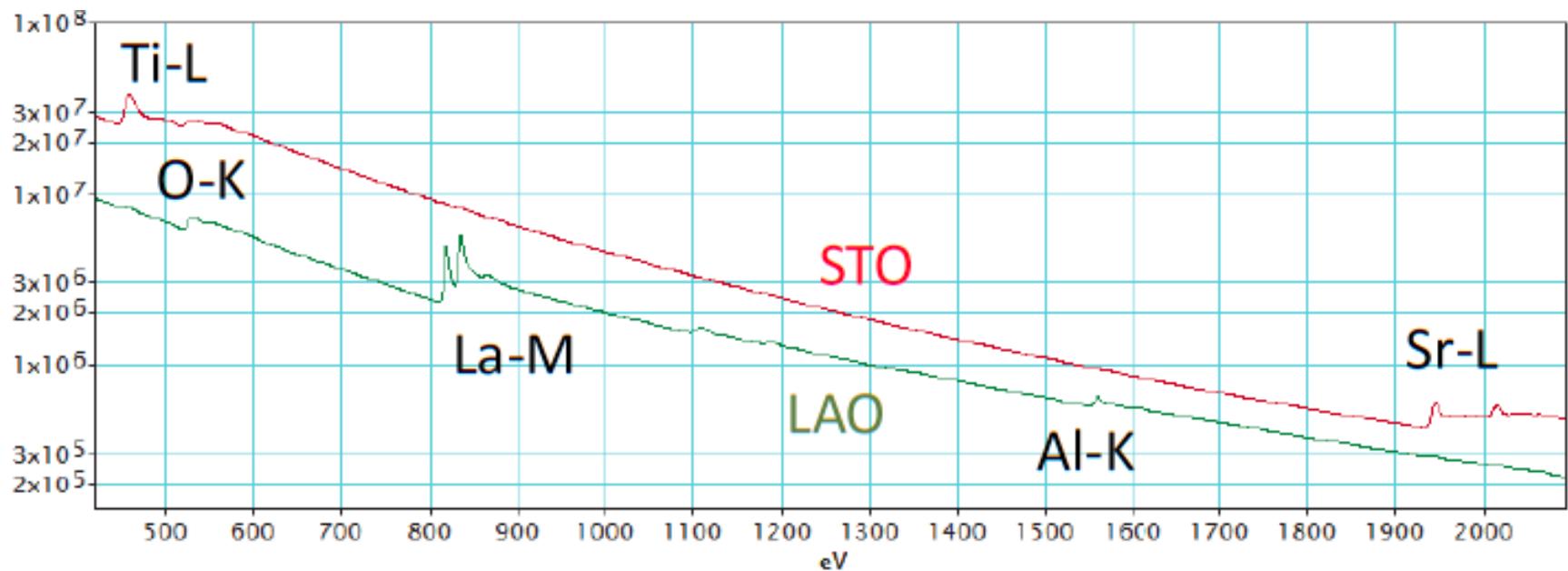
H 1s	13 eV
He 1s	24 eV
Li 1s	55 eV
B,C,N,O 1s	180 eV -- 530 eV
Mg, Al, Si 1s	1300 eV -- 1900 eV
P, S, Cl 1s	2100 eV -- 2900 eV
Mg, Al, Si 2p	50 eV -- 100 eV
P, S, Cl 2p	130 eV -- 200 eV
TM 1s	3900 eV -- 9000 eV
TM 2p	350 eV -- 900 eV
TM 3p	40 eV -- 80 eV

- X-section , - Probe intensity, - Damage...

EELS domain is usually limited from 1 to 2000 eV....

EELS characteristic core-level signals in STO/LAO

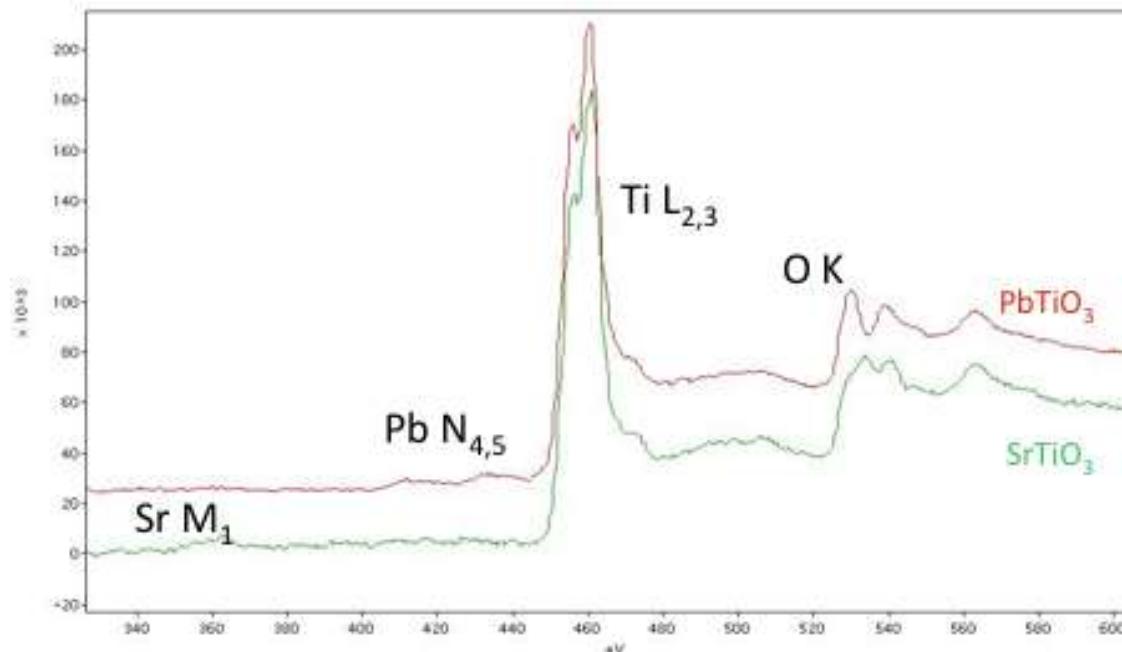
Log scale



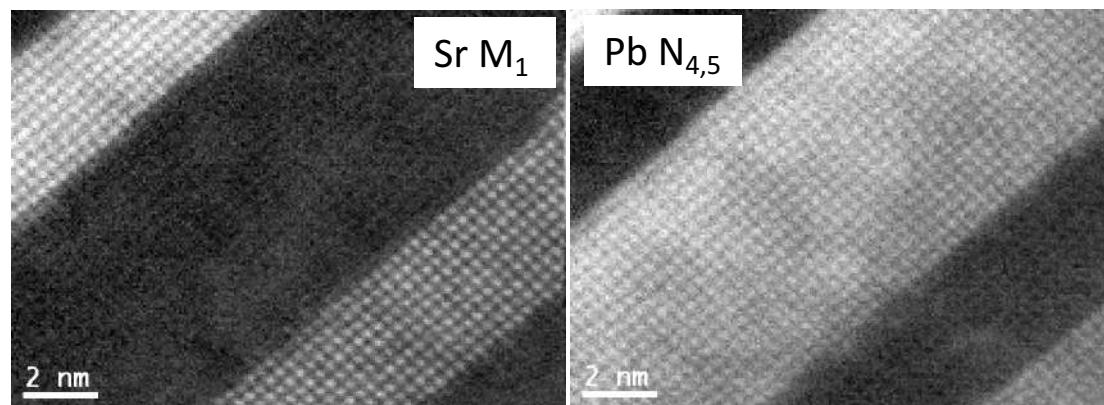
Courtesy G. Tieri & A. Gloter (CNRS-Orsay-France);
sample JM. Triscone (U. Geneva).

Large intensity dynamics and spectral range

Core-loss EELS
Elemental
analysis of
PTO/STO
superlattices



Mapping the
intensity of the
characteristic
edges



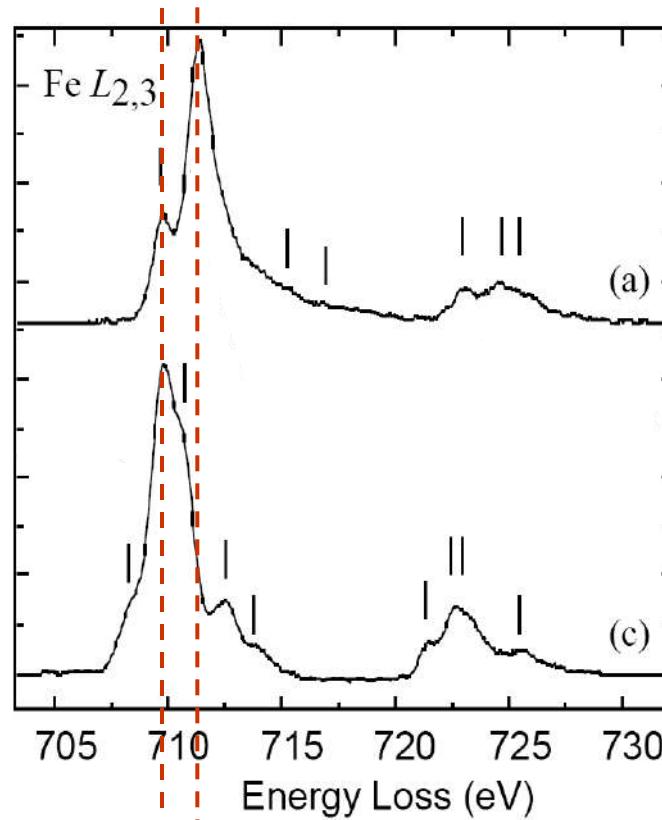
Negligible Sr-Pb interdiffusion in $\text{SrTiO}_3\text{-PbTiO}_3$ superlattices as evidenced by the EELS Sr M_1 (346-370 eV) and $\text{Pb N}_{4,5}$ (403-440 eV) signals.

STEM-HAADF, pixels size 9pm, 200 keV, 60pA, spots at 0.86 angstrom are seen in FFT.
STEM-EELS, pixel size 64pm, 5 ms per spectrum, 625 meV/ch, 200 keV, 60pA.

Courtesy K. March, A. Gloter, M. Kociak, M. Tence, A. Torres-Pardo (CNRS-Orsay-France); sample JM. Triscone (U. Geneva).

EELS at the core-loss is sensitive to TM charge (now not so far from XAS)

Fe L_{2,3} edges
 $2p^63d^n \rightarrow 2p^53d^{n+1}$



$Fe^{3+} (d^5)$
 Fe_2O_3 hematite

$Fe^{2+} (d^6)$
 $FeCO_3$ siderite

Gloter et al., Ultramicroscopy
96, 385, (2003)

Ferrous-ferric iron

Chemical shift of L3 line maximum of 1.6 - 2 eV

TM 2p edges (TiL_{2,3}) : Origin of the peaks

Classic scheme for 3d orbitals crystal field splitting
and 2p / L_{2,3} excitation

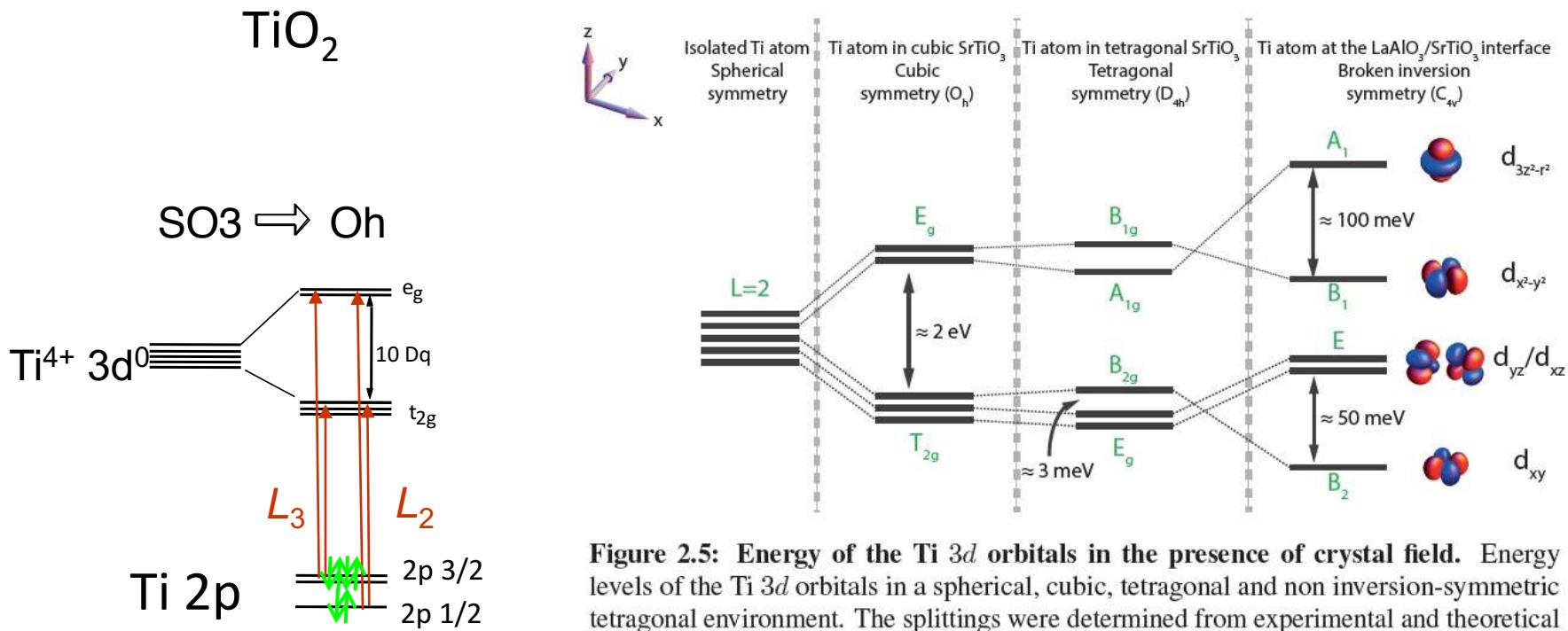
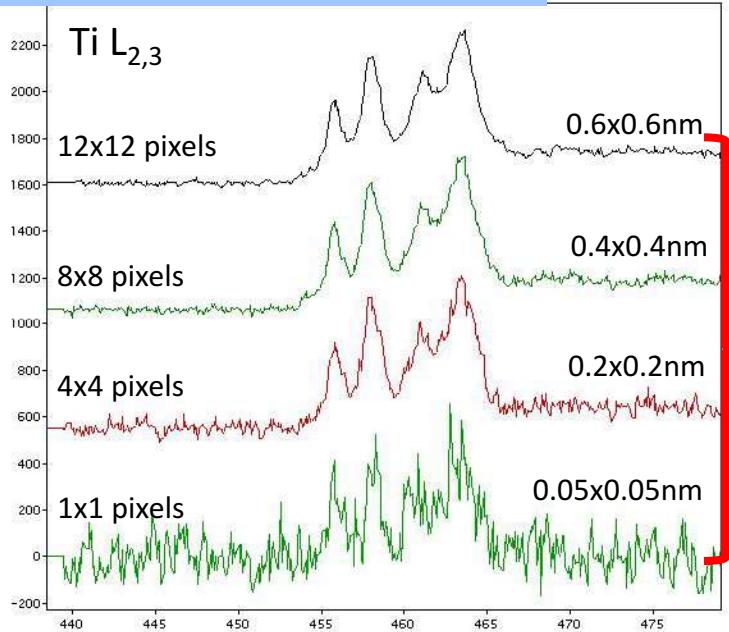


Figure 2.5: Energy of the Ti 3d orbitals in the presence of crystal field. Energy levels of the Ti 3d orbitals in a spherical, cubic, tetragonal and non inversion-symmetric tetragonal environment. The splittings were determined from experimental and theoretical studies [50–55]. The green labels refer to the irreducible representation to which the orbital(s) belong(s).

EELS principles: fine structures mapping



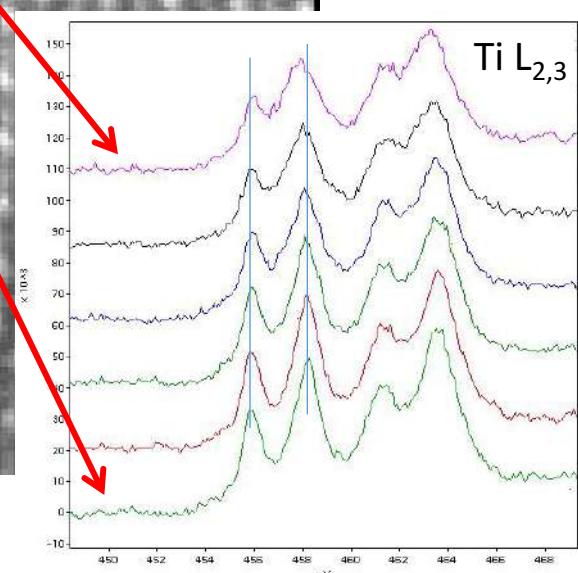
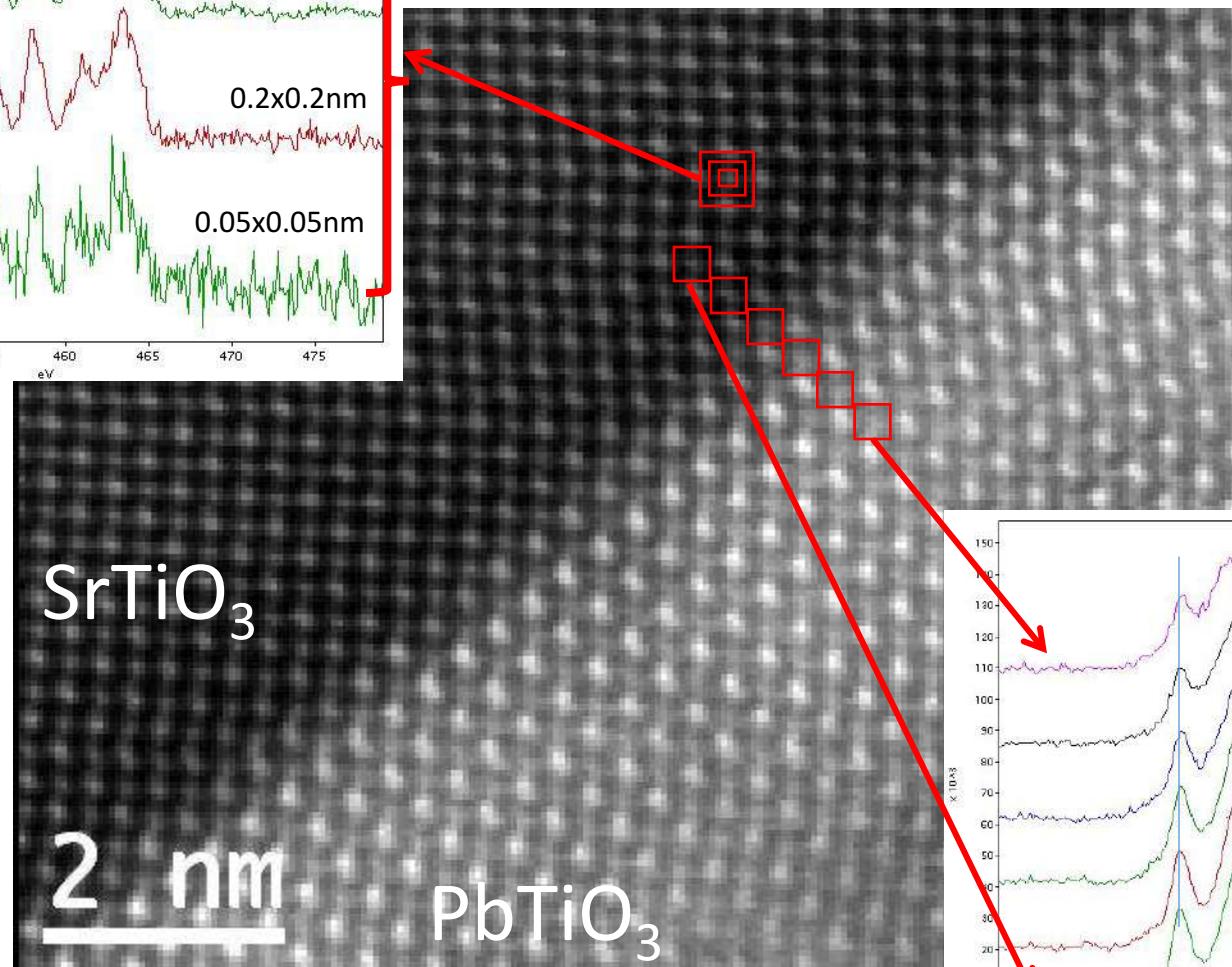
Measuring the crystal field splitting across a $\text{SrTiO}_3\text{-PbTiO}_3$ interface by the EELS $\text{Ti } L_{2,3}$ edges fine structures evolution.

206x161 spim images, 50pm pixel size, 10 ms per spectrum, 90 meV/ch, 200 keV, 60pA, raw EELS data.
Courtesy A. Gloter, M. Kociak, M. Tence, A. Torres-Pardo (CNRS-Orsay-France), sample JM. Triscone (U. Geneva).

A.Torres-Pardo et al., Phys. Rev. B 84, 220102 (2011).

P. Zubko et al. Nano Letters 12(6), 2846 (2012).

Coll. J.M. Triscone
U. Geneva



Spectromicroscopy up to 100.000 spectra

Unit cell resolved fine structures on $\text{Ti } L_{2,3}$ edges

Thin film : unusual (magnetic) properties

Take one magnetic
One non-magnetic

(La,Sr)MnO₃
(Ferro-magnet)

SrTiO₃ (non
magnetic)

Make heterostructures

(La,Sr)MnO₃
(10uc)

SrTiO₃
substrate

Still ferro-magnetic

Make smaller heterostructures

(La,Sr)MnO₃(4uc)

SrTiO₃ (7uc)

(La,Sr)MnO₃(4uc)

SrTiO₃ (7uc)

(La,Sr)MnO₃(4uc)

SrTiO₃ (7uc)

(La,Sr)MnO₃(4uc)

SrTiO₃ substrate

Non magnetic

Interfacial Control of Magnetic Properties at LaMnO₃/LaNiO₃ Interfaces

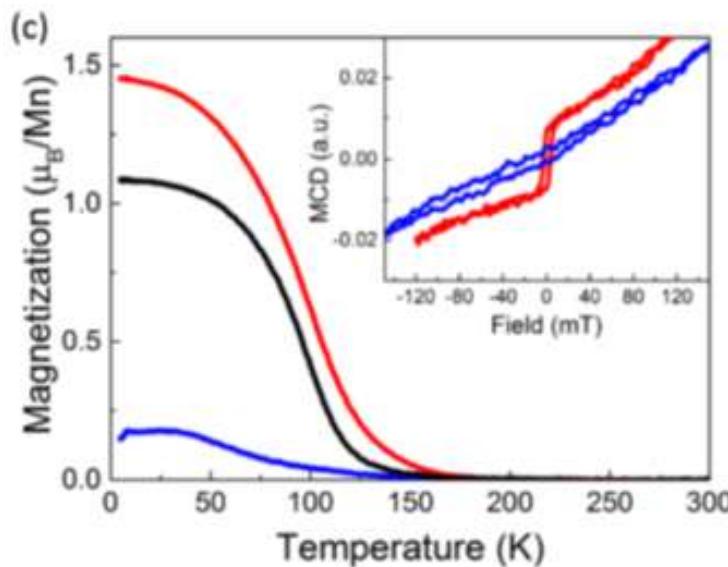
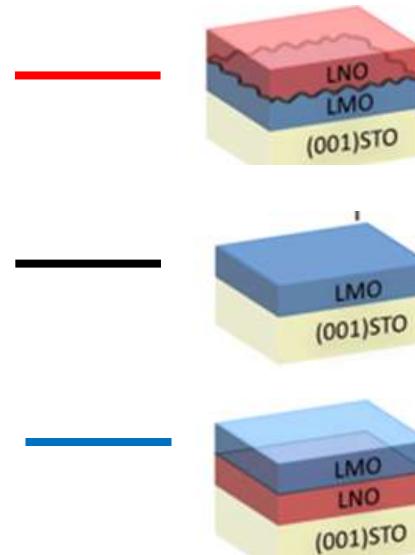
M. Gibert,^{*,†} M. Viret,^{†,‡} A. Torres-Pardo,[§] C. Piamonteze,^{||} P. Zubko,[†] N. Jaouen,[⊥] J.-M. Tonnerre,[#] A. Mougin,[§] J. Fowlie,[†] S. Catalano,[†] A. Gloter,[§] O. Stéphan,[§] and J.-M. Triscone[†]

Nano letters 15, 7355–7361 (2015)

LaMnO₃ is Mn³⁺, Mott insulator with A-type antiferromagnetism in bulk form, becomes ferromagnetic in thin film grown on STO.

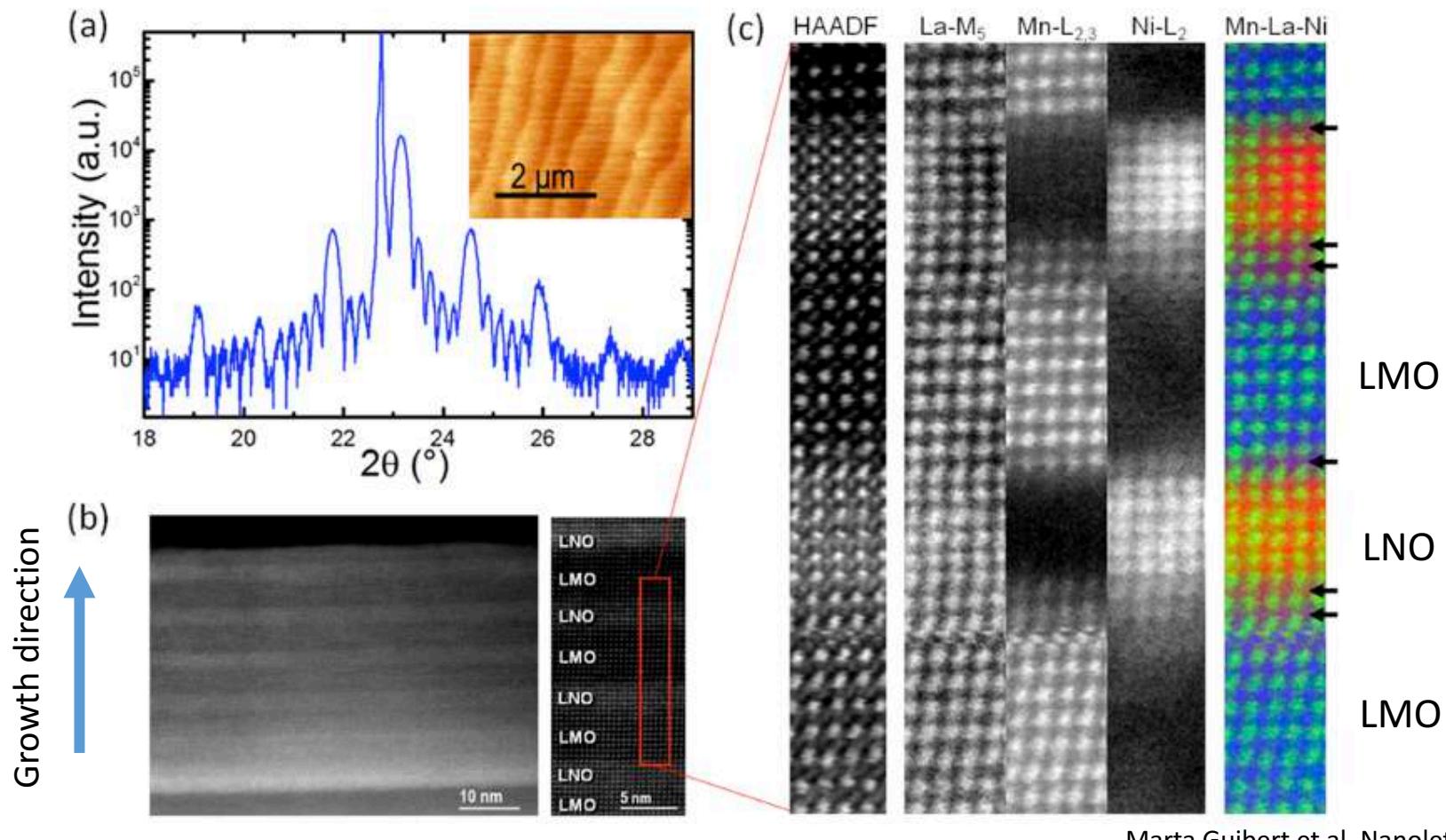
LaNiO₃ is Ni³⁺, paramagnetic metal in bulk form, could become insulating and/or AFM in thin films.

Magnetization versus temperature of
 $(7\text{u.c.LNO}/7\text{u.c.LMO})_1//(001)\text{STO}$ (red),
 $(7\text{u.c.LMO}/7\text{u.c.LNO})_1//(001)\text{STO}$ (blue) and
 $7\text{u.c.LNO}//(001)\text{STO}$ (grey).
A 28u.c. LMO film//(001)STO (black) is also shown.



LMO FM is decreased or enhanced in the bi-layers

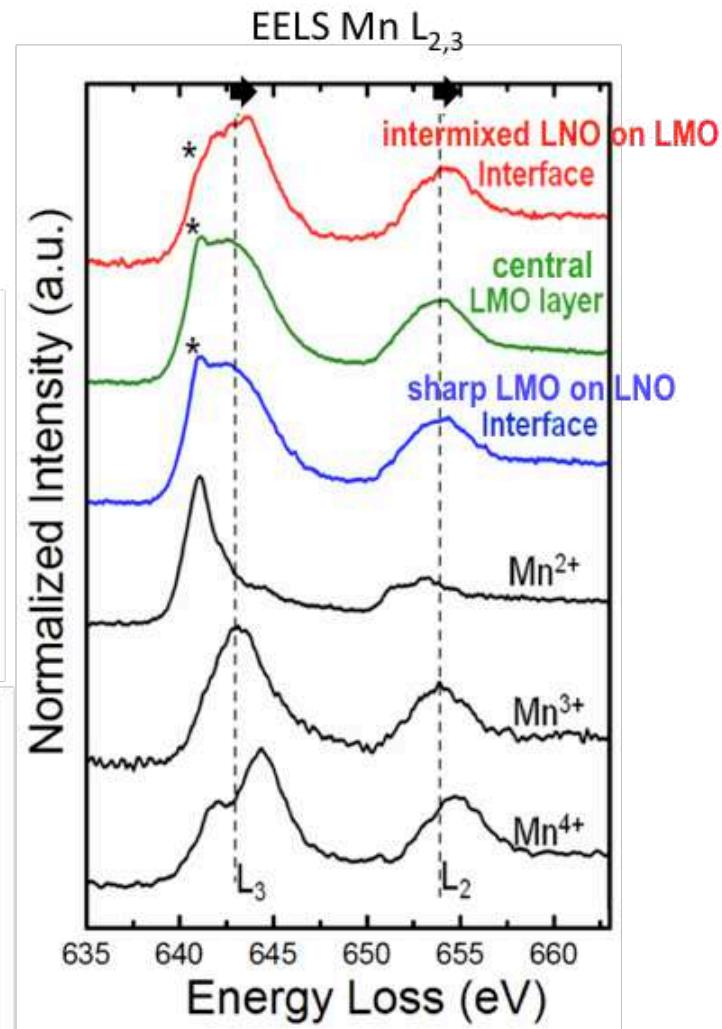
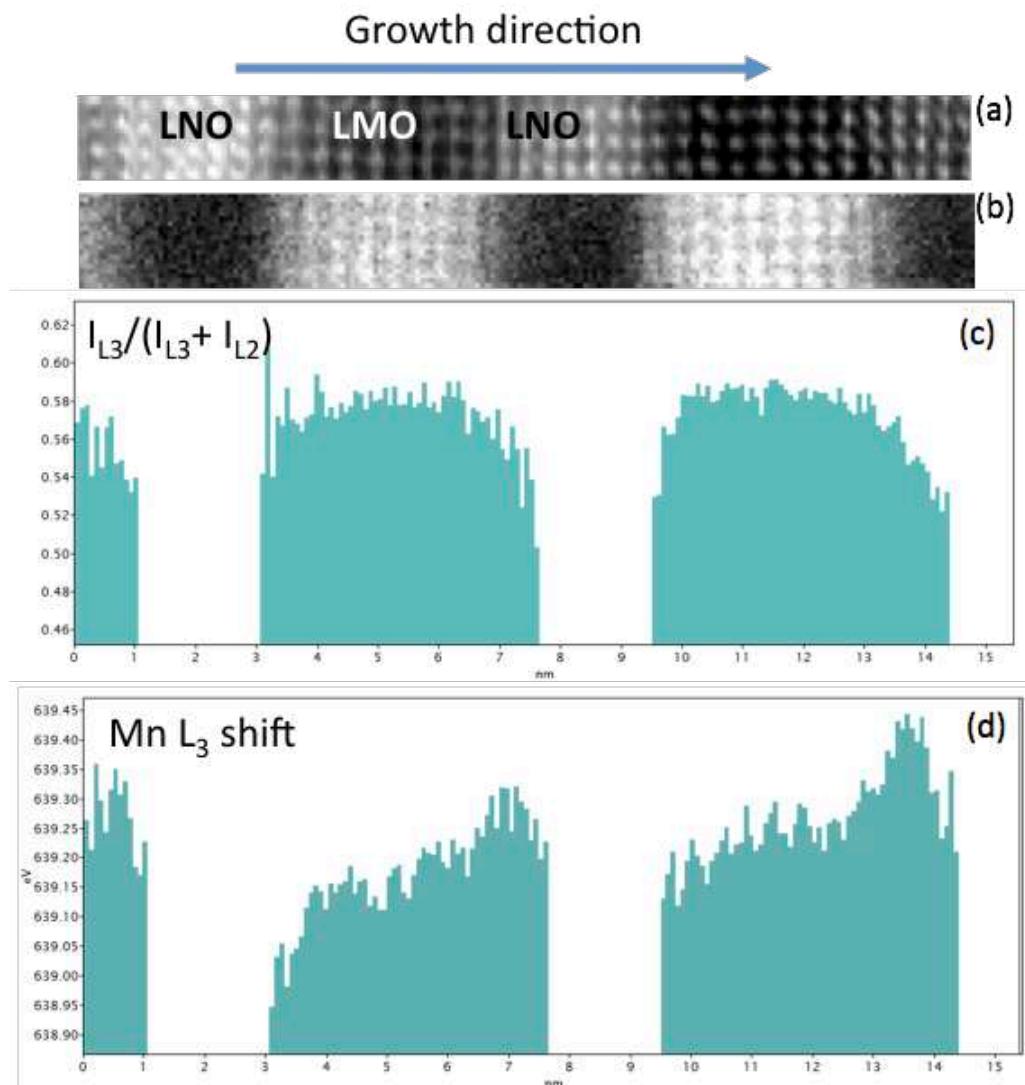
Structure of the LNO/LMO interfaces in superlattices



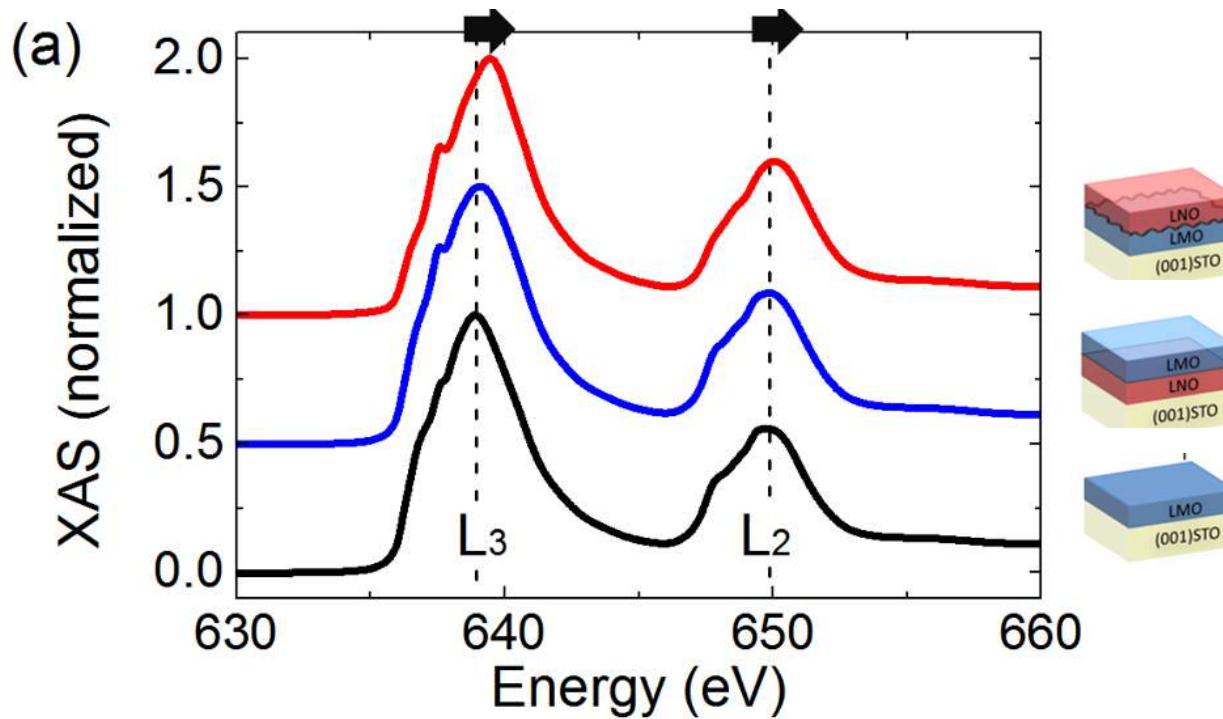
Marta Guibert et al. Nanoletter 2016

- (a) XRD scans for a (8LNO/8LMO)₆ superlattice grown on (001)STO. Inset: corresponding atomic force microscopy image.
- (b) Low magnification TEM image. (c) EELS measurements showing the interfacial structural asymmetry (indicated by arrows).

LMO/LNO interface is rougher compared to the LNO/LMO interface

EELS Mn L_{2,3} at the LNO/LMO interfaces

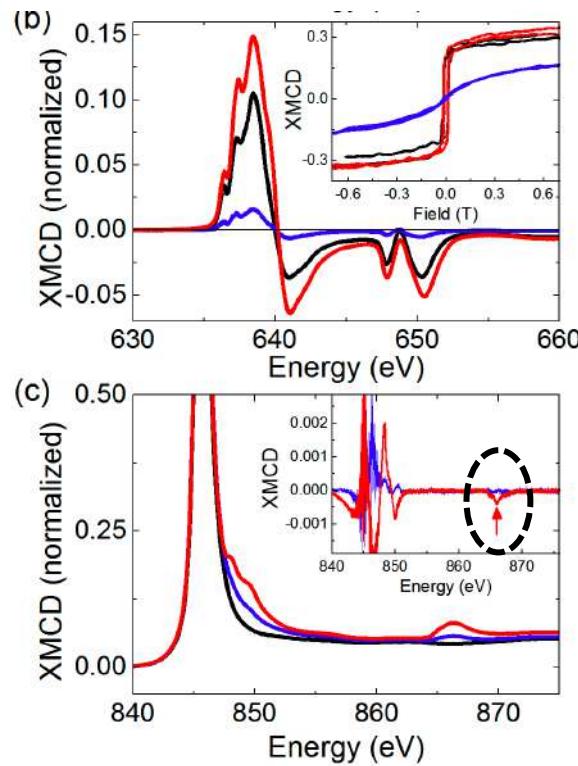
Both energy shifts and branching ratio indicated introduction of more holes (Mn⁴⁺) in the rougher LMO/LNO interface compared to the LNO/LMO interface

XAS Mn L_{2,3} of the STO/LMO , STO/LNO/LMO and STO/LMO/LNO

XAS measurements at 2K in 30° grazing incidence in total electron yield (TEY) configuration after field cooling in 0.05T with the field parallel to the sample plan. (C. Piamonetze, PSI, Suisse)

Electron transfer from Mn³⁺-Ni³⁺ to Ni²⁺-Mn⁴⁺ for STO/LMO/LNO

XAS spectral anisotropies (XMCD, XLD)



X-ray Magnetic Circular Dichroism (XMCD) shows magnetic Mn (and Ni for the rough interface)

X-ray Linear Dichroism (XLD) is also observed,

- XLD_{300K} at room temperature evidence orbital anisotropies
- (XLD_{2K_1T} - XLD_{300K}) and (XLD_{2K_0T}-XLD_{300K}) provide information about FM and AFM ordering

EELS – XAS – XMCD – XLD indicate that

1) (STO [100])-LMO-LNO is rough with Mn⁴⁺-Ni²⁺ charge distribution at the interface.

The Ni²⁺- Mn⁴⁺ interact ferromagnetically according to Goodenough-Kanamori rules, like in the case of the insulating ferromagnetic double-perovskite La₂NiMnO₆

2) (STO [100])-LNO-LMO is sharp, with Mn³⁺-Ni³⁺ charge distribution at the interface.

Presence of a canted-antiferromagnetic order in the LMO layer

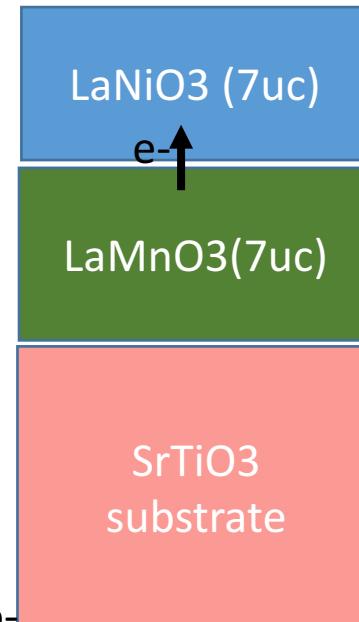
Thin film : unusual (magnetic) properties

Take basically
2 « non » ferro-magnetic

Make heterostructures

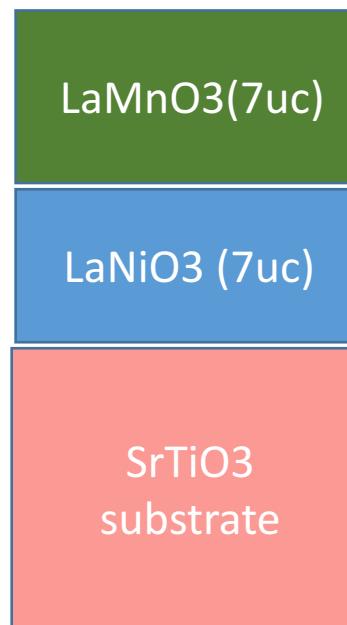


Becomes more ferro-magnetic



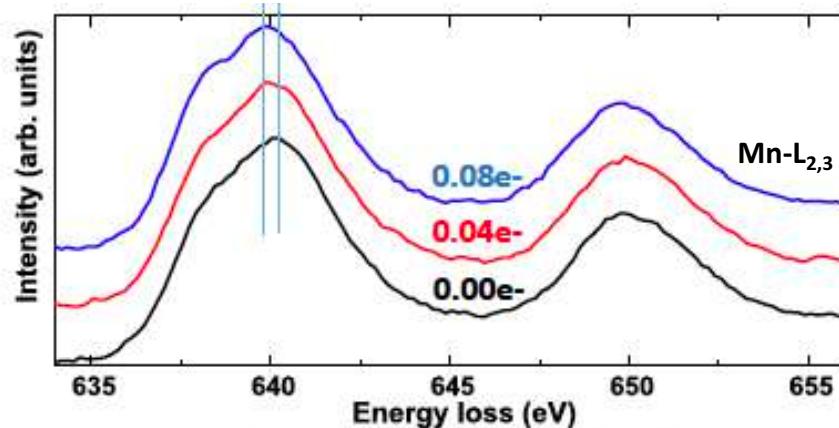
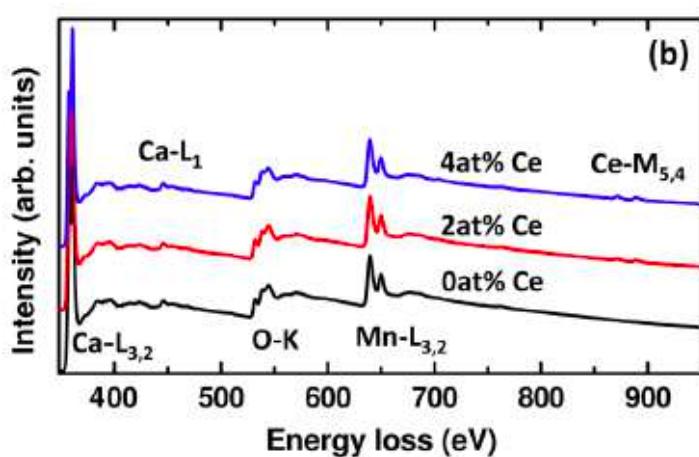
Mn⁴⁺-Ni²⁺
At the diffused
interface

Becomes non
magnetic

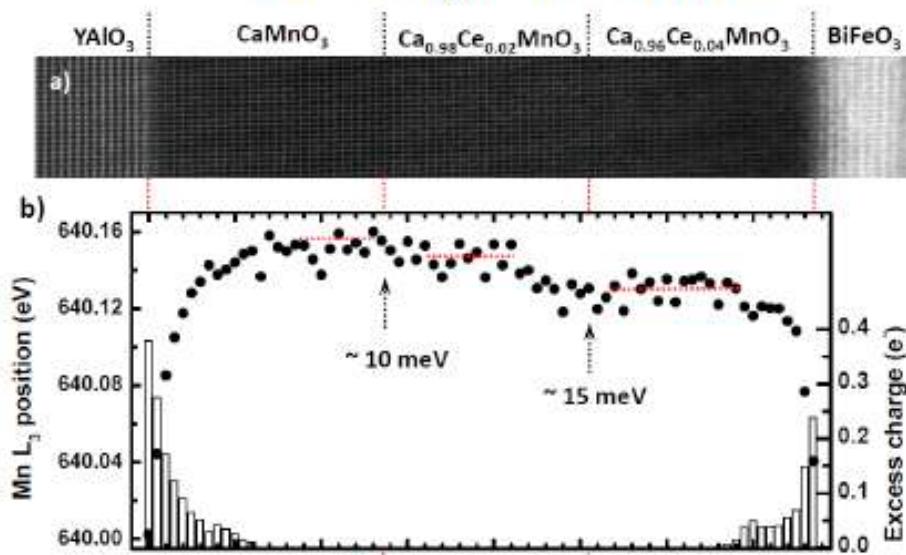
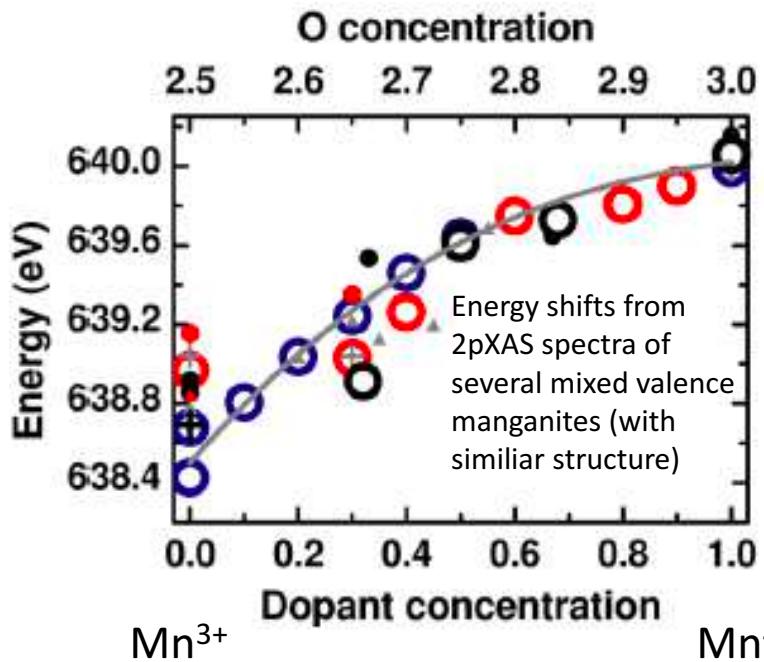


Stays
Mn³⁺-Ni³⁺
at the interface

Charge quantification at interfaces, which sensitivity?



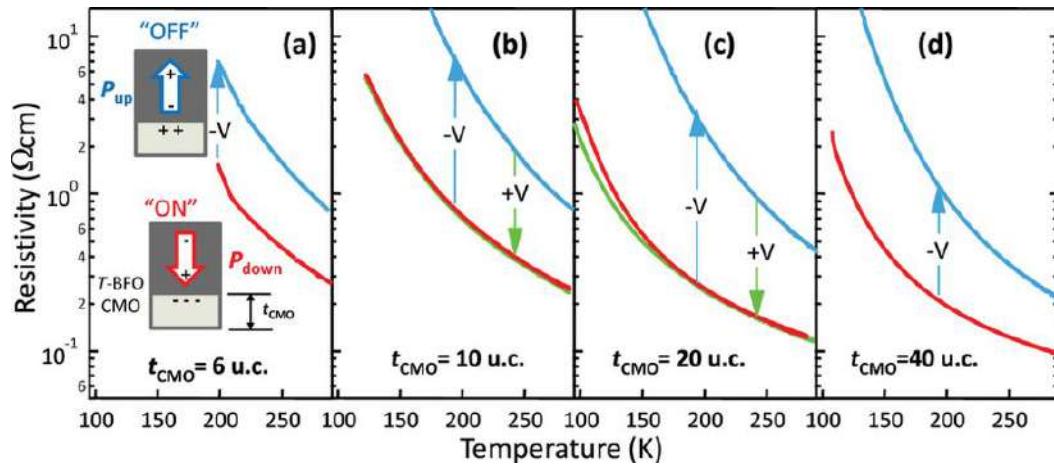
Excess charges at the interfaces



Ferroelectric control of a Mott insulator the $\text{Ca}_{1-x}\text{Ce}_x\text{MnO}_3/\text{BiFeO}_3$ system



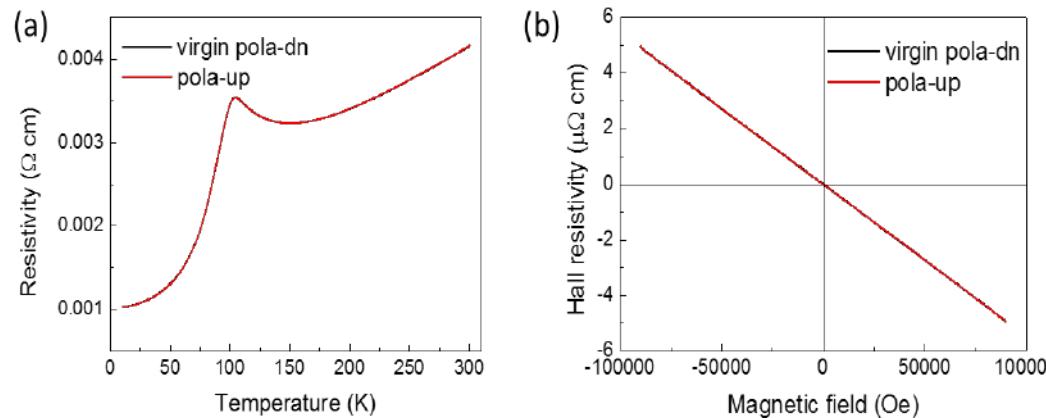
4 nm BFO



H. Yamada et al.
Scientific Reports 3,
2834 (2013).

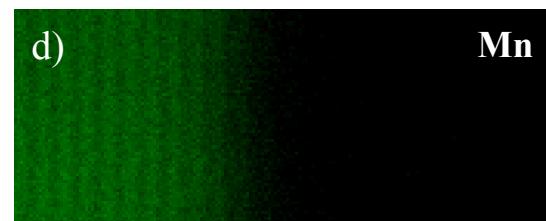
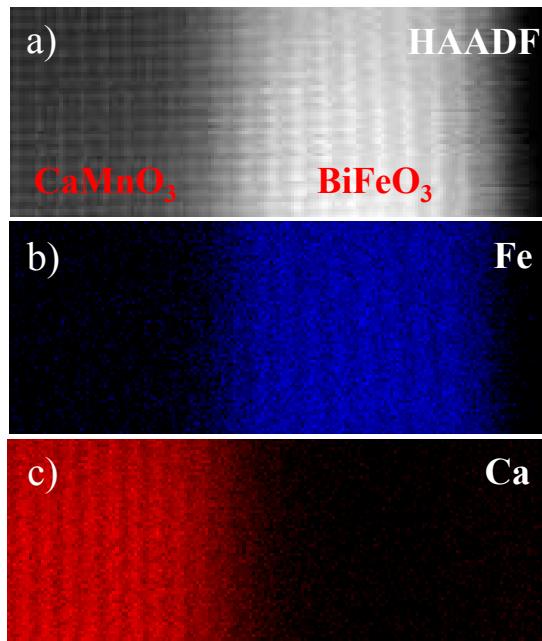
M. Marinova et al.
Nano letters 15 2533
(2015).

30 nm BFO

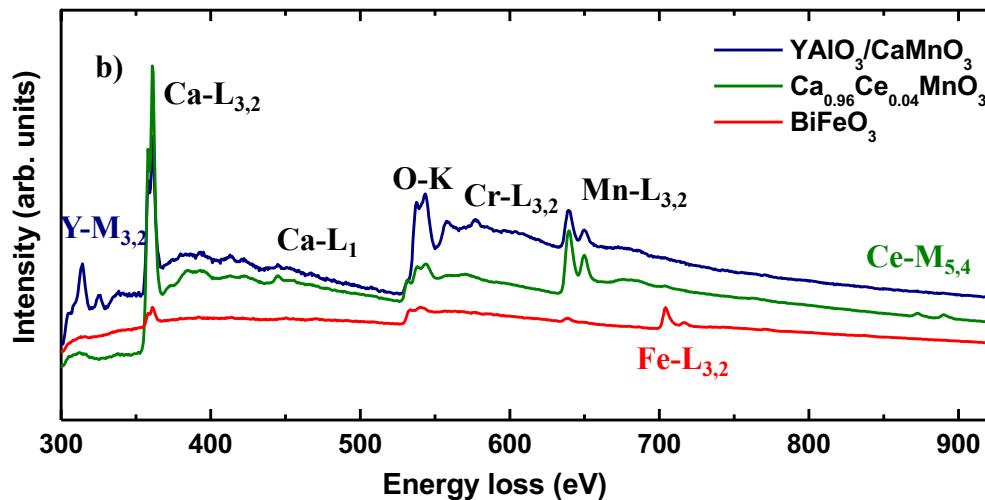


X. Li, in preparation.
No switching between
different polarisation states,
why?

Coll. V. Garcia et al.,
UMPhy Thales

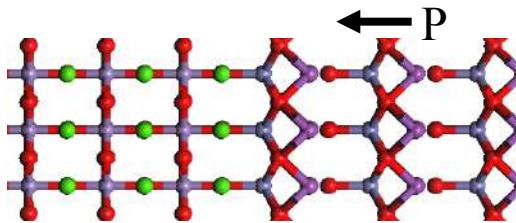


Interface termination planes ($\text{CaMnO}_3/\text{BiFeO}_3$)



Perovskite ABO_3 : AO plane and BO_2 plane

**The observed termination plane sequence is
 $(\text{CaO})^0-(\text{MnO}_2)^0-(\text{CaO})^0-(\text{FeO}_2)^{-1}-(\text{BiO})^{+1}-(\text{FeO}_2)^{-1}$**



and not $(\text{MnO}_2)^0-(\text{CaO})^0-(\text{MnO}_2)^0-(\text{BiO})^{+1}-(\text{FeO}_2)^{-1}-(\text{BiO})^{+1}$

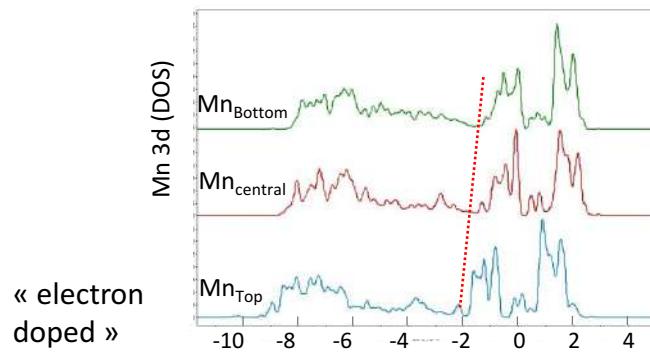
M. Marinova, et al. Nanoletters 2015.

Coll. V. Garcia, S. Fusil, A. Barthélémy, M. Bibes, UMPHY Thales

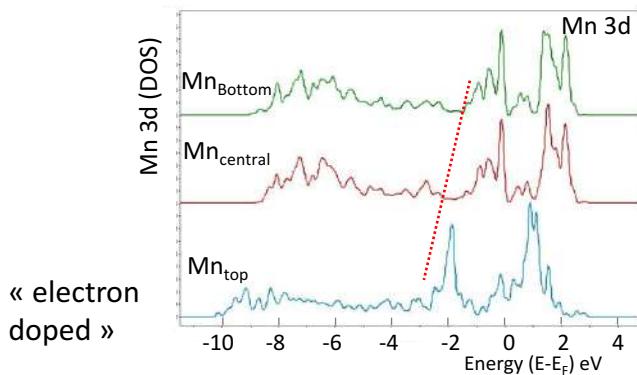
Ab-initio modeling and role of the interface plane termination

P point towards a...

.....Fe terminated-BFO



.....Bi terminated-BFO



Fe terminated is 1.5 eV lower in total energy
 It has less electrons in Mn eg orbitals
 It has larger BFO FE structural distortion

Polarization pointing toward CMO injects e-

« electron doped »

Bi-terminated

BFO

P

Fe-terminated

CMO

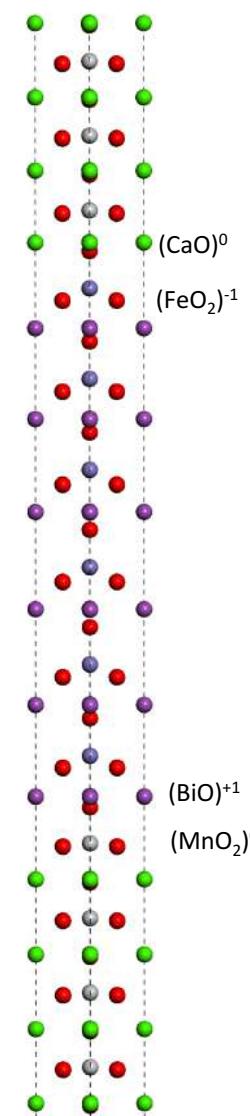
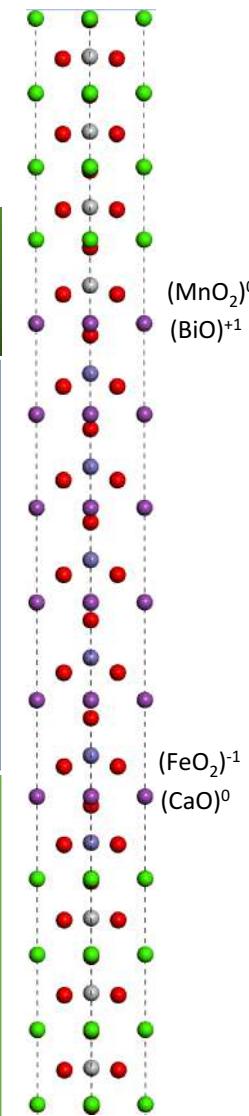
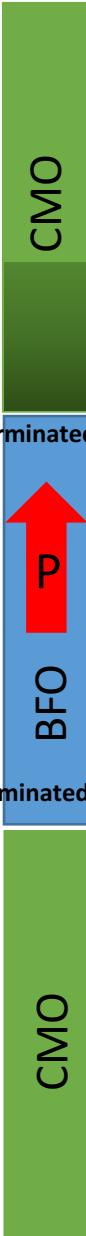
Fe-terminated

BFO

P

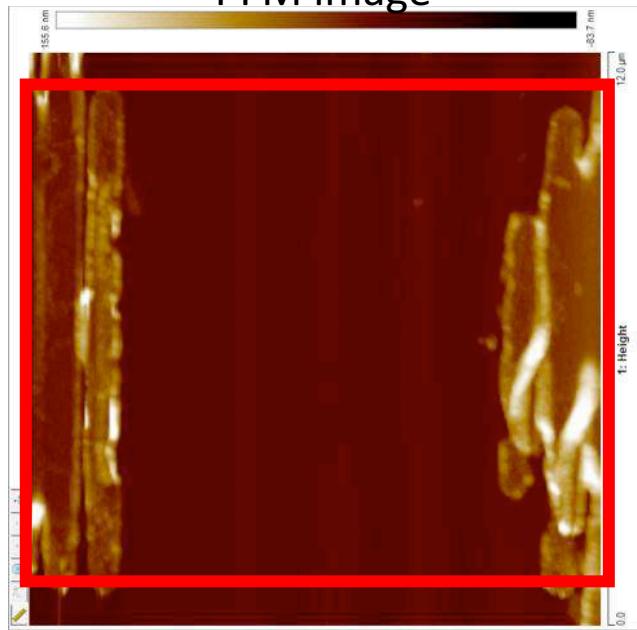
Bi-terminated

CMO

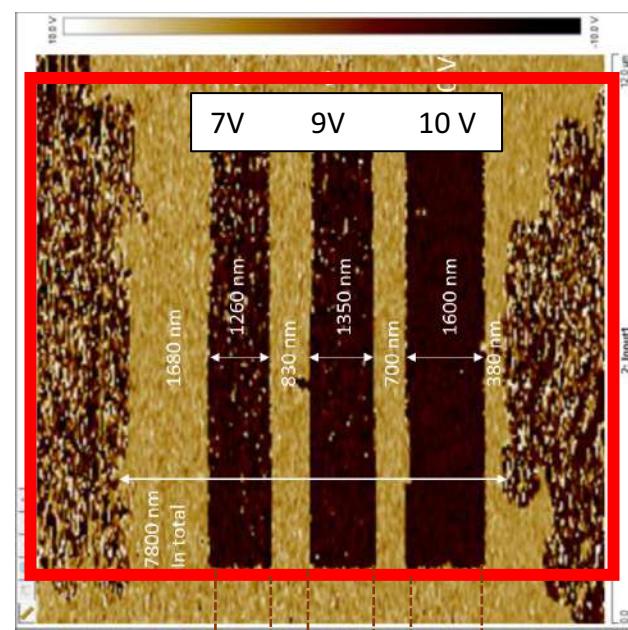


Ex-situ polling by electrode or piezo-force microscope + FIB sample preparation

PFM image

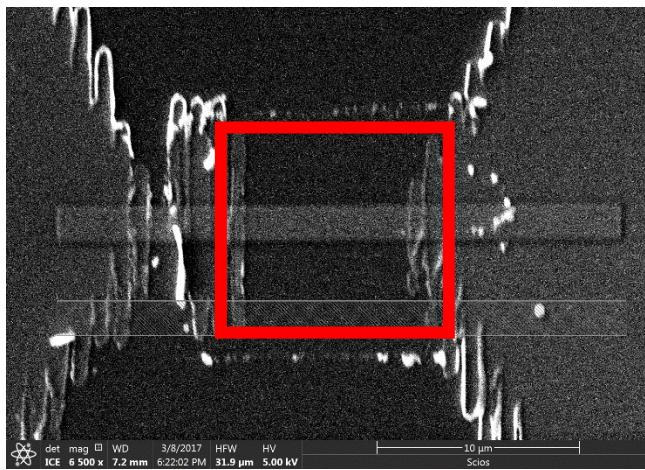


PFM phase image

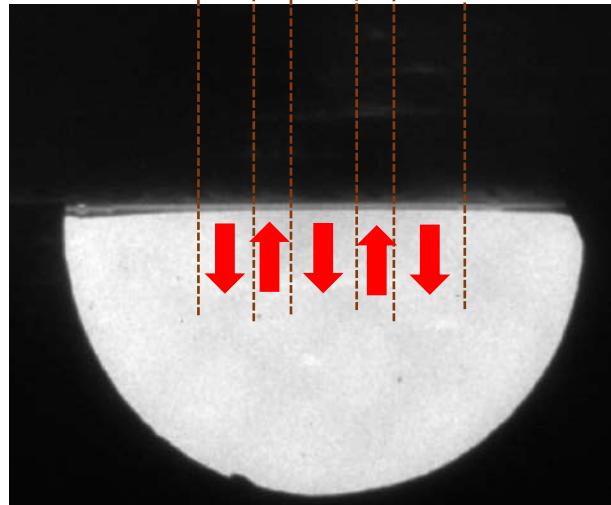


Apparent
electrical
switching in
PFM

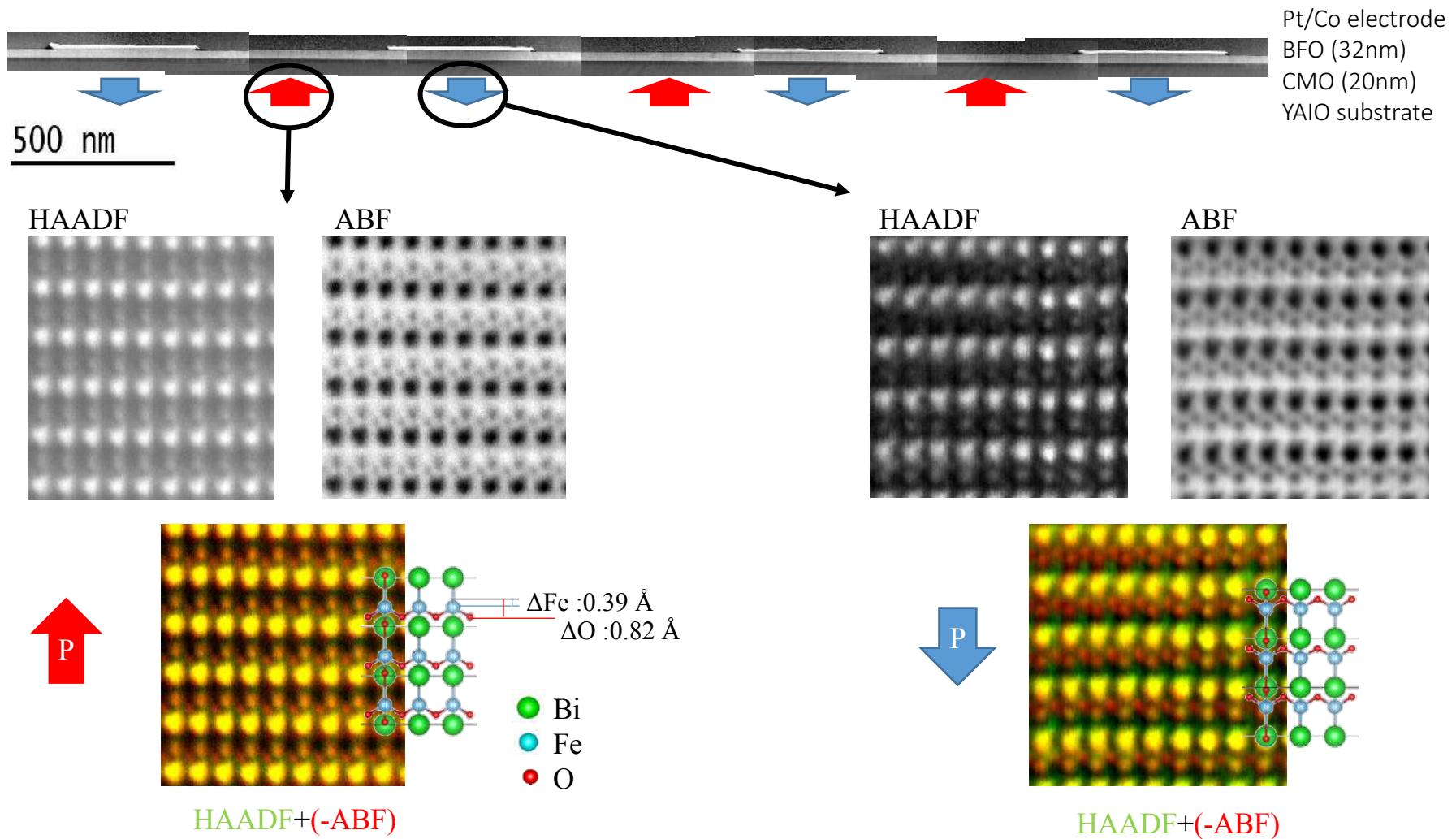
SEM image



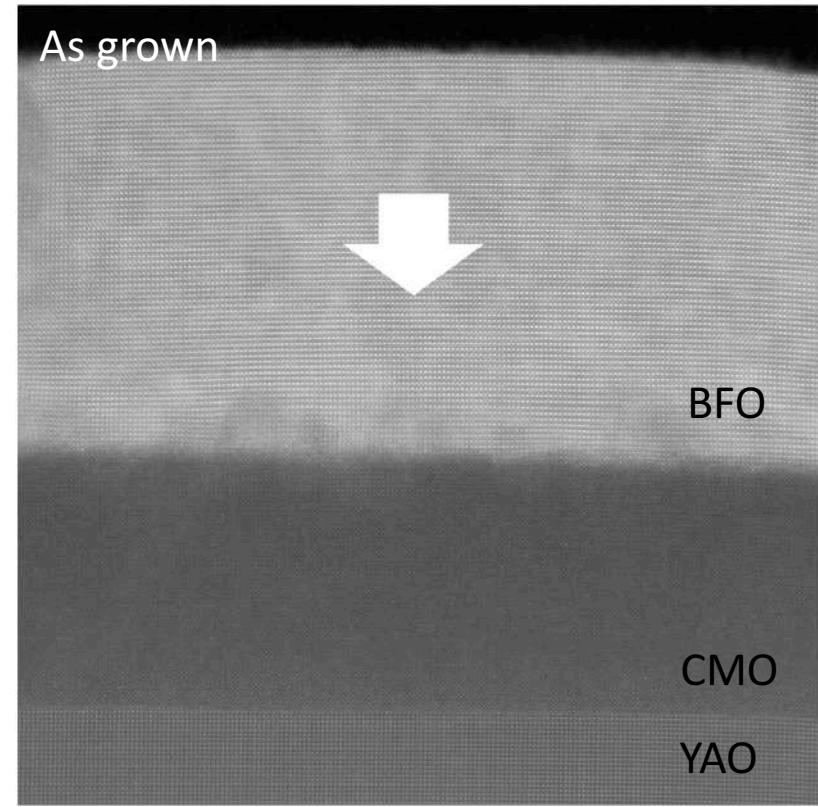
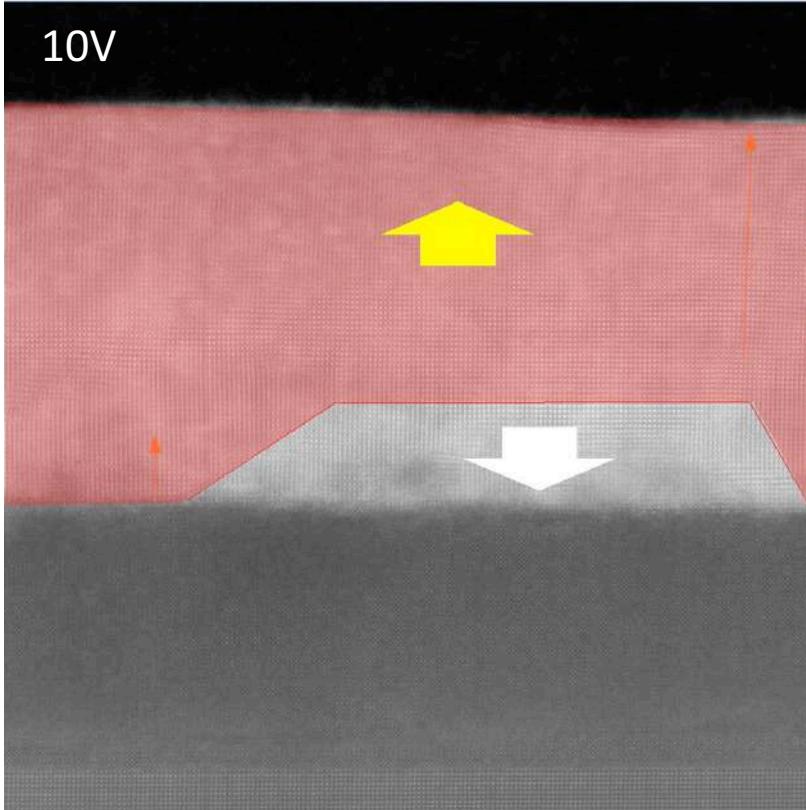
STEM image



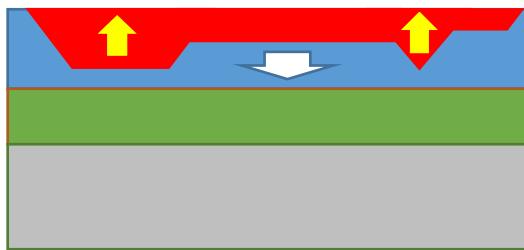
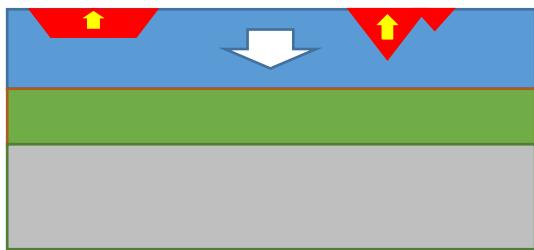
Observation of polarization in BiFeO₃ combining ABF and HAADF imaging



BFO domains



→ Trends with increasing positive bias →

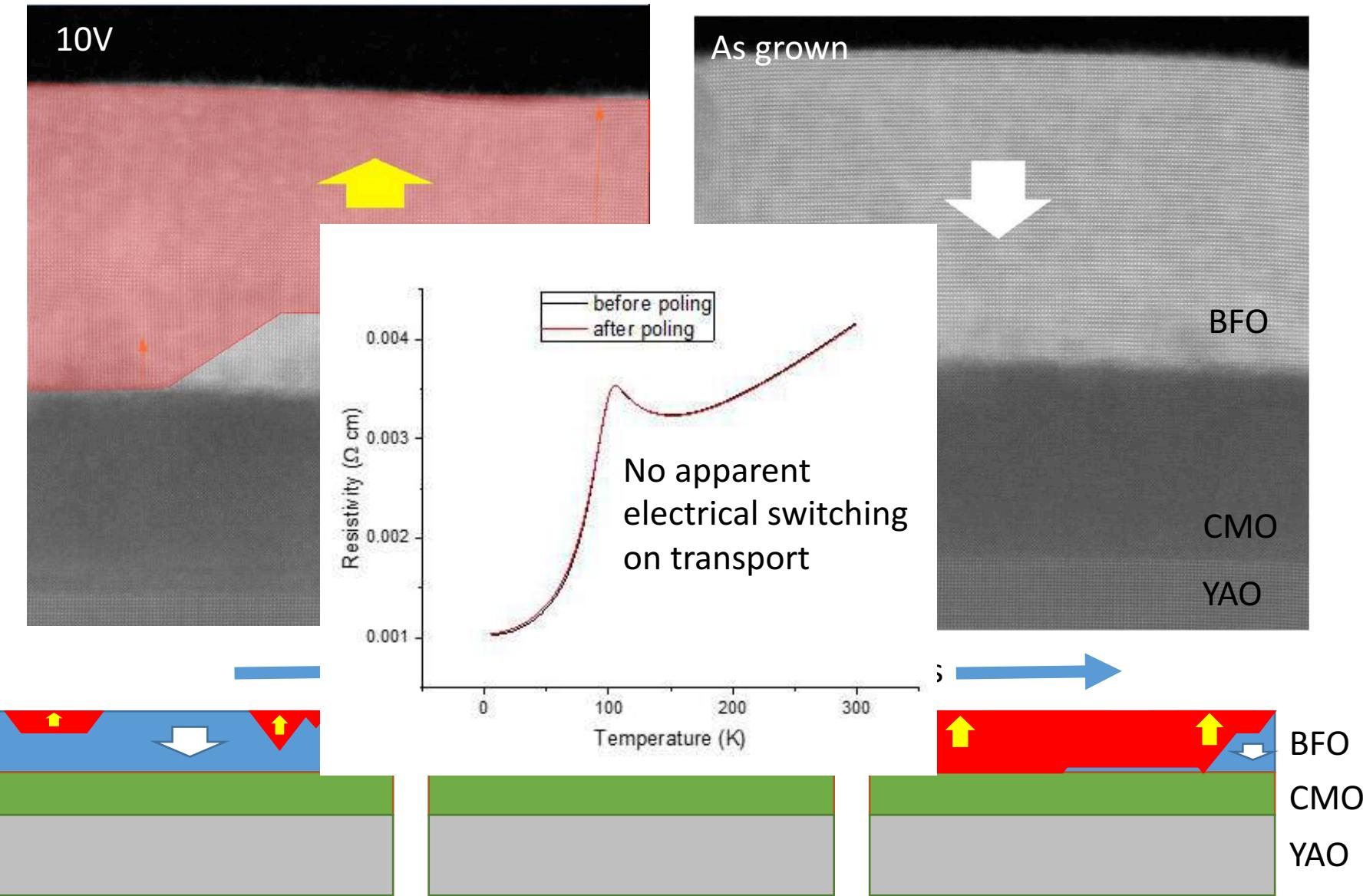


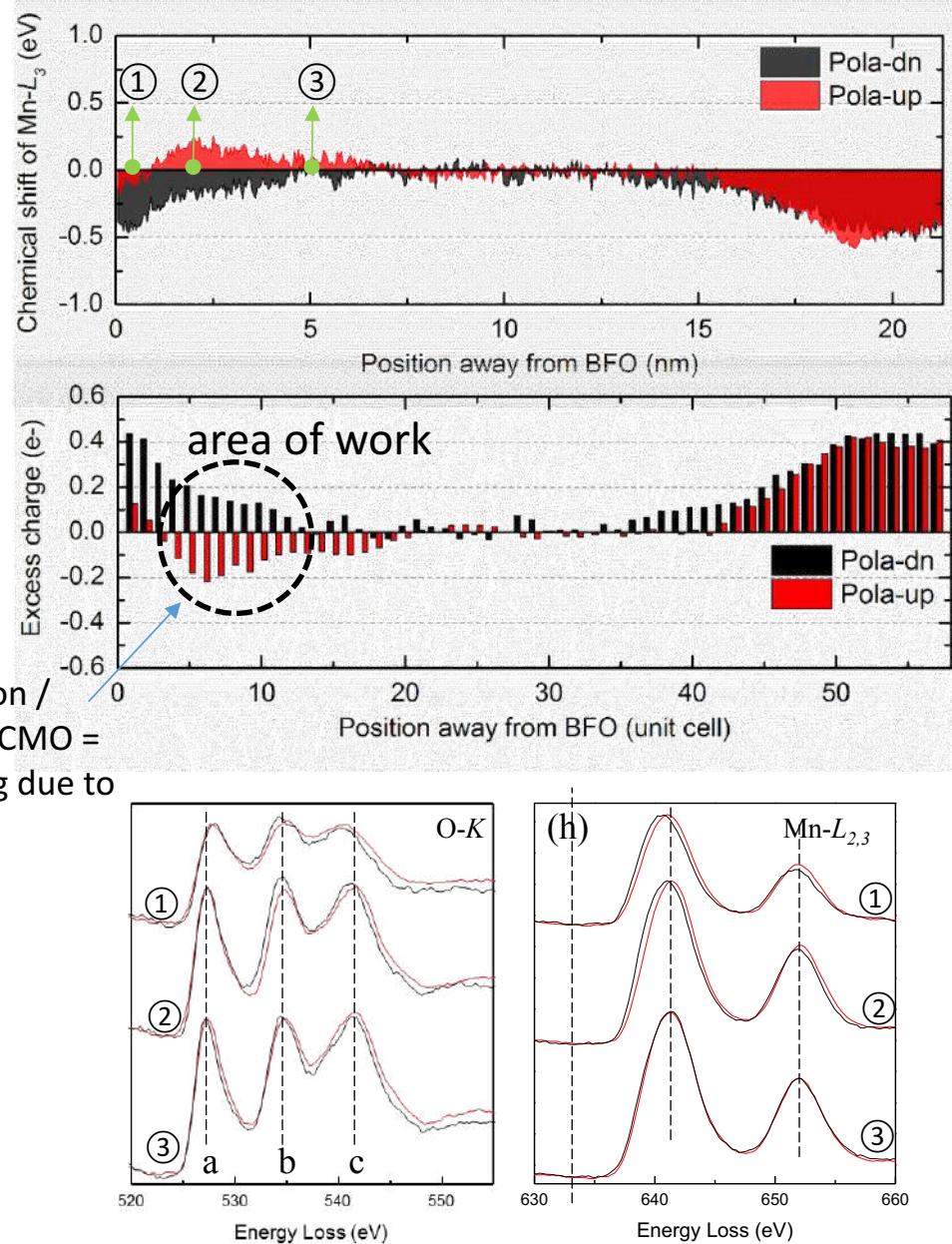
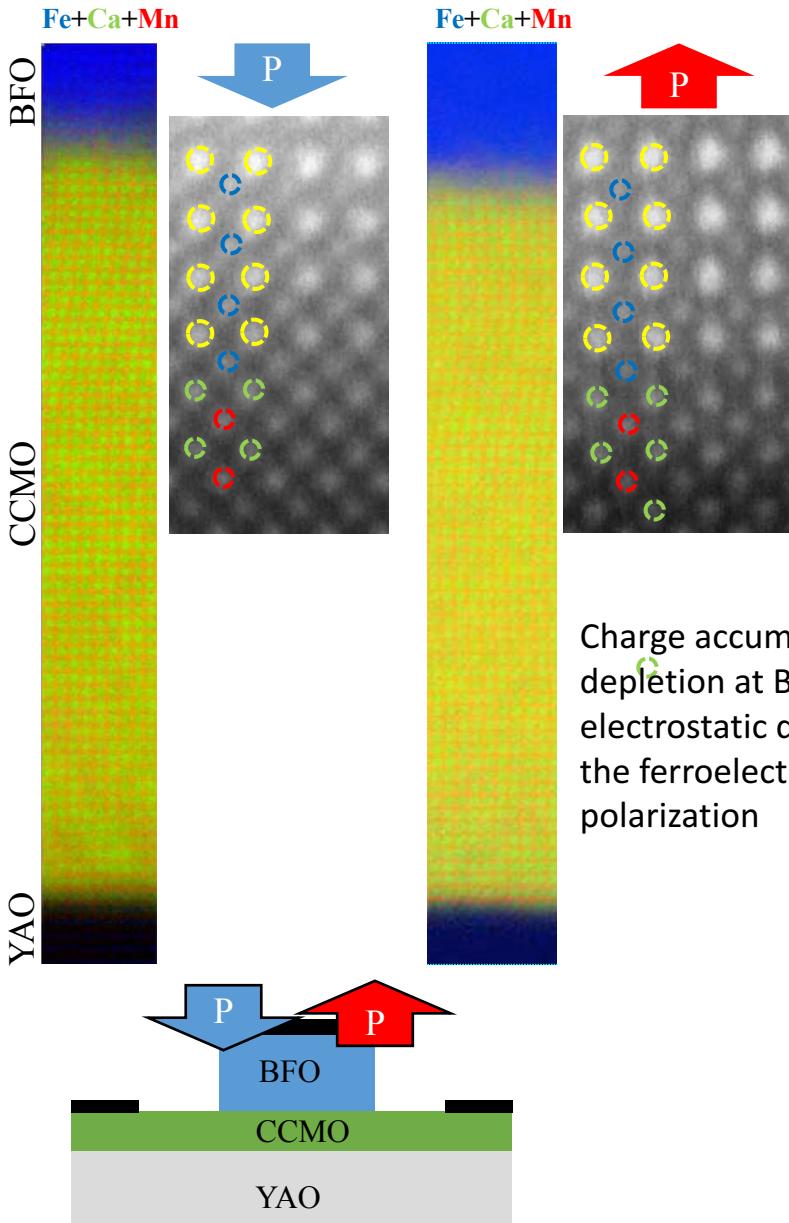
BFO

CMO

YAO

BFO domains



Charges Ce doped CMO (EELS Mn $L_{2,3}$) for BFO pol-up and pol-dn

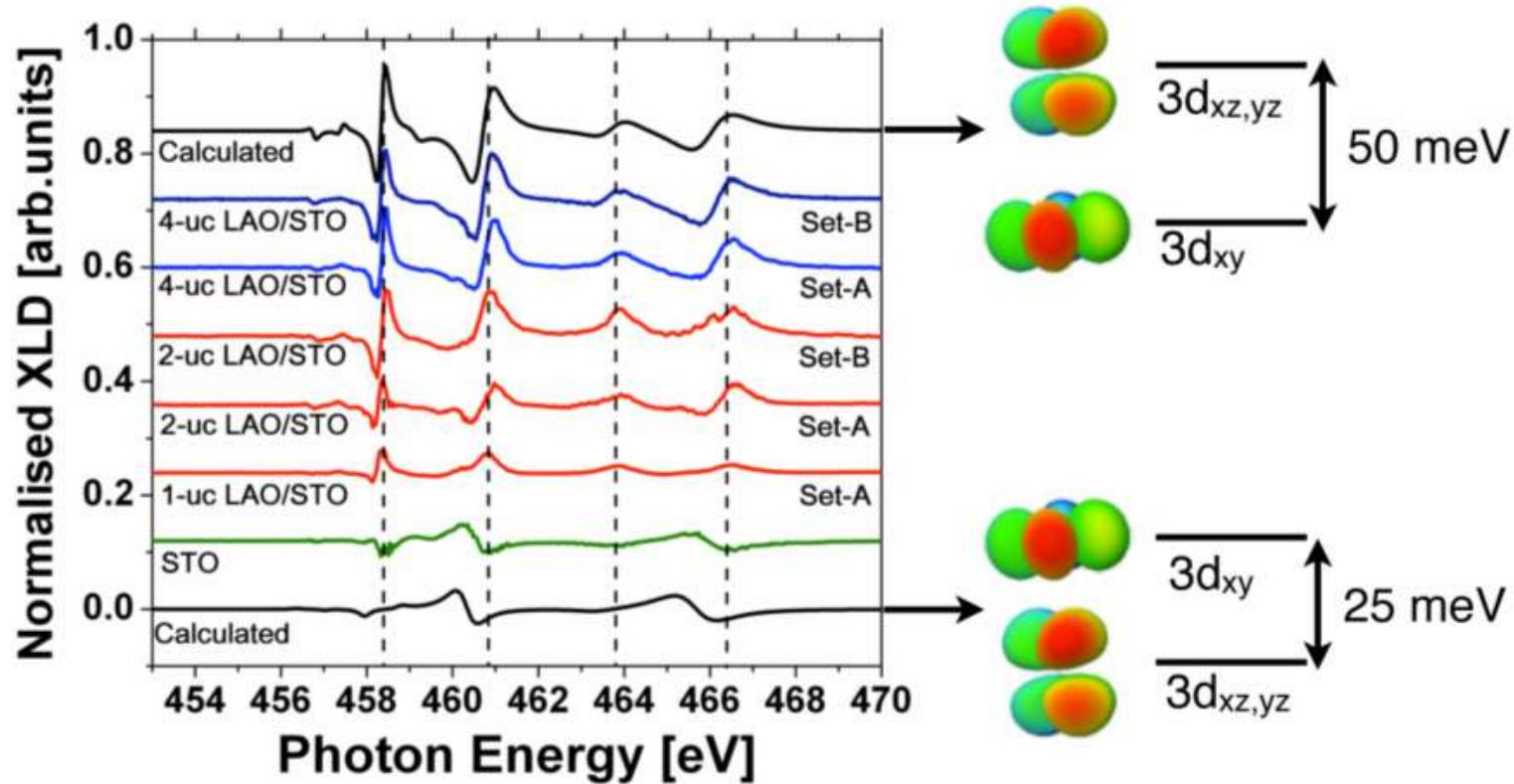
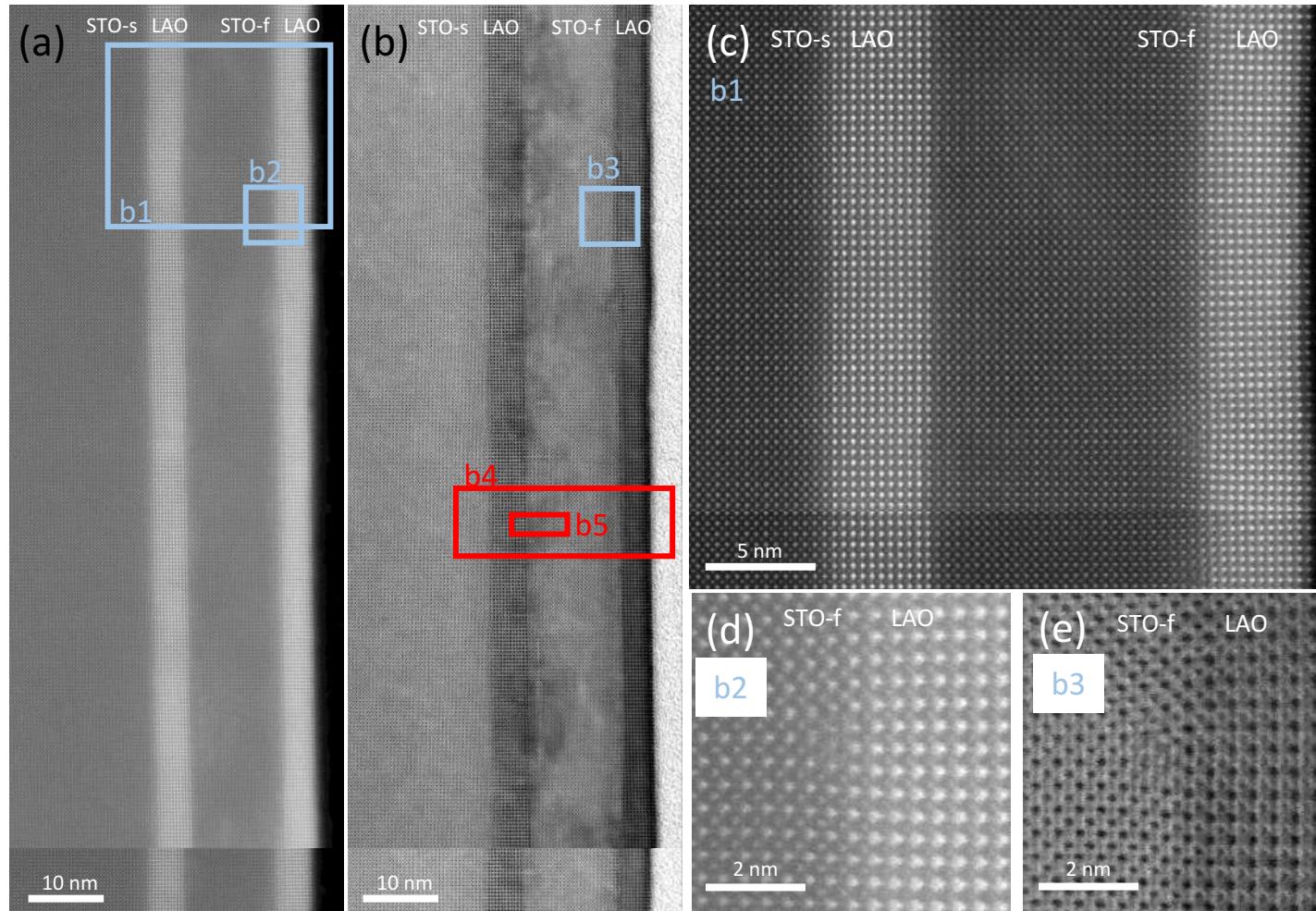
Orbitals hierarchy at the LaAlO₃/SrTiO₃ interface

Figure 1. X-ray linear dichroism (XLD) spectra around the titanium L_{2,3} absorption edge of SrTiO₃ (green line) and LaAlO₃/SrTiO₃ bilayers characterised by a LaAlO₃ thickness below (red lines) and above (blue lines) the critical thickness of four unit cells. Data from two sample sets are shown. Black lines are calculations which reproduce the data on STO (bottom) and LAO/STO (top) using multiplet atomic model calculations with point charge crystal field. On the right a schematic of the orbital splitting needed to reproduce the data is depicted, showing the inversion of hierarchy between in-plane and out-of-plane t_{2g} orbitals at the LaAlO₃/SrTiO₃ interface.

2D gаз are known to occur at LaAlO₃/SrTiO₃ interface

STO-substrate / LAO-f / STO-f / LAO-f

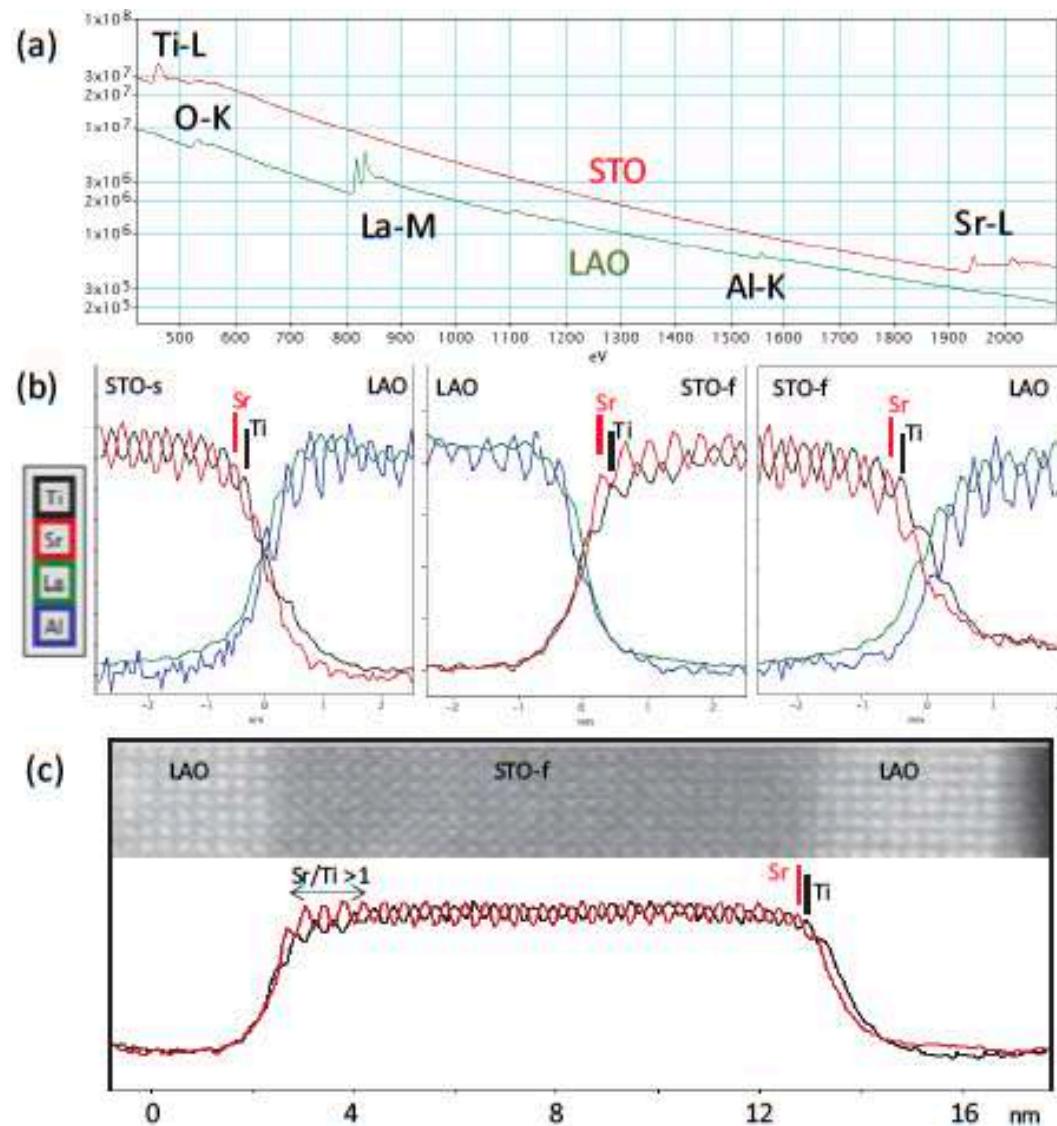


Coll. JM. Triscone
(U. Geneva)

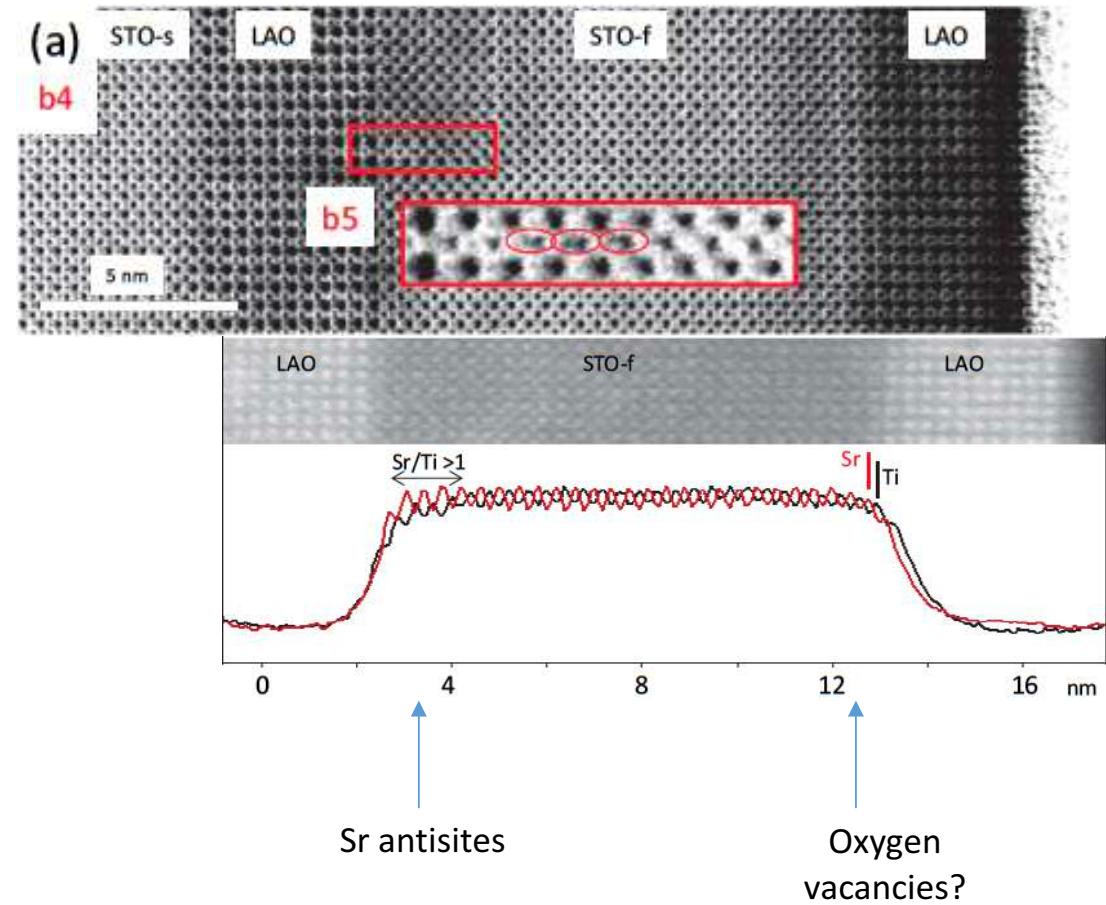
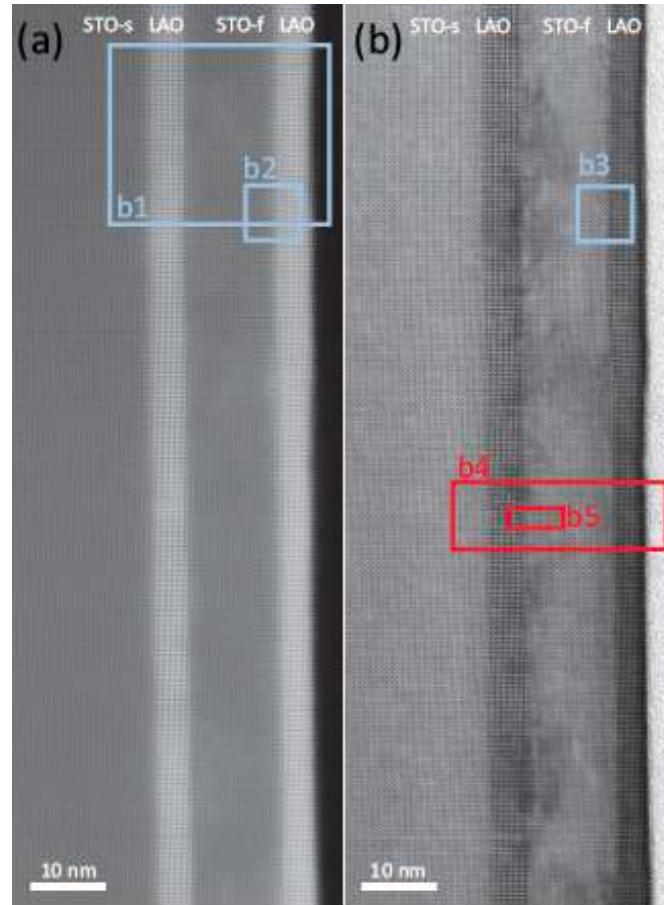
G. Tieri et al. submitted

More defects at the STO-film / LAO-film interfaces

Interface plane terminaison by elemental EELS mapping

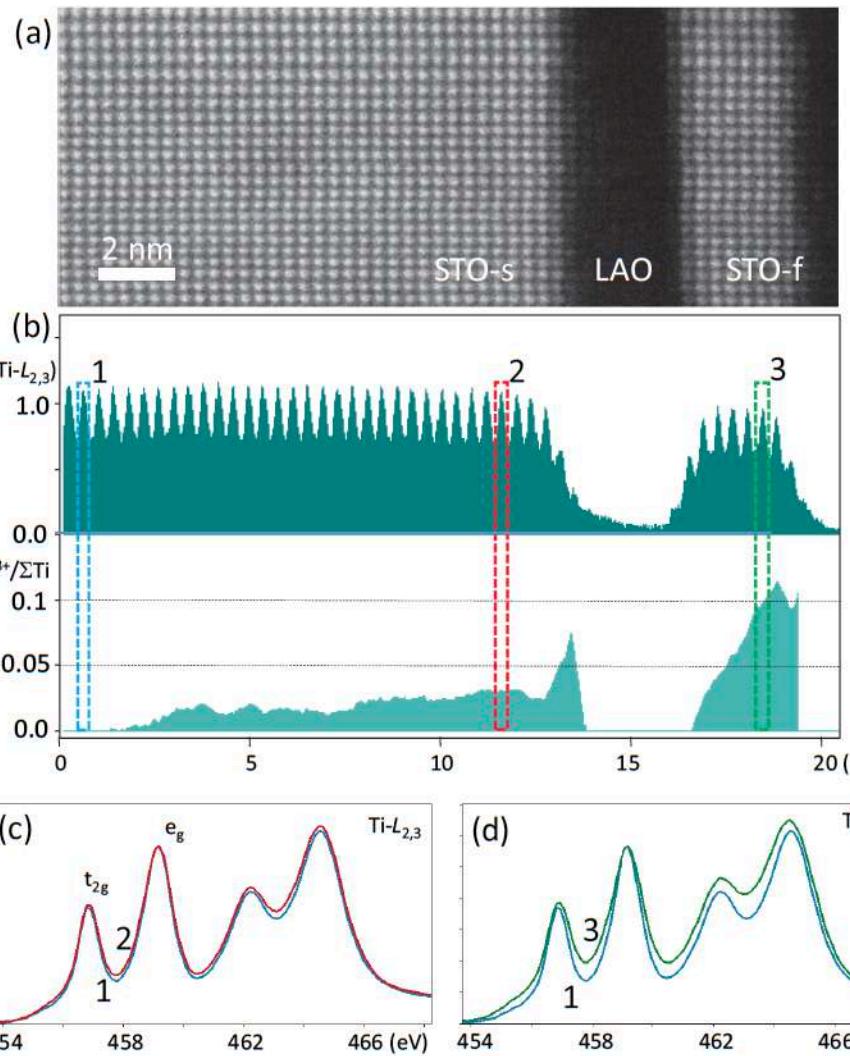


Defects and off stoichiometry



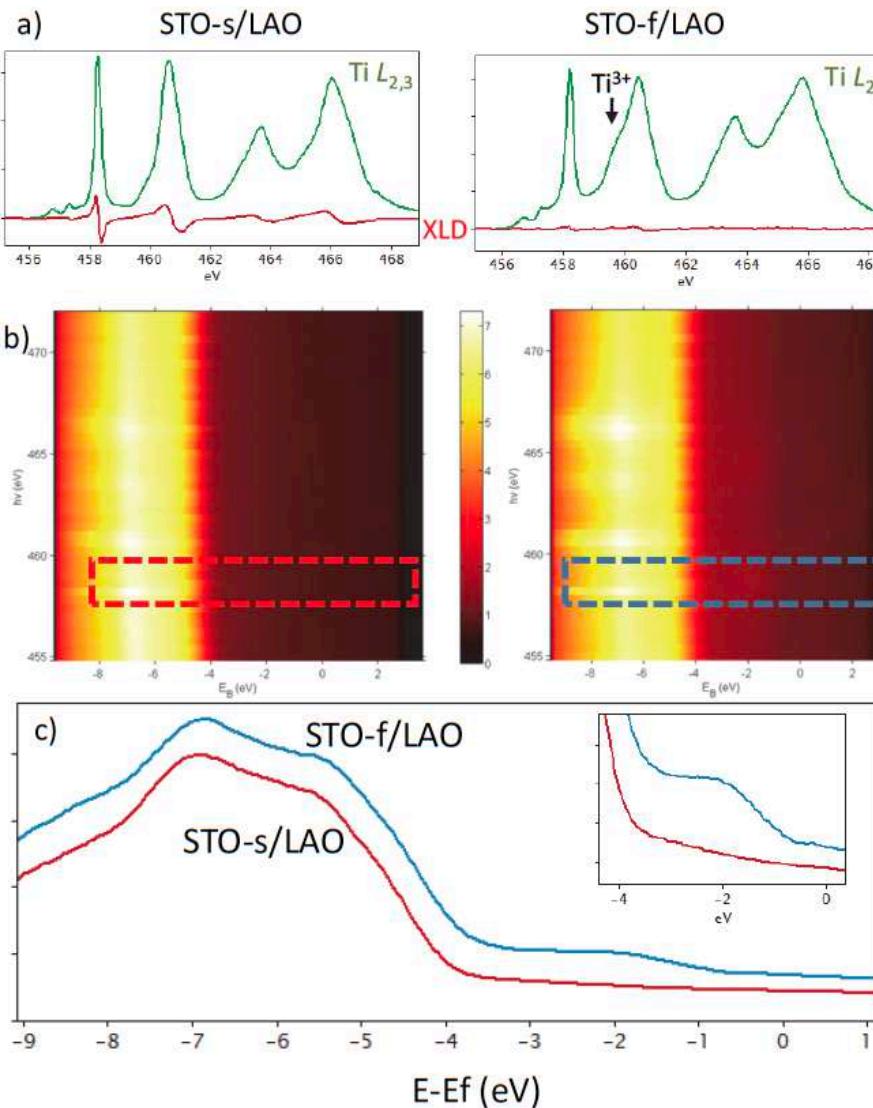
Mapping and quantifying Ti^{3+} states

Spectroscopic quantification of additional electrons at the interfaces



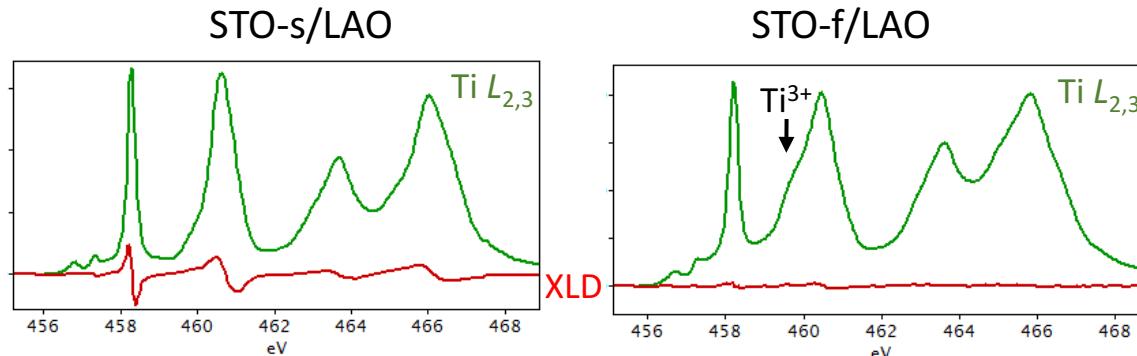
This is significantly higher than expected for the transport carrier densities, indicating that a significant amount of the carriers at these interfaces are localized

Spectroscopic probing of Ti³⁺ states by XAS and resonant PES



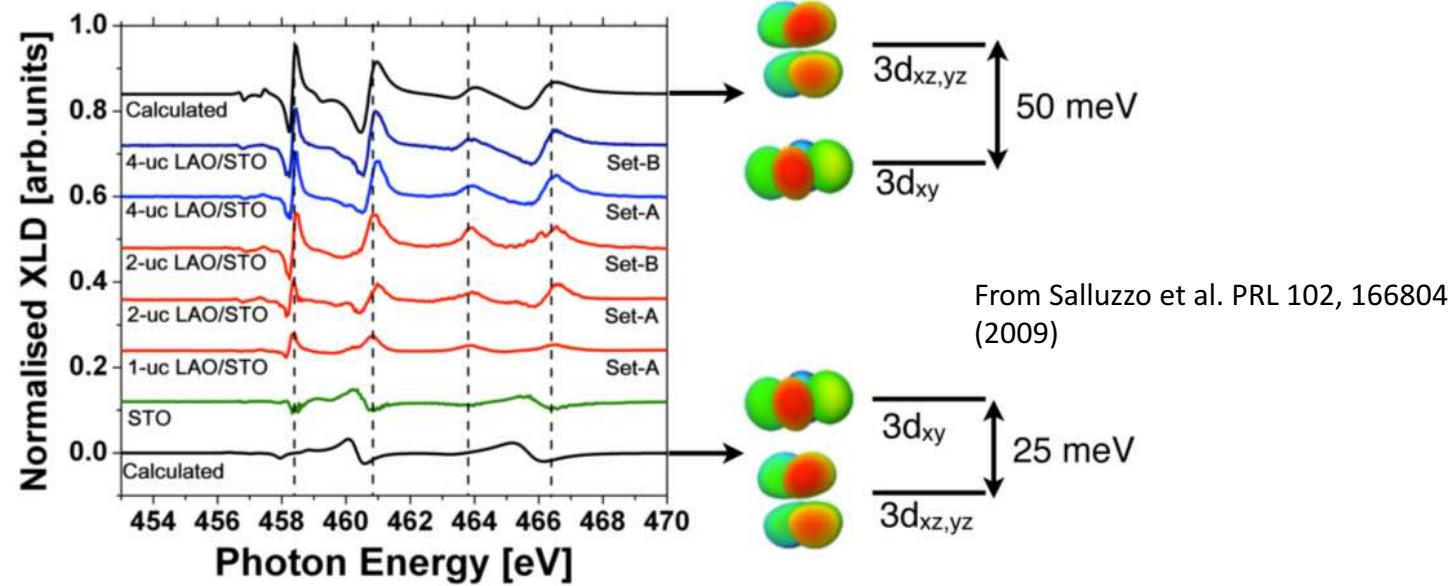
Larger Ti³⁺ signature at the STO-f/LAO interface

In gap states (charged V_O states?) at the STO-f/LAO interface

XLD Ti L_{2,3} for both interfaces

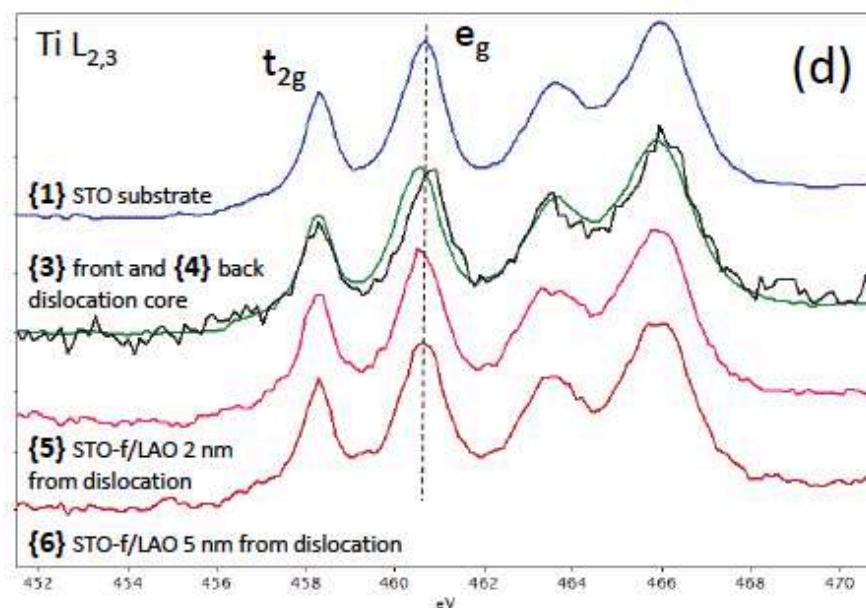
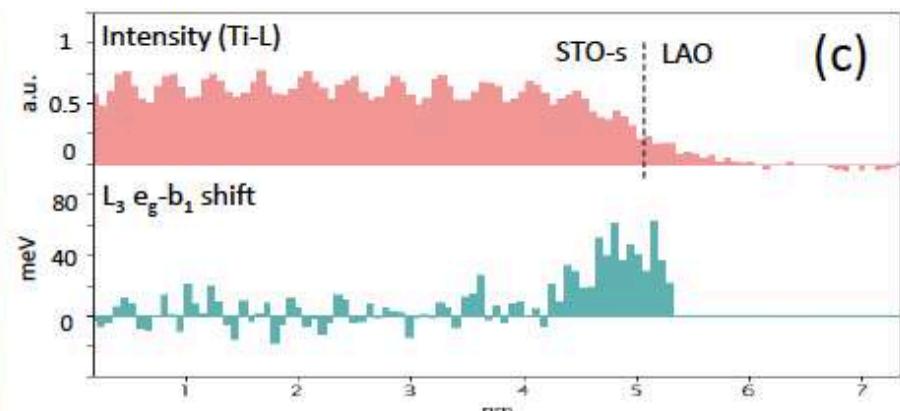
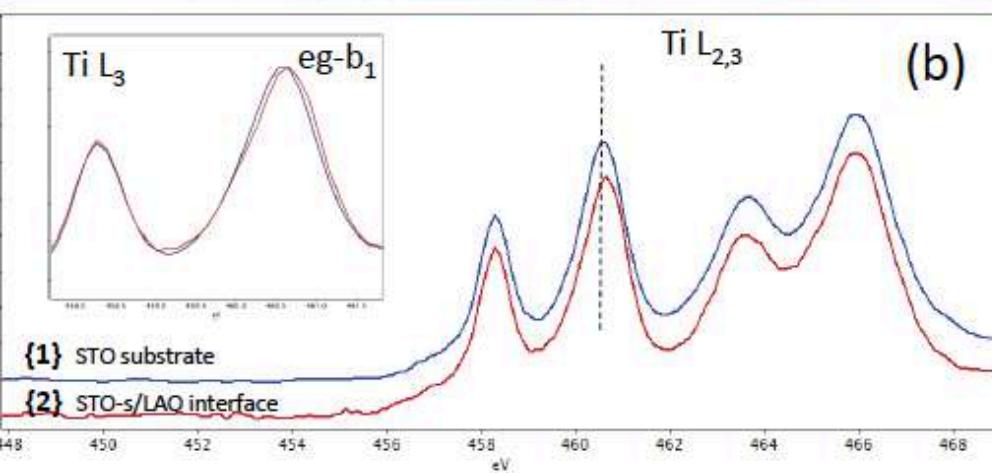
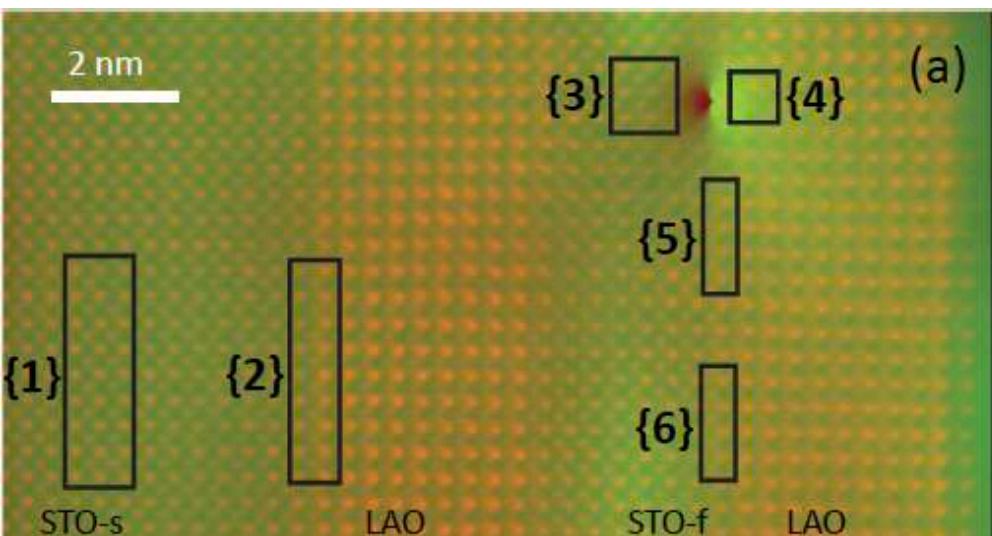
Intense dichroic signal at the STO-s/LAO whose sign confirms that the $3d_{xy}$ orbital lies at lower energy

No dichroic signal for the STO-f/LAO interface indicating that the orbital hierarchy is suppressed (or blurred)



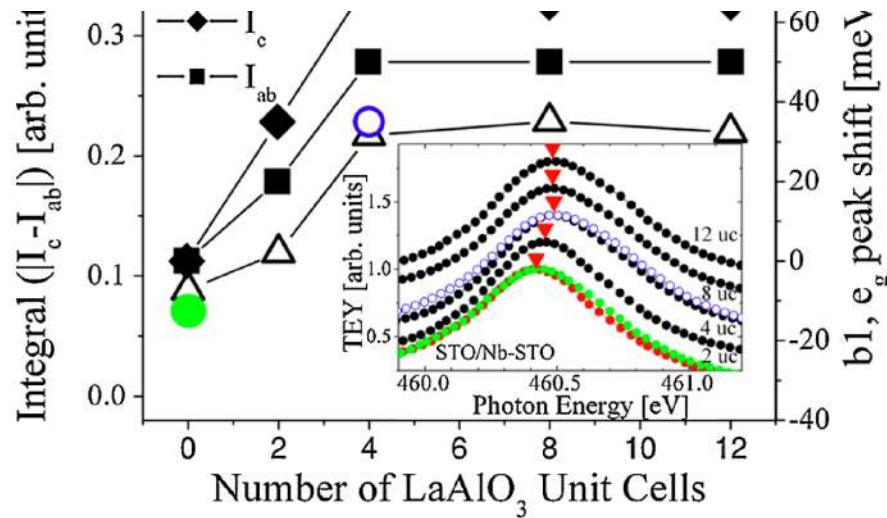
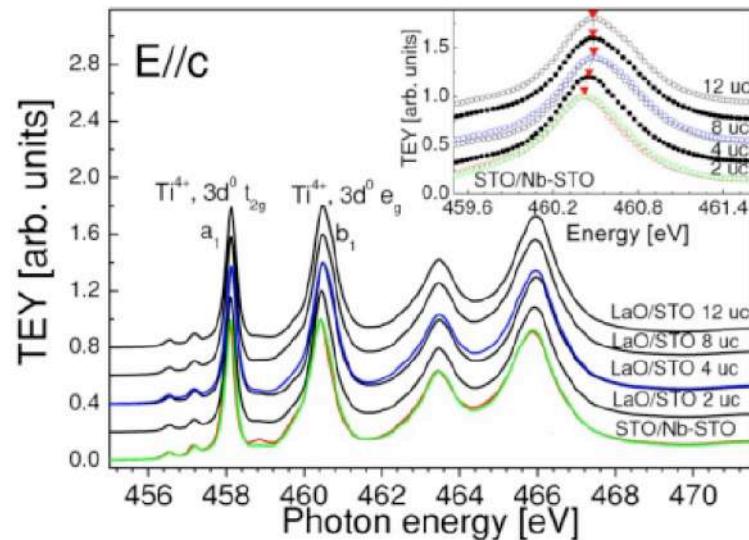
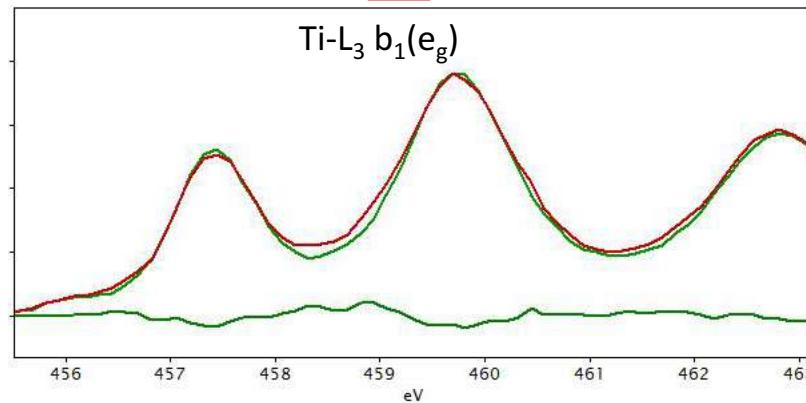
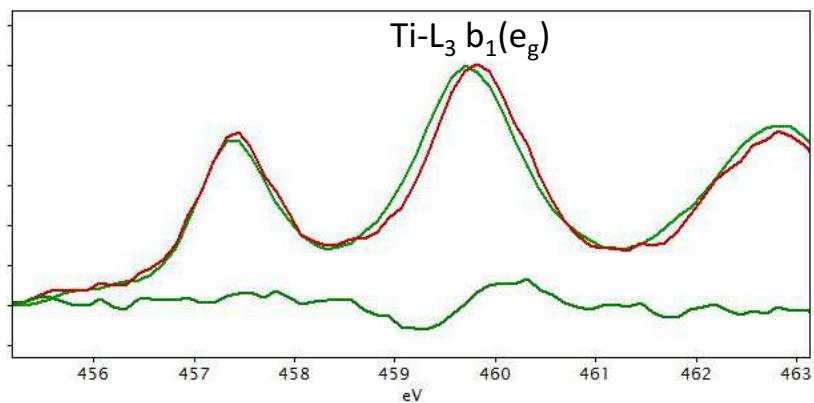
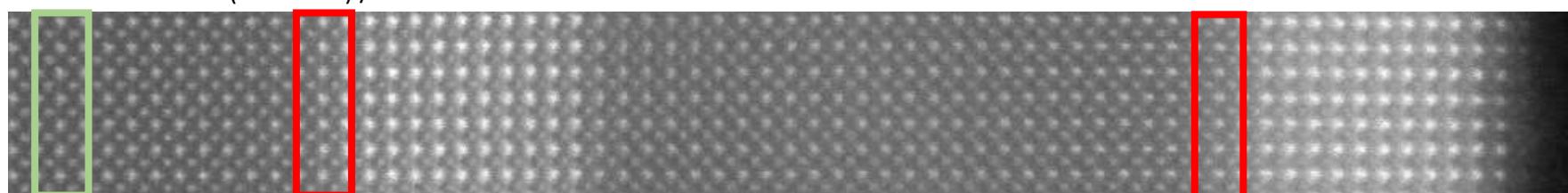
Different orbitals hierarchy at the different $\text{LaAlO}_3/\text{SrTiO}_3$ interfaces of the bi-layer ?

Isotropic EELS Ti L_{2,3} for both interfaces



Different orbitals hierarchy at the different $\text{LaAlO}_3/\text{SrTiO}_3$ interfaces of the bi-layer ?

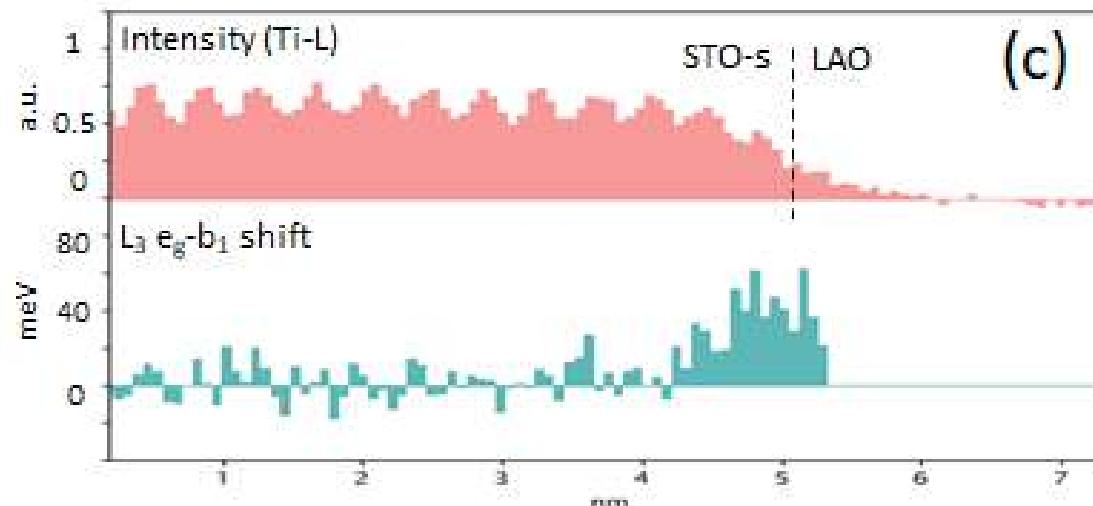
STO (substrate) / LAO



Real space extension of the « strong » orbital polarisation

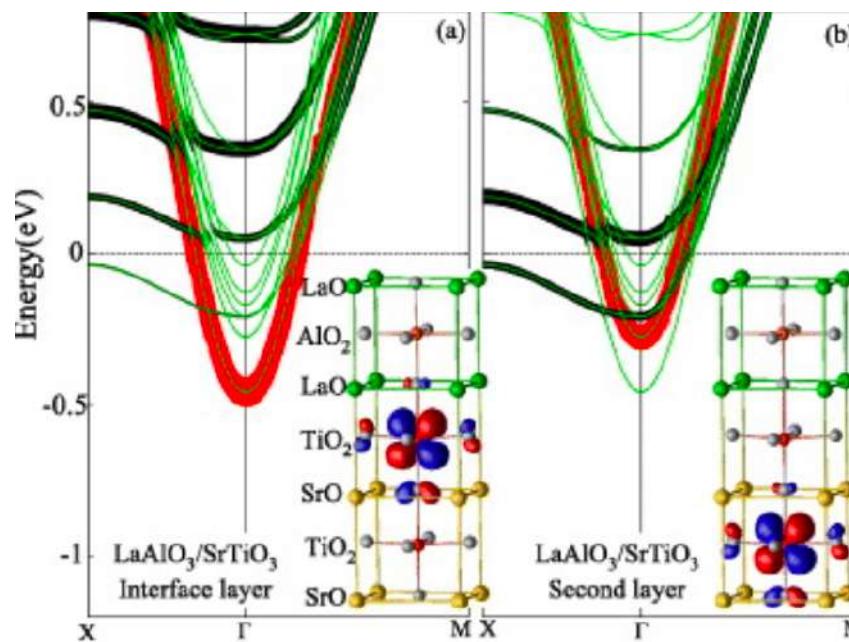
STEM-EELS

50 meV L3 eg shift is only measured for the Ti @ interface :
 SrO – TiO₂ - LaO



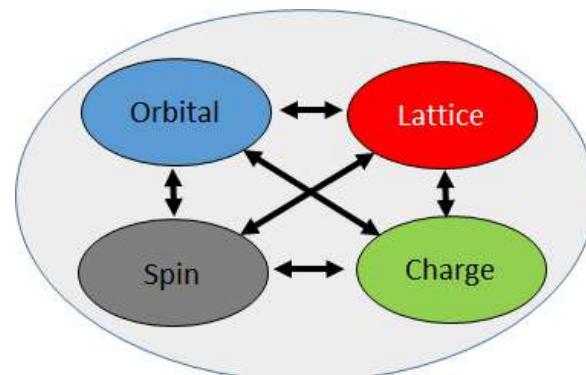
G. Tieri et al. submitted

Ab-initio

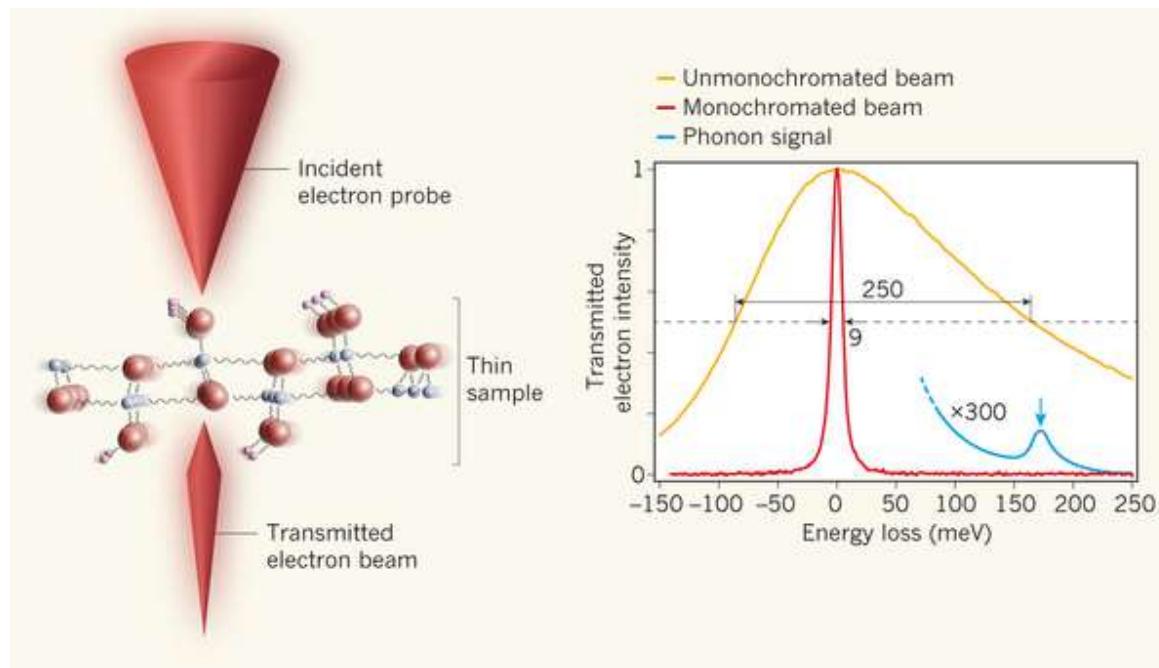


From Zhong et al., EPL 99 (2012) 37011

- 1) Octahedra rotation in $\text{LSMO}_3/\text{STO}_3$ superlattices
- 2) Charge control in manganite by ferroelectric switching in $\text{CaMnO}_3/\text{BiFeO}_3$ based Mott transistor.
- 3) Orbital ordering in $\text{LaAlO}_3/\text{SrTiO}_3$ bilayers
- 4) Low energy excitation in ABO_3 ($\text{B}=\text{TM}$) by EELS

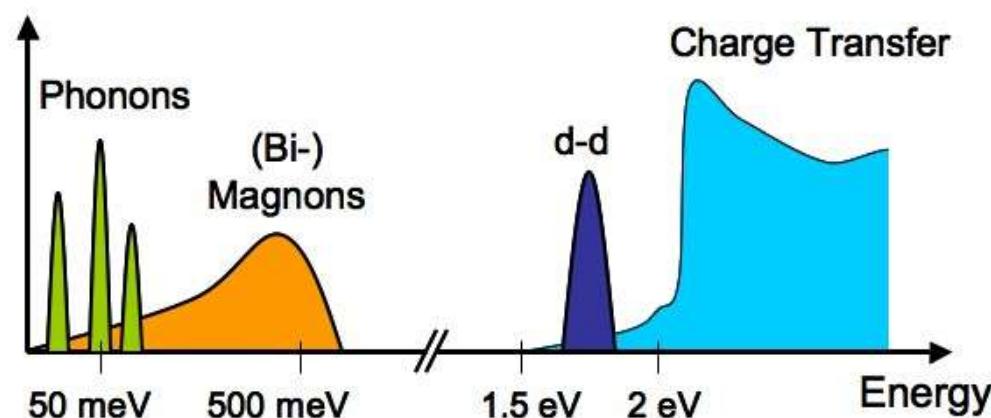


Higher spectral resolution STEM-EELS (keeping sub-nanometer spatial resolution)



LO phonon in h-BN, with a peak energy of 173meV

Krivanek et al. Nature **514**, 209 (2014).



High resolution EELS

50 -100 meV

5 meV

High Resolution imaging

1 Å

Sub-nm

Low temperature
(He & N₂)

NanoCL &
Light injection

1 meV

High currents

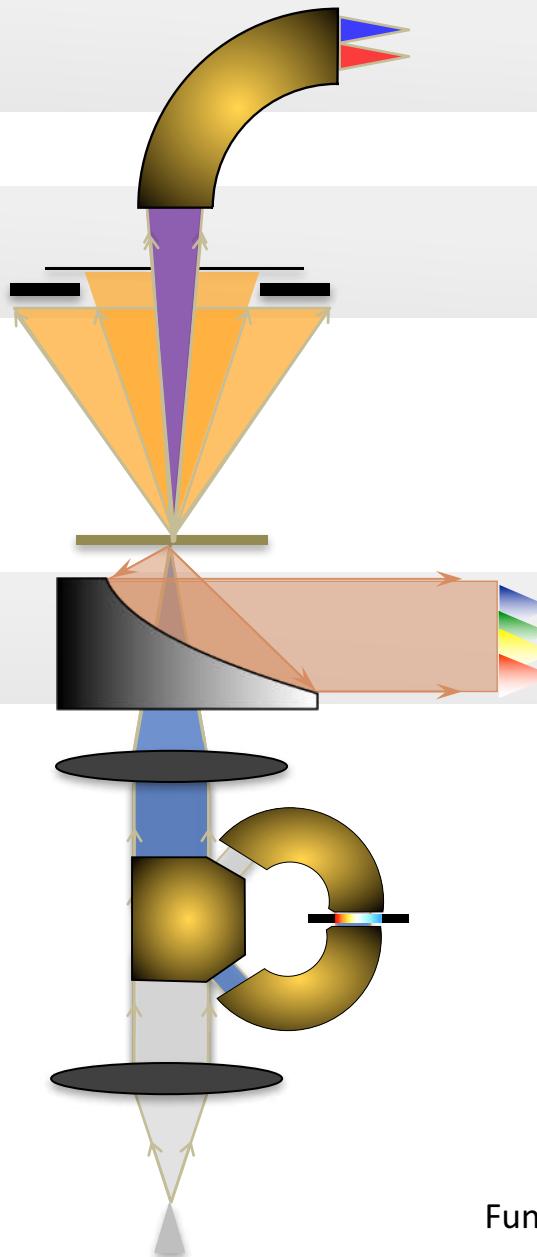
High resolution
monochromation

60 pA

10 pA @ 10 meV
(50 pA @ 50 meV)

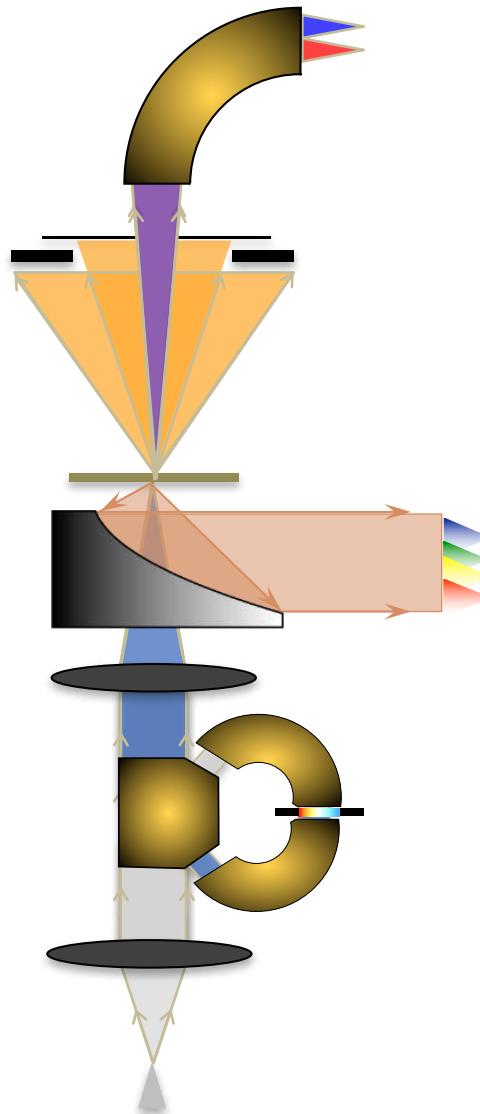
Other specs, especially for reciprocal
space measurements

Funded in 2010



Equipement d'excellence
TEhPOS
Transmission Electron Microscopy at Palaiseau Orsay Saclay

The CHROMATEM microscope



- NION Hermes 200M
 - 30-200 kV
 - Side entry stage
 - High-resolution EELS spectrometer

- HennyZ sample holders:
 - Variable temperature (LN_2 -600°C)
 - Electrical contacts

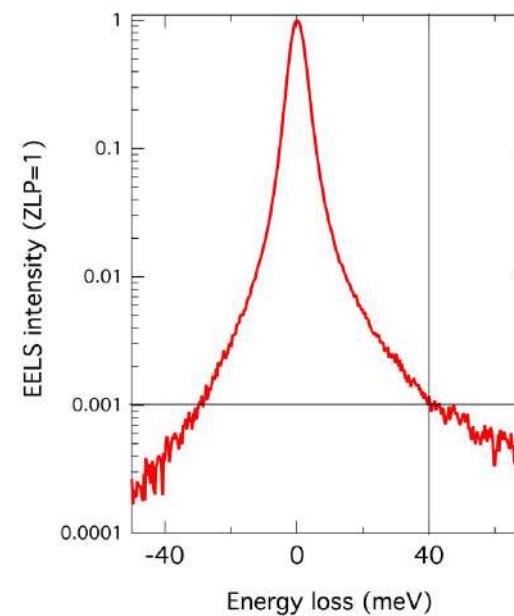
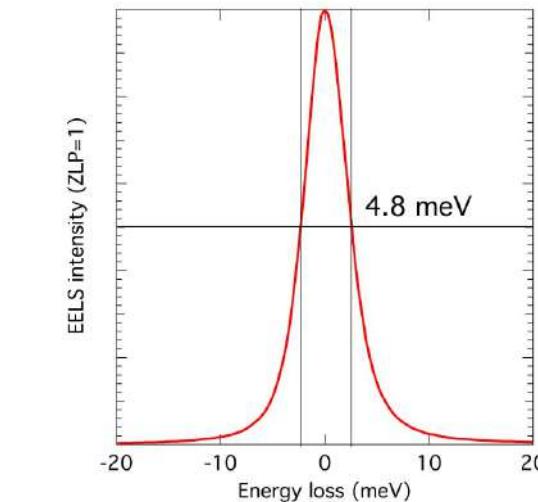
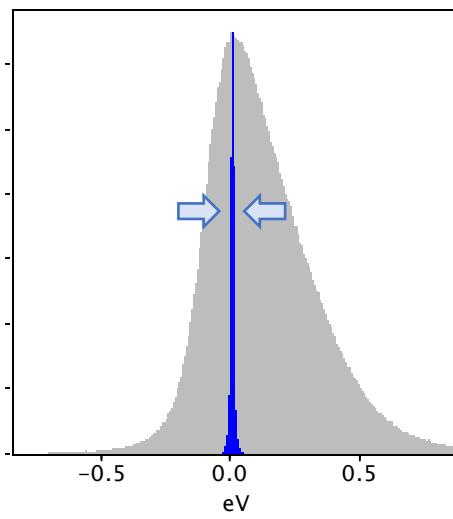
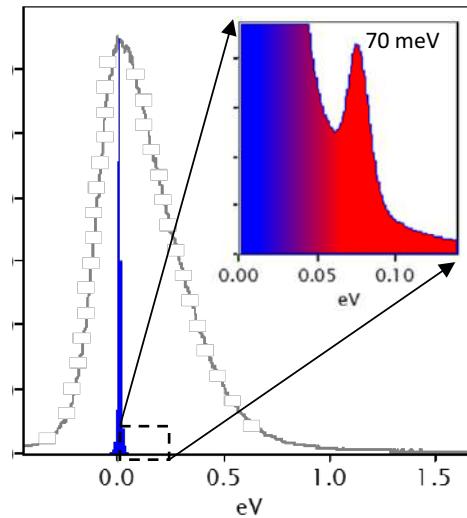
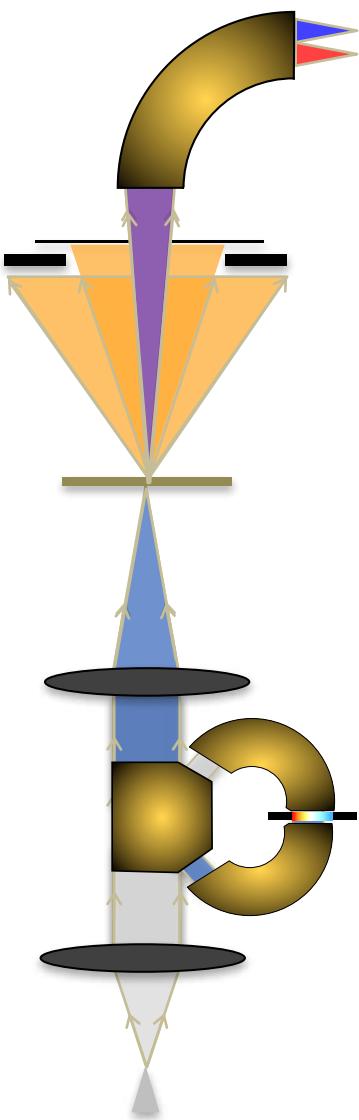
- Attolight light injection/detection system



© Photo credit – Cyril FRESILLON / LPS / CNRS Photothèque

Delivery early 2018

CHROMATEM: providing a privileged access to very low losses down to the IR



“Classical” localized plasmons are now routinely mapped by STEM –EELS or STEM CL

A very lively field

Wherever space and energy variations are entangled

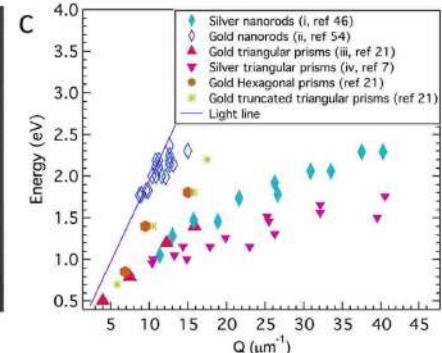
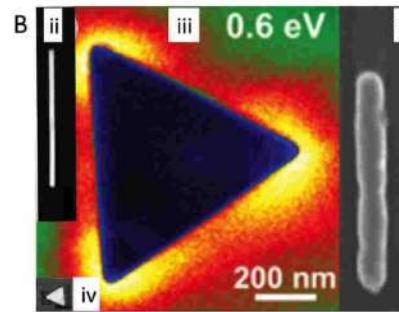
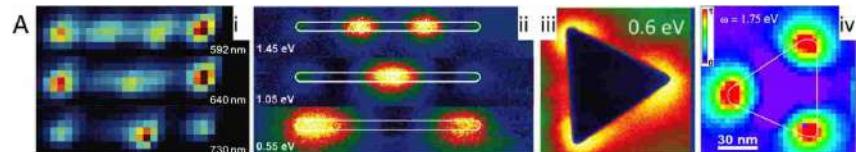
Nelayah, J. et al. Nat. Physics (2007)

Vesseur, E. J. R., et al. Nano Lett. 7, 2843–2846 (2007).

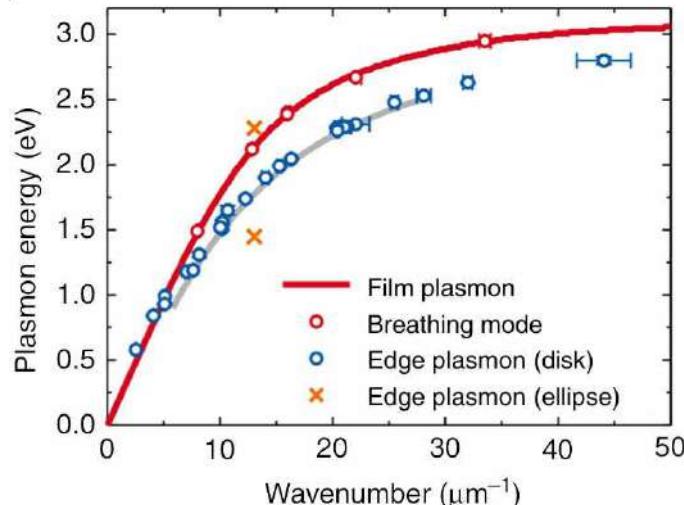
Gu, L. et al. Phys. Rev. B 83, (2011).

Nelayah, J. et al. Nano Lett. 10, 902–907 (2010).

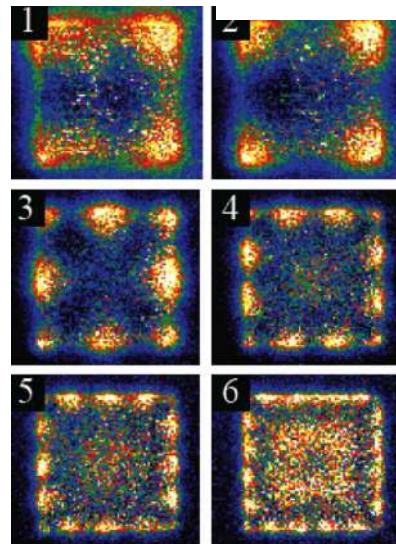
Rossouw, et al. Nano Lett. 11, 1499–1504 (2011).



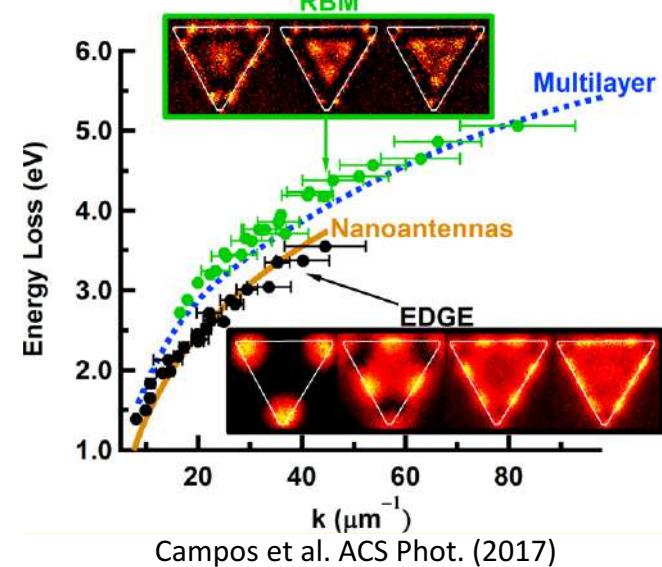
Plasmon classification



Schmidt, F.-P. et al. Nature Communications 5, (2014).



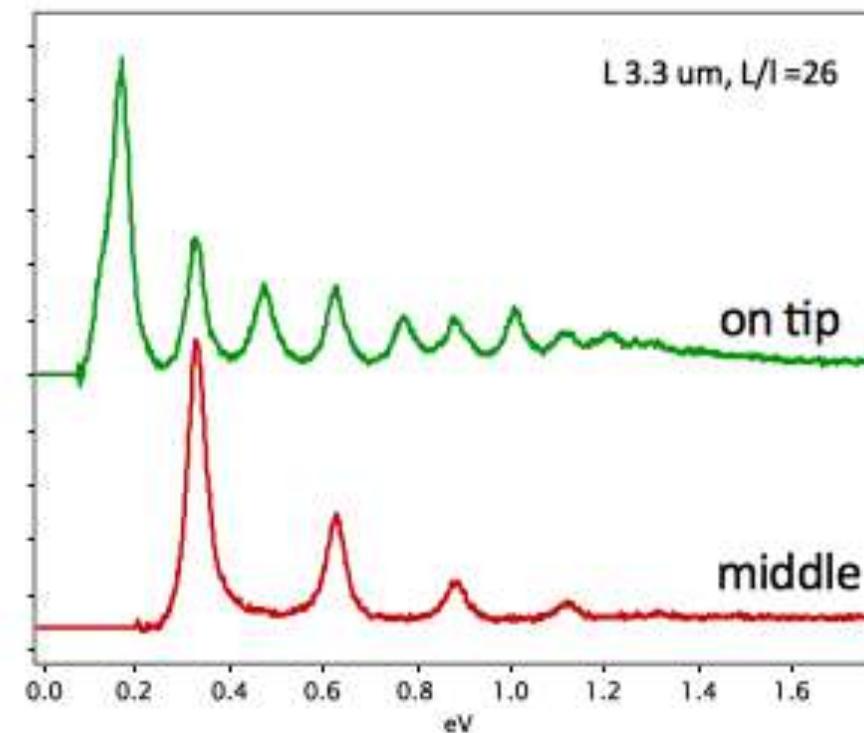
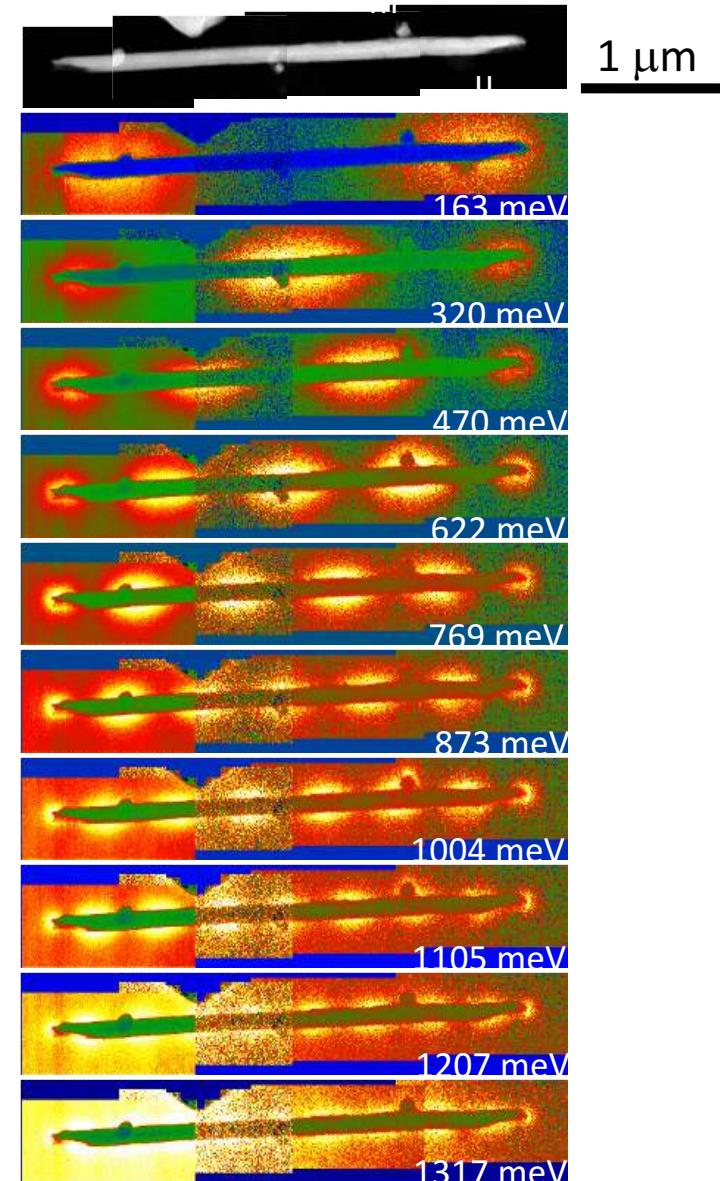
Bellido et al. ACS Phot. (2016)



Campos et al. ACS Phot. (2017)

Spatially resolved low energy excitation

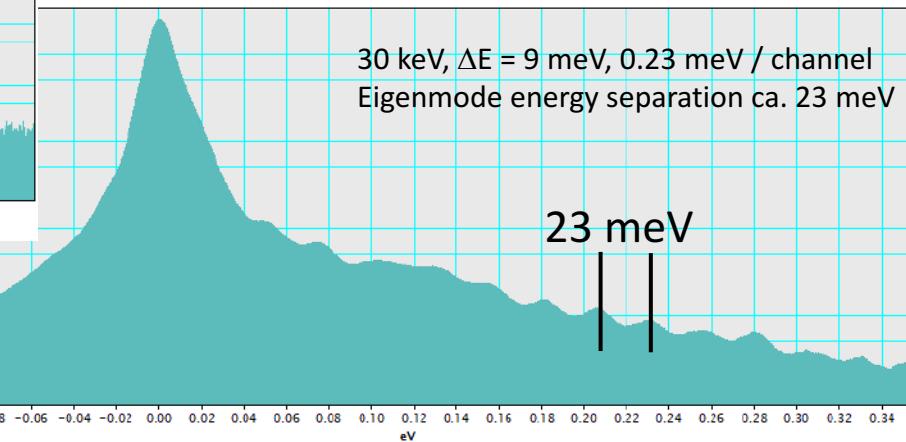
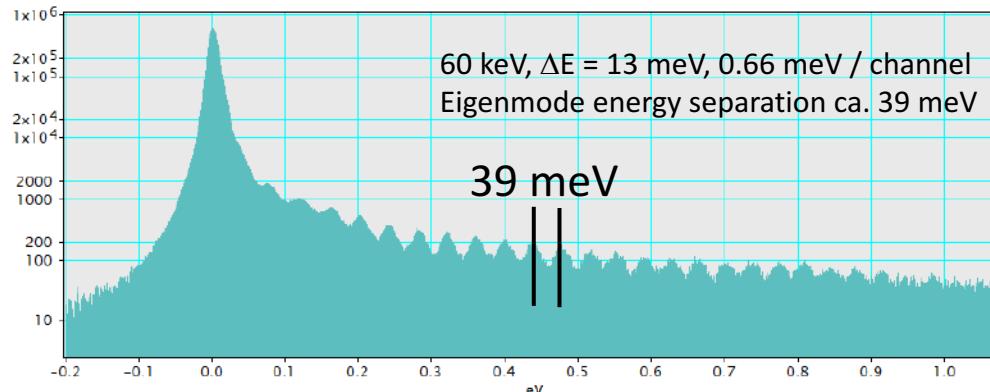
THz plasmons: the example of plasmon modes in a Cu nanowire



A. Gloter,
M. Kociak (LPS-Orsay)
SY Chen (NTUST)
S. Song (NCHU)
Unpublished

EELS in CHROMATEM microscope -> down to 5 meV energy resolution @ 30 keV

- Plasmon modes on « long » metallic nano-rods

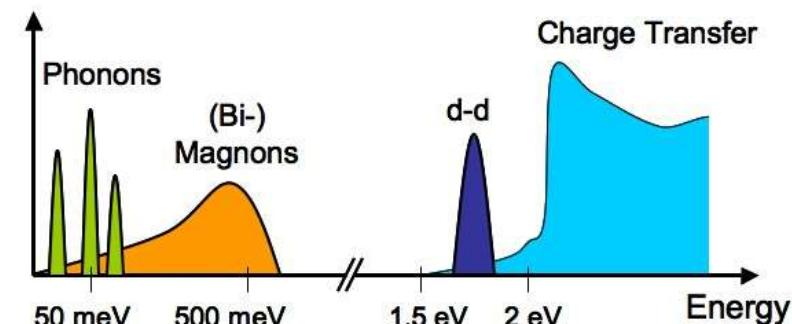


E. Tseng, A. Gloter, L. Tizei, X. Li, O. Stéphan, M. Kociak
and NTUST-TW collaborators (in preparation 2019)

- « TMO oxides »

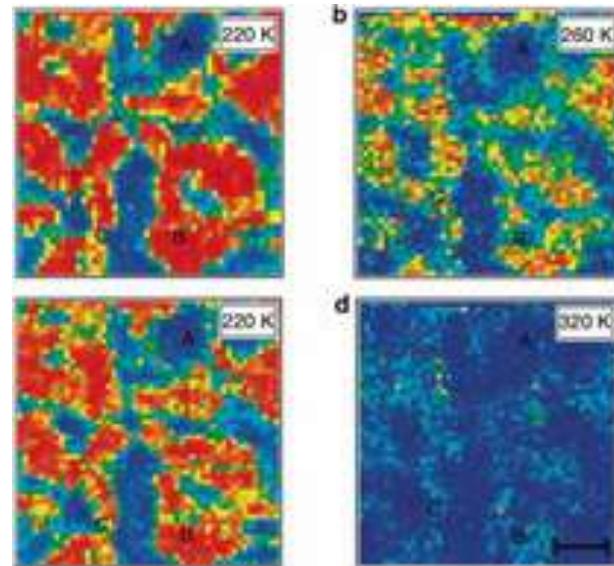
Access (*and map with nanometer scale resolution !!!*) the low energy electronic excitations in transition-metal oxide based nanostructures

Phonons, magnons ??, d-d, p-d, plasmons



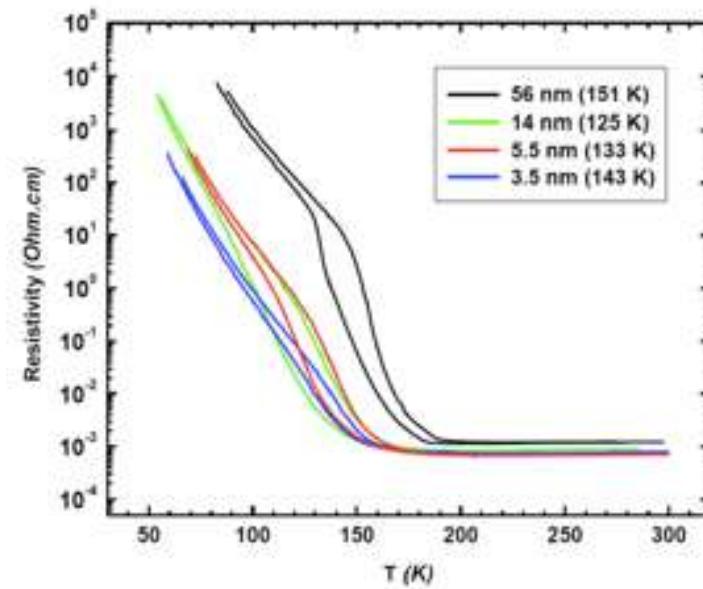
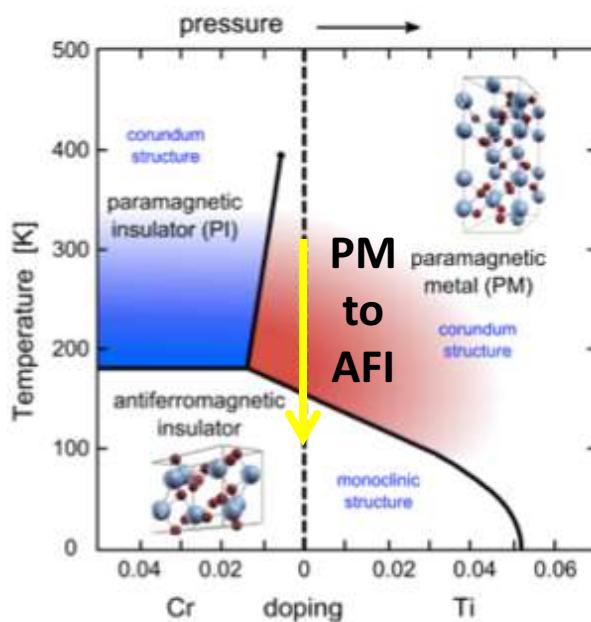
Ament et al. Rev. Mod. Phys., 83, 2011
RIXS review

MIT transitions at the nanoscale: V₂O₃ as an archetypal MIT system



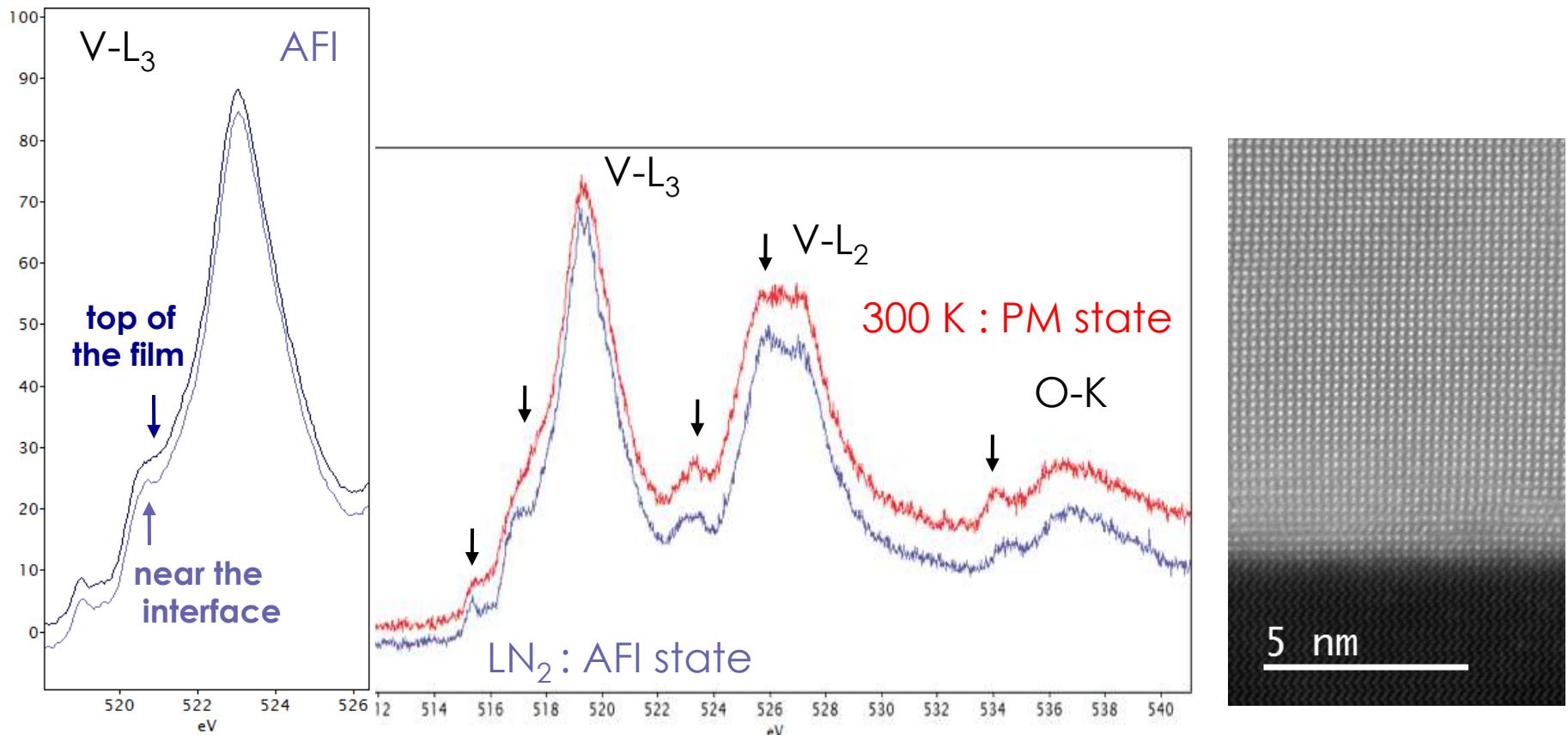
MIT in the V₂O₃ thin films at 137 K

Lupi, S. et al. *Nature Communications* **1**, 105 (2010).



L. Dillemans et al. *APL* **104** (2014) 071902

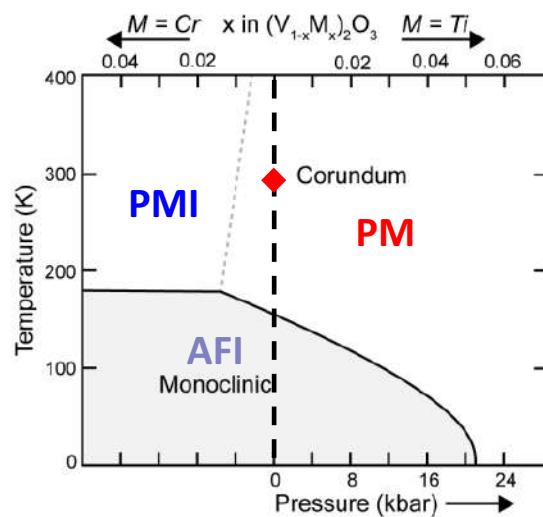
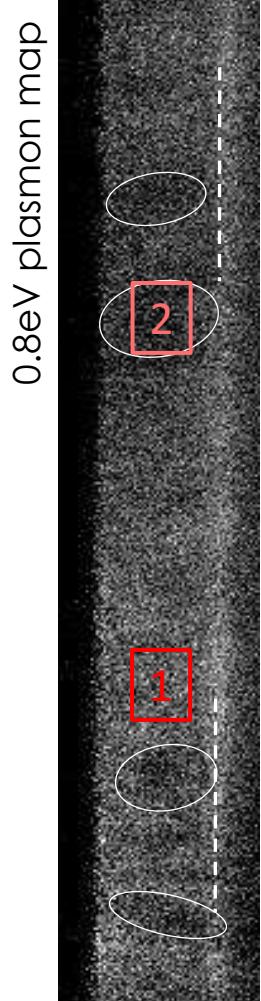
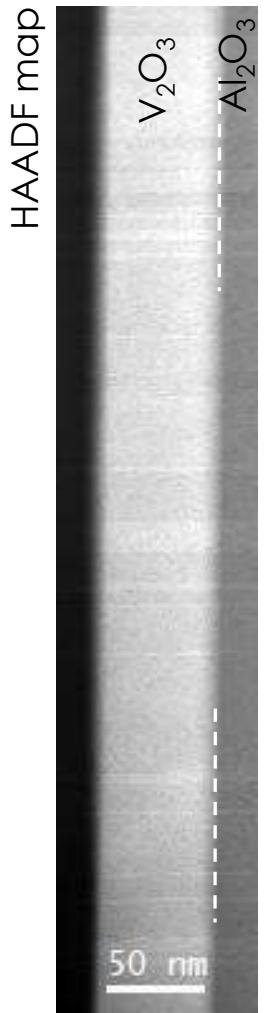
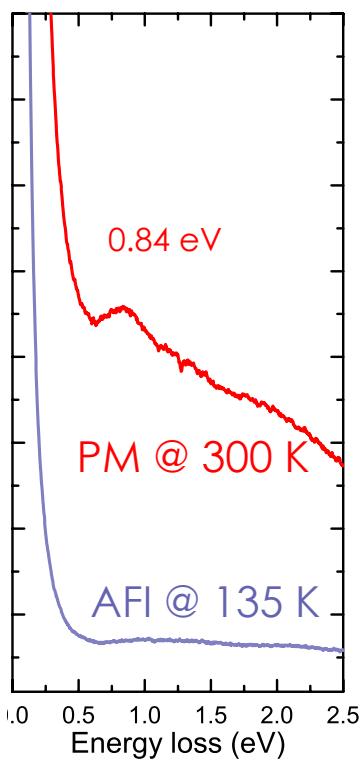
Thermally-induced metal-insulator transition Core-loss excitations probed through the low-T MIT



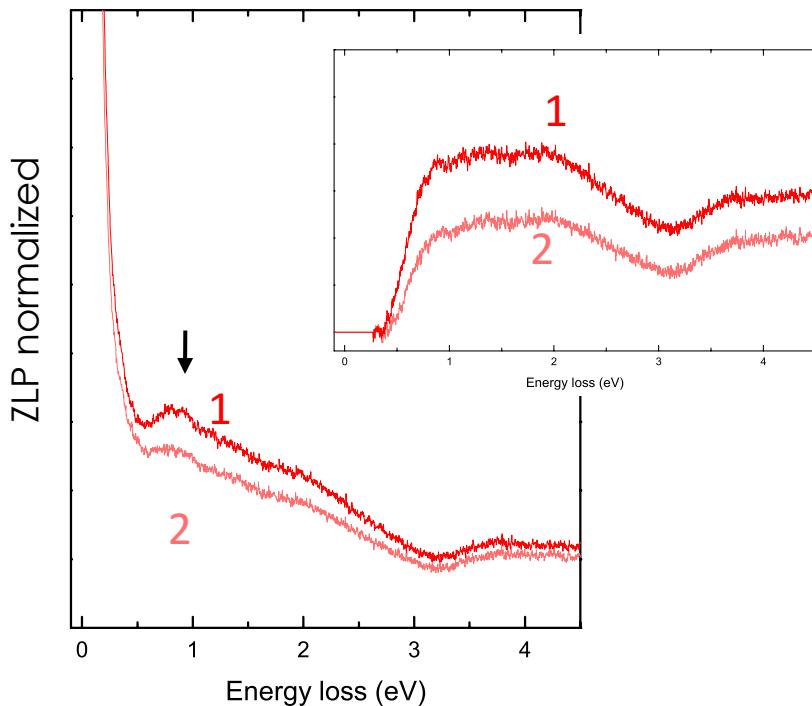
L. Bocher & X. Li, unpublished (2018)
 Coll. M. Menghini et al. KU Leuven
 See Laura Bocher's talk tomorrow

➔ PM and AFI spectroscopic signatures
 as in Abe et al. JJAP 38 (1999)

Mapping the metallic state



→ local inhomogeneities in the film



Raw data

- (80x400) px with 12ms/px < 7min in total

Acknowledgements

STEM group @ LPS-Orsay

<https://www.stem.lps.u-psud.fr/>

Main collaborators are:

J.M. Triscone et al University of Geneva

V. Garcia, S. Fusil, M. Bibes, A. Barthelemy et al, UMPHY CNRS-Thales

P. van Haken, Max Planck Institute for Solid State Research, Stuttgart

L. Liz-Marzan, CICBiomaGUNE, San Sebastian

A few reviews :

C. Colliex et al., Capturing the signature of single atoms with the tiny probe of a STEM,
Ultramicroscopy, 123 (2012) 80

A. Gloter et al., Atomically resolved mapping of EELS fine structures, *Materials Science in Semiconductor Processing*, 65 (2017) 2

Stephen J. Pennycook, Peter D. Nellist, Scanning Transmission Electron Microscopy: Imaging and Analysis, Springer Science & Business Media (2011)



ESTEEM 3 in a nutshell



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